

Heavy Quark Masses (from QCD Sum Rules) and their impact on the muon $g-2$

Pere Masjuan
Universitat Autònoma de Barcelona
(masjuan@ifae.es)

Work ongoing in collaboration with
Jens Erler and Hubert Spiesberger
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UAB
Universitat Autònoma
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Outline

- Motivation and Introduction
- Using Sum Rules to extract m_Q
 - overview
 - our proposal for *charm* and *bottom*
- *Impact on the muon $g-2$*
- Conclusions and outlook

Motivation: why precise m_q ?

$$\text{Higgs decay} \sim \overline{m_b}(M_H)^2$$

$$\Gamma(B \rightarrow X_u l \nu) \sim G_F^2 m_b^5 |V_{ub}|^2$$

$$\Gamma(B \rightarrow X_c l \nu) \sim G_F^2 m_b^5 f(m_c^2/m_b^2) |V_{cb}|^2$$

$$\begin{aligned} B &\rightarrow K^{(*)} l l \\ B &\rightarrow D^{(*)} l \nu \end{aligned} \quad (\text{pQCD contributions on FFs depend on } m_q)$$

Yukawa unification

[Baer et al '00]

$$\frac{\delta m_b}{m_b} \sim \frac{\delta m_t}{m_t} \quad \text{if } \delta m_t \sim 1\text{GeV} \Rightarrow \delta m_b \sim 25\text{MeV}$$

Motivation: why precise m_Q ?

Υ -spectroscopy

$$m(\Upsilon(1S)) = 2M_b - C\alpha^2 M_b + \dots$$

Lattice QCD

$$M_{H^{(*)}} = m_h + \bar{\Lambda} + \frac{\mu_\pi^2}{2m_h} - d_{H^{(*)}} \frac{\mu_G^2(m_h)}{2m_h} + \mathcal{O}(m_h^{-2})$$

QCD Sum Rules

$$\int \frac{ds}{s^{n+1}} R_q(s) \sim \left(\frac{1}{m_q} \right)^{2n}$$

Motivation: why precise m_Q ?

Snapshot from PDG

| <i>VALUE (GeV)</i> | <i>DOCUMENT ID</i> | <i>TECN</i> |
|--|-----------------------|-------------|
| 1.27 ± 0.02 | OUR EVALUATION | |
| 1.266 ±0.006 | 1 NARISON 2020 | THEO |
| 1.290 ^{+0.077} _{-0.053} | 2 ABRAMOWICZ 2018 | HERA |
| 1.273 ±0.010 | 3 BAZAVOV 2018 | LATT |
| 1.2737 ±0.0077 | 4 LYTLE 2018 | LATT |
| 1.223 ±0.033 | 5 PESET 2018 | THEO |
| 1.279 ±0.008 | 6 CHETYRKIN 2017 | THEO |
| 1.272 ±0.008 | 7 ERLER 2017 | THEO |
| 1.246 ±0.023 | 8 KIYO 2016 | THEO |
| 1.288 ±0.020 | 9 DEHNADI 2015 | THEO |
| 1.348 ±0.046 | 10 CARRASCO 2014 | LATT |
| 1.24 ±0.03 ^{+0.03} _{-0.07} | 11 ALEKHIN 2013 | THEO |
| 1.159 ±0.075 | 12 SAMOYLOV 2013 | NOMD |
| 1.278 ±0.009 | 13 BODENSTEIN 2011 | THEO |
| 1.28 ^{+0.07} _{-0.06} | 14 LASCHKA 2011 | THEO |
| 1.196 ±0.059 ±0.050 | 15 AUBERT 2010A | BABR |
| 1.25 ±0.04 | 16 SIGNER 2009 | THEO |

Motivation: why precise m_Q ?

Snapshot from PDG

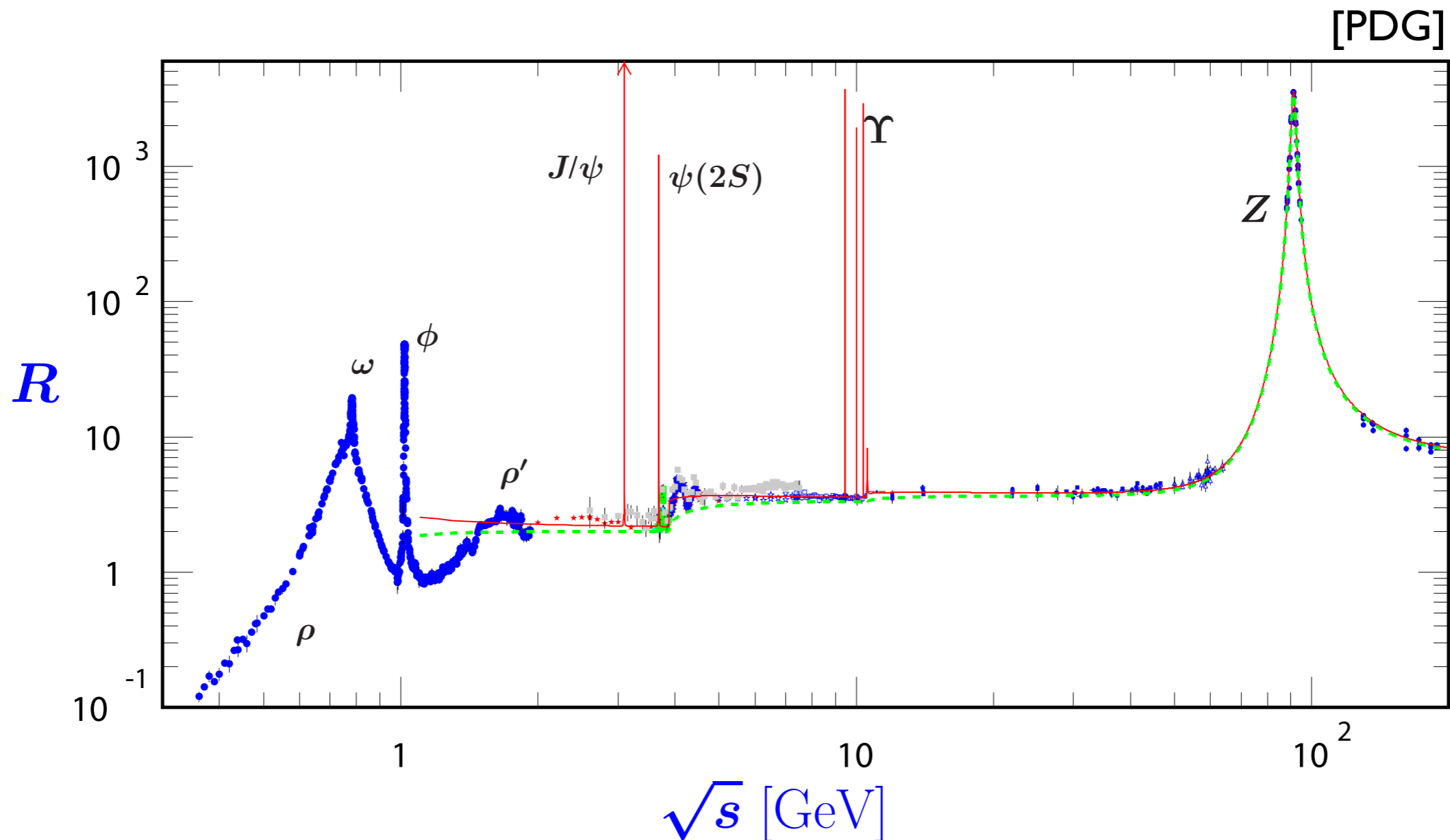
| VALUE (GeV) | DOCUMENT ID | TECN |
|--|---|------|
| $4.18^{+0.03}_{-0.02}$ | OUR EVALUATION of \overline{MS} Mass. | |
| 4.197 ± 0.008 | 1 NARISON 2020 | THEO |
| 4.049 $^{+0.138}_{-0.118}$ | 2 ABRAMOWICZ 2018 | HERA |
| 4.195 ± 0.014 | 3 BAZAVOV 2018 | LATT |
| 4.186 ± 0.037 | 4 PESET 2018 | THEO |
| 4.197 ± 0.022 | 5 KIYO 2016 | THEO |
| 4.183 ± 0.037 | 6 ALBERTI 2015 | THEO |
| 4.203 $^{+0.016}_{-0.034}$ | 7 BENEKE 2015 | THEO |
| 4.196 ± 0.023 | 8 COLQUHOUN 2015 | LATT |
| 4.176 ± 0.023 | 9 DEHNADI 2015 | THEO |
| 4.21 ± 0.11 | 10 BERNARDONI 2014 | LATT |
| 4.169 $\pm 0.002 \pm 0.008$ | 11 PENIN 2014 | THEO |
| 4.166 ± 0.043 | 12 LEE 2013O | LATT |
| 4.247 ± 0.034 | 13 LUCHA 2013 | THEO |
| 4.171 ± 0.009 | 14 BODENSTEIN 2012 | THEO |
| 4.29 ± 0.14 | 15 DIMOPOULOS 2012 | LATT |
| 4.18 $^{+0.05}_{-0.04}$ | 16 LASCHKA 2011 | THEO |
| 4.186 $\pm 0.044 \pm 0.015$ | 17 AUBERT 2010A | BABR |
| 4.163 ± 0.016 | 18 CHETYRKIN 2009 | THEO |
| 4.243 ± 0.049 | 19 SCHWANDA 2008 | BELL |

QCD Sum Rules

QCD Sum Rules

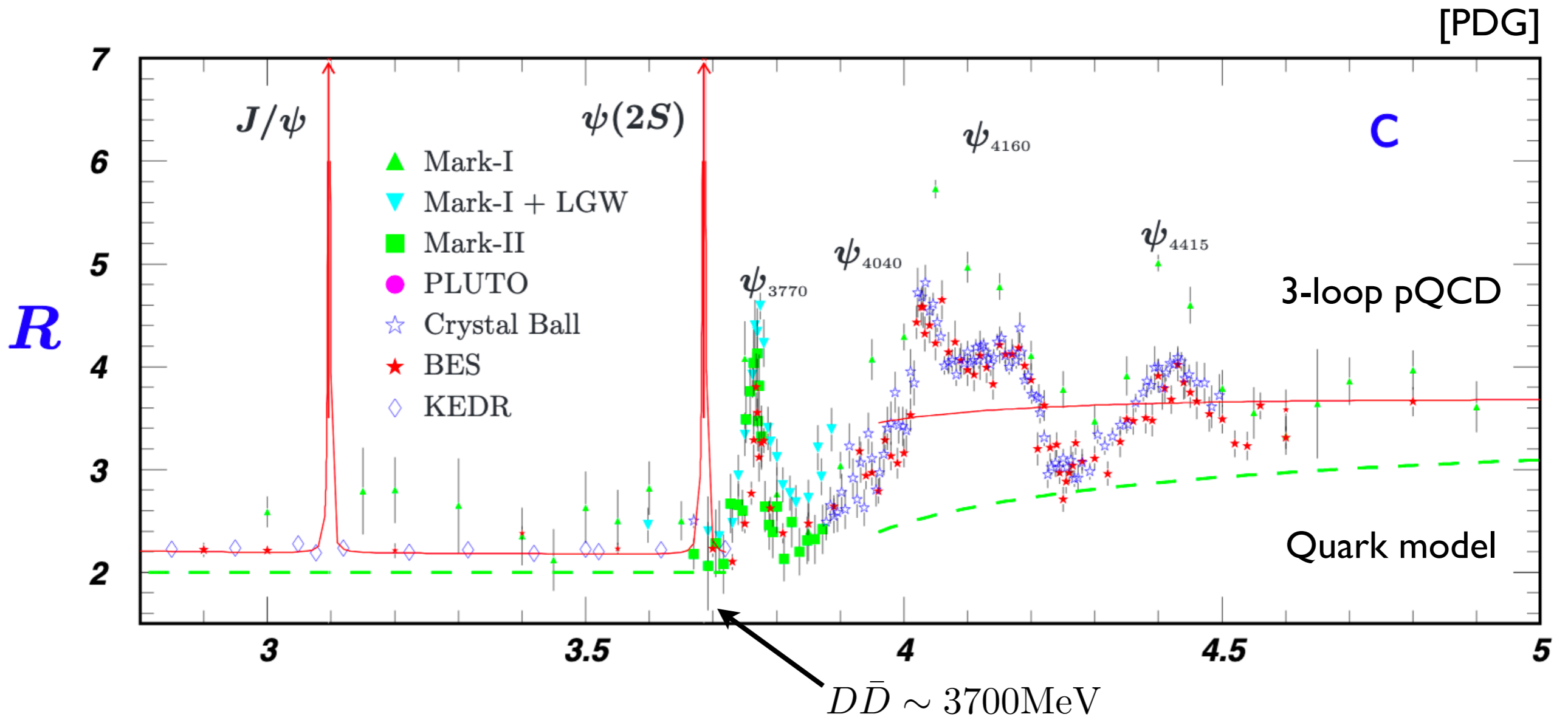
$$R(s) = \frac{\sigma(e^+e^- \rightarrow \text{hadrons})}{\sigma(e^+e^- \rightarrow \mu^+\mu^-)}$$

$$\sigma(e^+e^- \rightarrow \mu^+\mu^-) = 4\pi\alpha_{\text{em}}(s)^2/3s$$



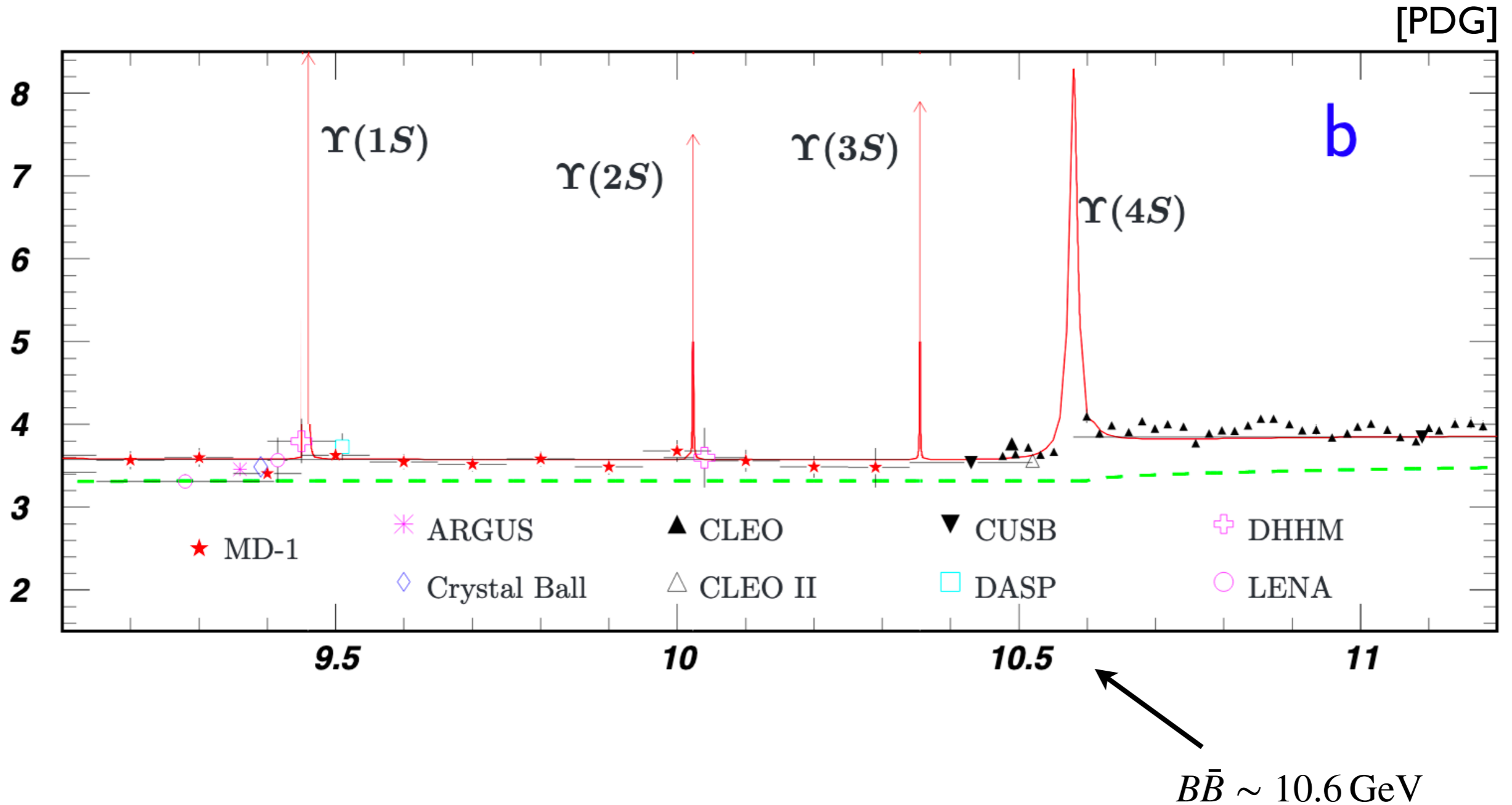
QCD Sum Rules

Zoom into the open-charm threshold



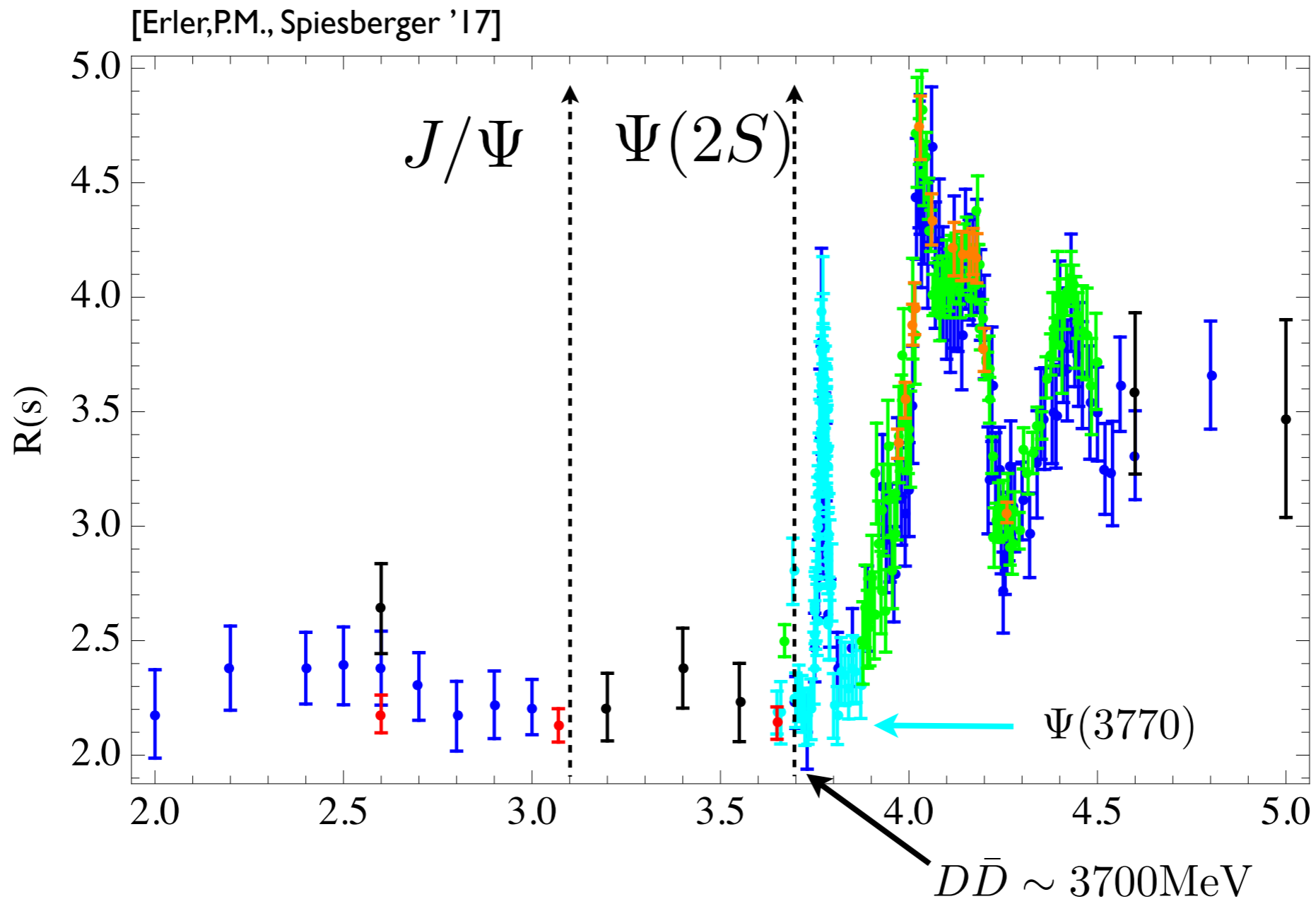
QCD Sum Rules

Zoom into the open-bottom threshold



QCD Sum Rules

Zoom into the open-charm threshold

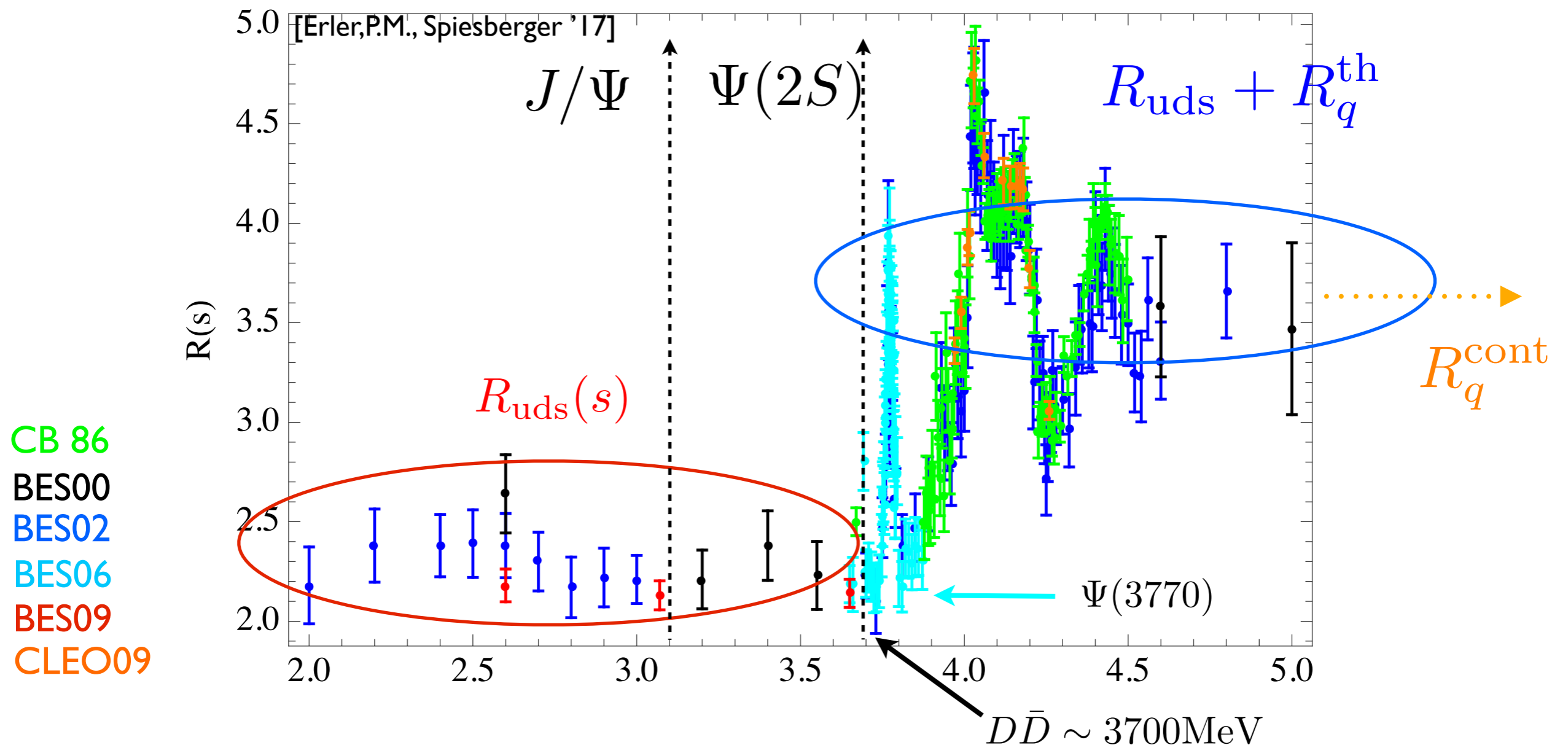


QCD Sum Rules

Zoom into the open-charm threshold

$$R(s) = R_{\text{uds}}(s) + R_q(s)$$

$$R_q(s) = R_q^{\text{Res}}(s) + R_q^{\text{th}}(s) + R_q^{\text{cont}}(s)$$



QCD Sum Rules

Using the optical theorem:

[SVZ,'79]

$$R(s) = 12\pi \text{Im}[\Pi(s + i\epsilon)]$$

$\Pi_q(s)$ is the correlator of two heavy-quark vector currents which can be calculated in pQCD order by order in α_s and satisfies a Dispersion Relation:

$$12\pi^2 \frac{\hat{\Pi}_q(0) - \hat{\Pi}_q(-t)}{t} = \int_{4m_q^2}^{\infty} \frac{ds}{s} \frac{R_q(s)}{s+t} \quad \hat{\Pi}_q(s) \text{ in } \overline{MS}$$

For $t \rightarrow 0$

$$\mathcal{M}_n := \frac{12\pi^2}{n!} \left. \frac{d^n}{dt^n} \hat{\Pi}_q(t) \right|_{t=0} = \int_{4m_q^2}^{\infty} \frac{ds}{s^{n+1}} R_q(s)$$

QCD Sum Rules

$\hat{\Pi}_q(s)$ can be Taylor expanded:

$$\Pi_q(t) = Q_q^2 \frac{3}{16\pi^2} \sum_{n \geq 0} \bar{C}_n \left(\frac{t}{4\hat{m}_q^2} \right)^n$$

QCD Sum Rules

$\hat{\Pi}_q(s)$ can be Taylor expanded:

$$\Pi_q(t) = Q_q^2 \frac{3}{16\pi^2} \sum_{n \geq 0} \bar{C}_n \left(\frac{t}{4\hat{m}_q^2} \right)^n$$

$$\mathcal{M}_n^{\text{pQCD}} = \frac{9}{4} Q_q^2 \left(\frac{1}{2\hat{m}_q(\hat{m}_q)} \right)^{2n} \bar{C}_n$$

$$\bar{C}_n = \bar{C}_n^{(0)} + \left(\frac{\hat{\alpha}}{\pi} \right) \bar{C}_n^{(1)} + \left(\frac{\hat{\alpha}}{\pi} \right)^2 \bar{C}_n^{(2)} + \left(\frac{\hat{\alpha}}{\pi} \right)^3 \bar{C}_n^{(3)} + \mathcal{O} \left(\frac{\hat{\alpha}}{\pi} \right)^4$$

[Maier et al, '08]

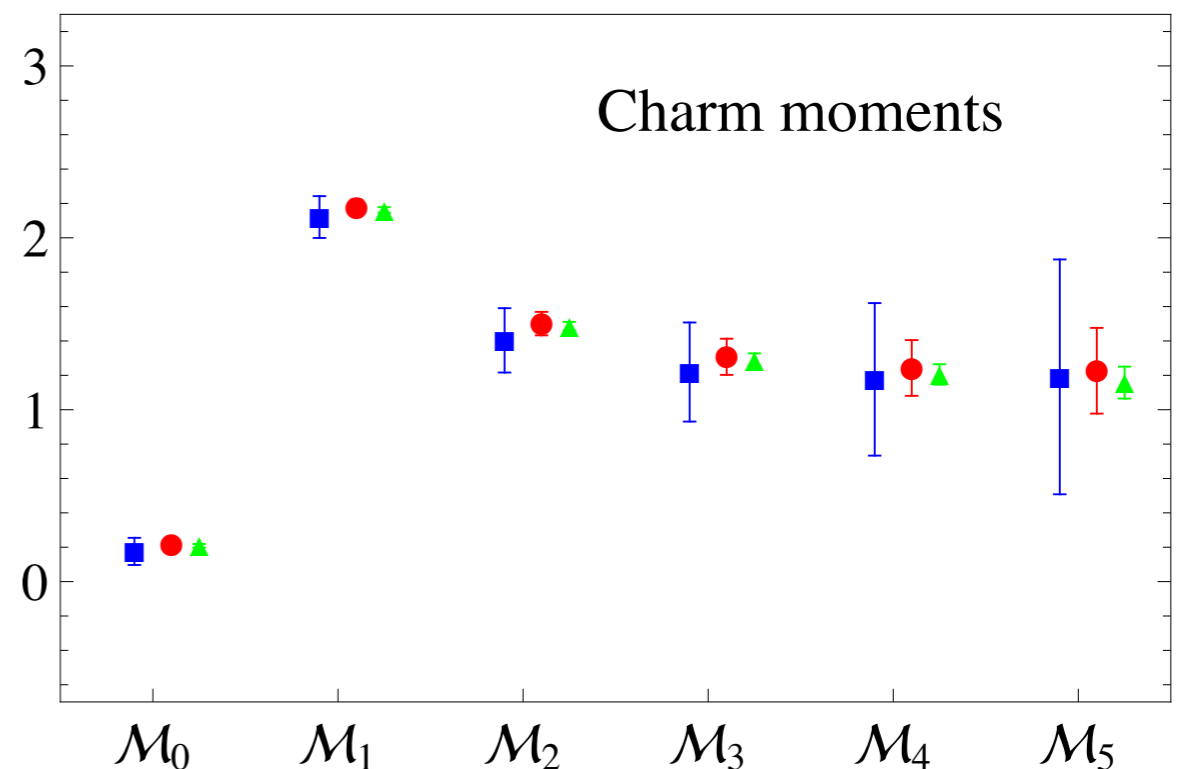
[Chetyrkin, Steinhauser'06]

[Melnikov, Ritberger'03]

[Kiyo et al '09]

[Hoang et al '09]

[Greynat et al '09]



$$\hat{\alpha} = \hat{\alpha}(\overline{m}_q)$$

QCD Sum Rules

Sum Rules:

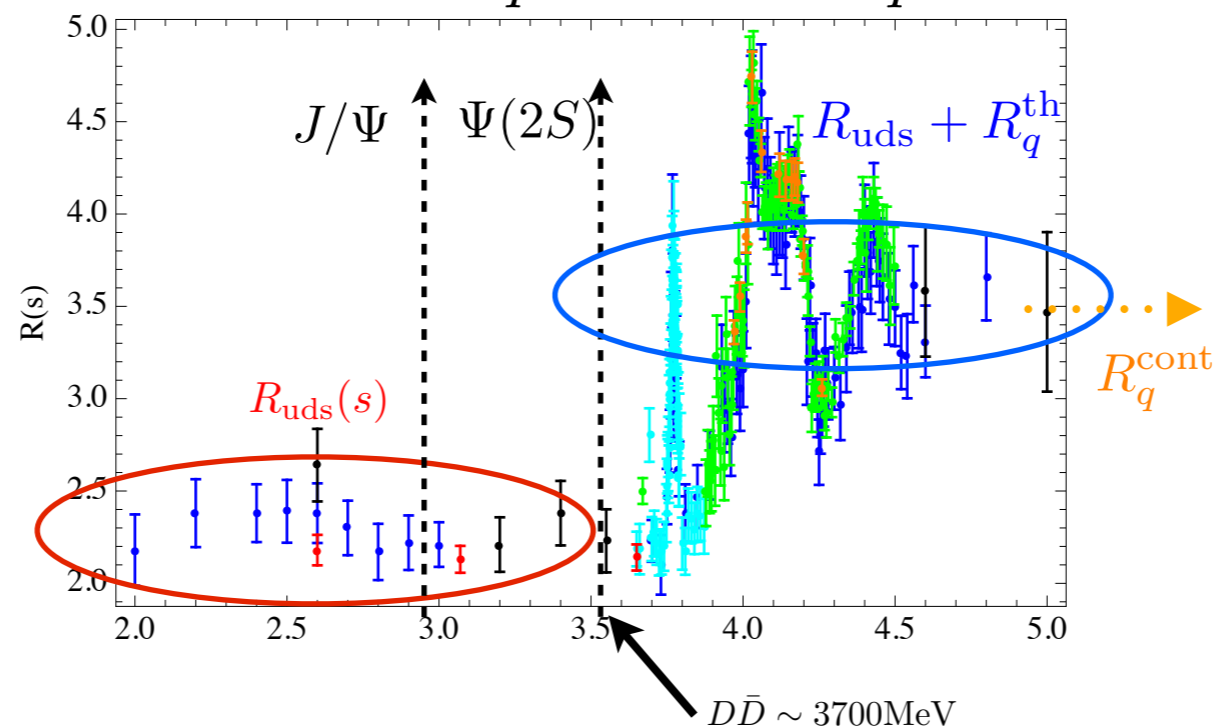
$$\mathcal{M}_n = \int_{4m_q^2}^{\infty} \frac{ds}{s^{n+1}} R_q(s)$$

L.h.s. from theory

$$\mathcal{M}_n^{\text{pQCD}} = \frac{9}{4} Q_q^2 \left(\frac{1}{2\hat{m}_q(\hat{m}_q)} \right)^{2n} \bar{C}_n$$

R.h.s. from experiment

$$R_q(s) = R_q^{\text{Res}}(s) + R_q^{\text{th}}(s) + R_q^{\text{cont}}(s)$$



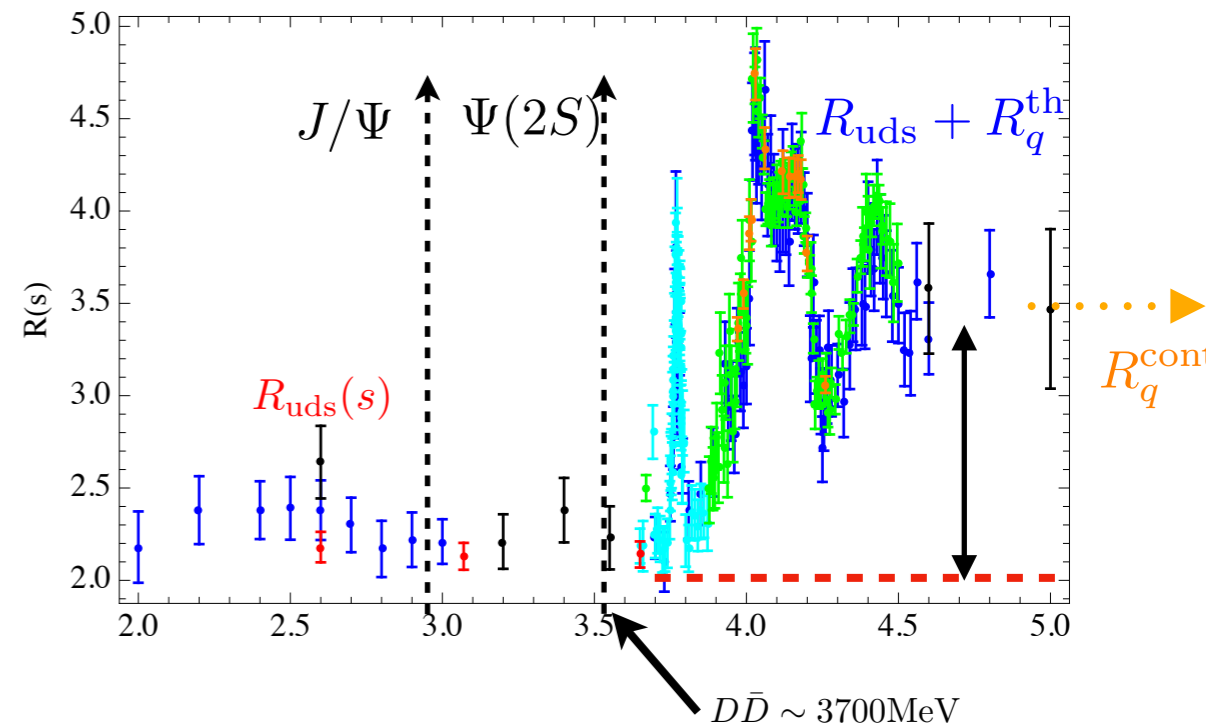
QCD Sum Rules

$$R_q(s) = R_q^{\text{Res}}(s) + R_q^{\text{th}}(s) + R_q^{\text{cont}}(s)$$

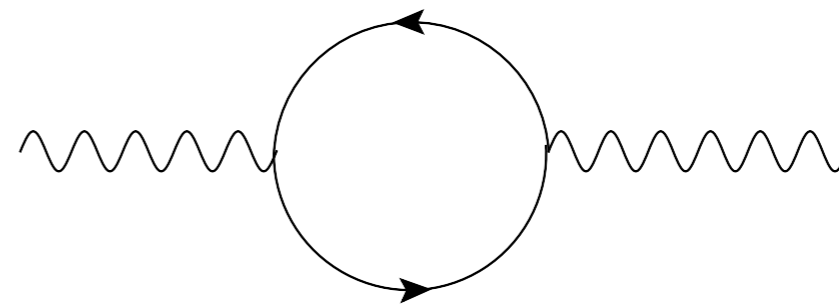
$$R_q^{\text{Res}}(s) = \frac{9\pi M_R \Gamma_R^e}{\alpha_{\text{em}}^2(M_R)} \delta(s - M_R^2)$$

$$R_q^{\text{th}}(s) = R_q(s) - R_{\text{background}} \quad (2M_D \leq \sqrt{s} \leq 4.8\text{GeV})$$

$$R_q^{\text{cont}}(s) \quad \text{calculated using pQCD} \quad (\sqrt{s} \geq 4.8\text{GeV})$$



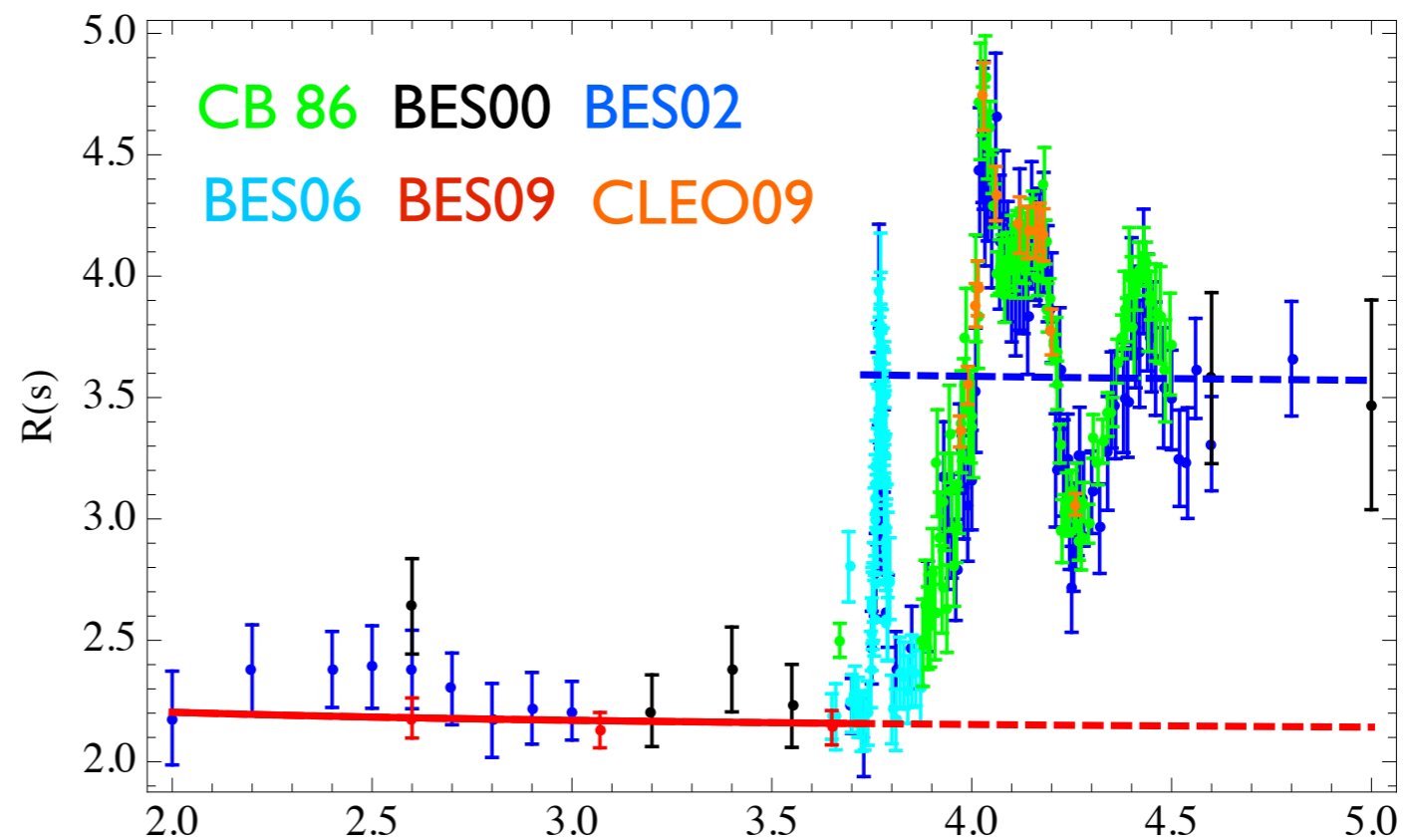
$$(2M_D \leq \sqrt{s} \leq 4.8\text{GeV})$$



Background

$$R_{\text{background}} = R_{\text{uds}} + R_{\text{uds}(\text{cb})} + R_{\text{sing}} + R_{\text{QED}}$$

Light flavor
contribution in
charm region

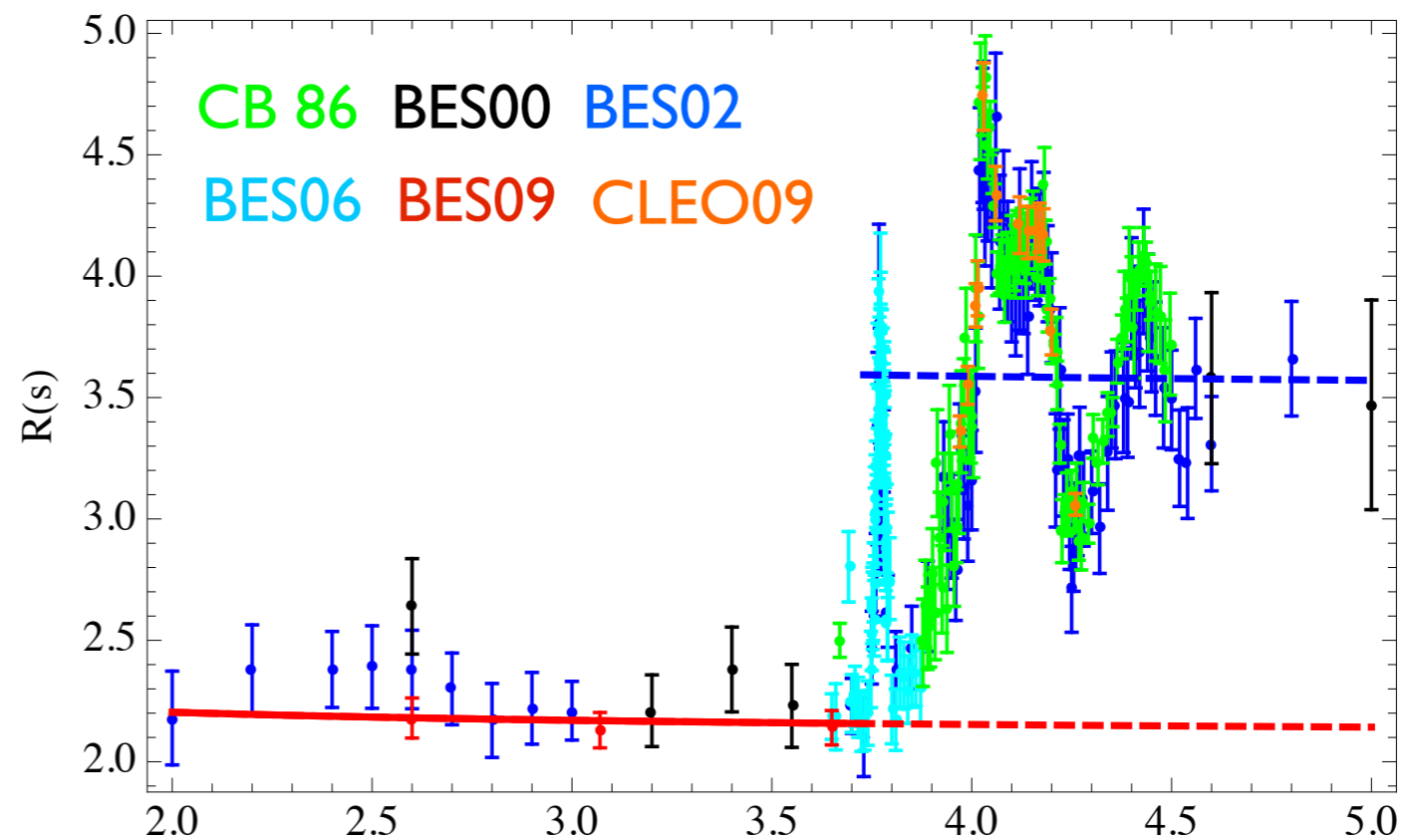
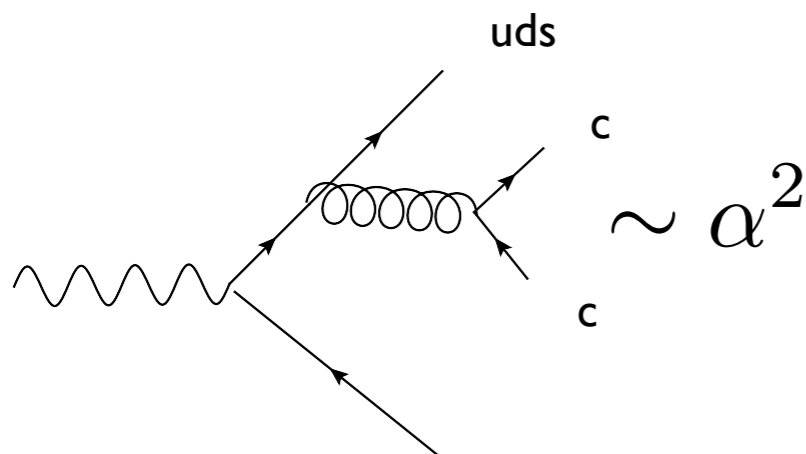


Using pQCD below threshold, calculate R, and extrapolate

Background

$$R_{\text{background}} = R_{\text{uds}} + R_{\text{uds}(cb)} + R_{\text{sing}} + R_{\text{QED}}$$

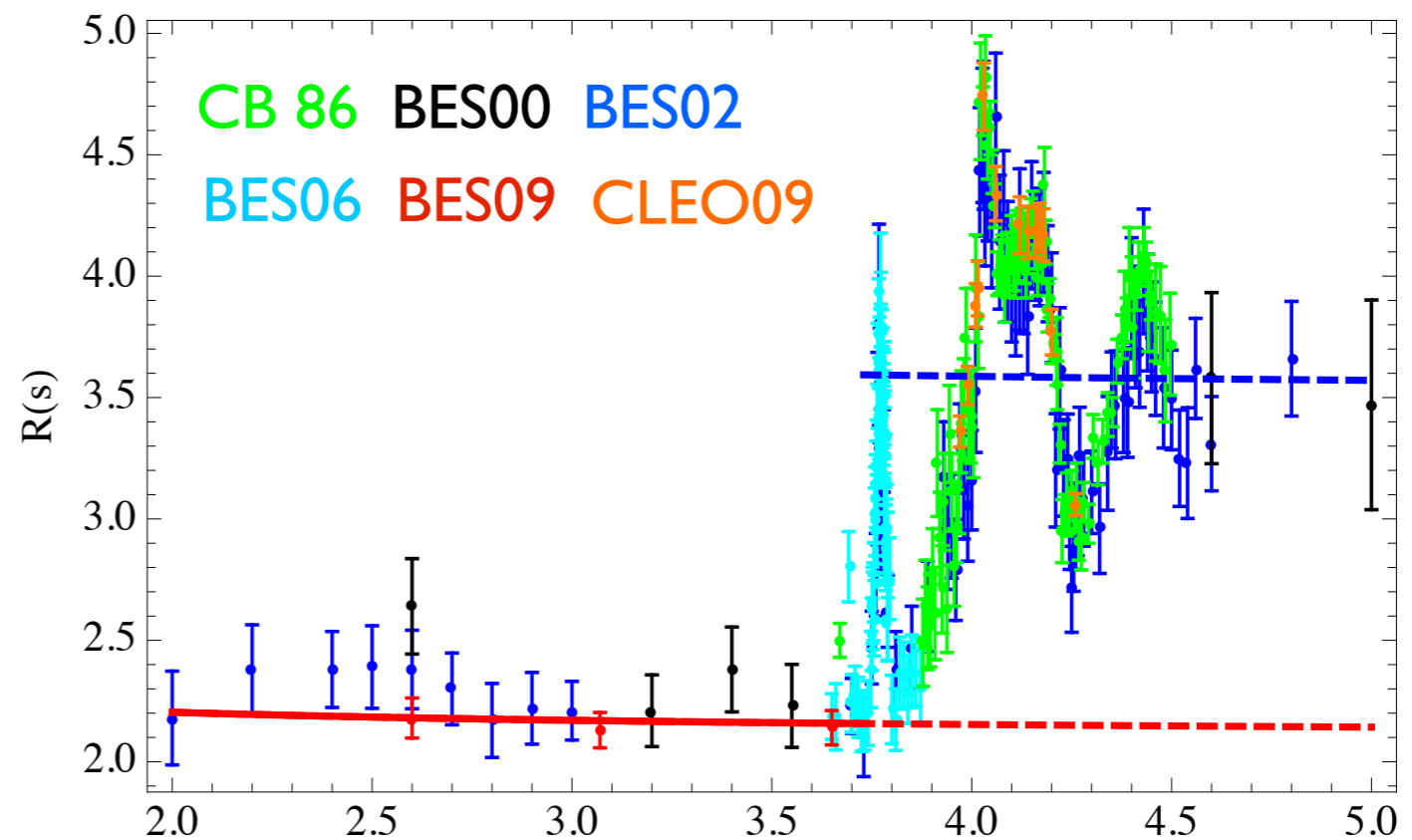
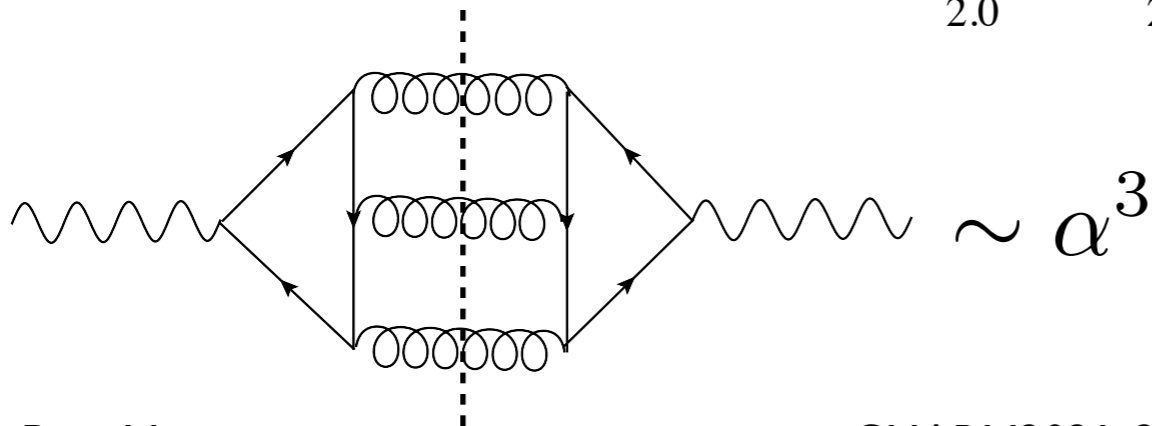
Light flavor
contribution in
charm region
+
secondary
production



Background

$$R_{\text{background}} = R_{\text{uds}} + R_{\text{uds}(\text{cb})} + R_{\text{sing}} + R_{\text{QED}}$$

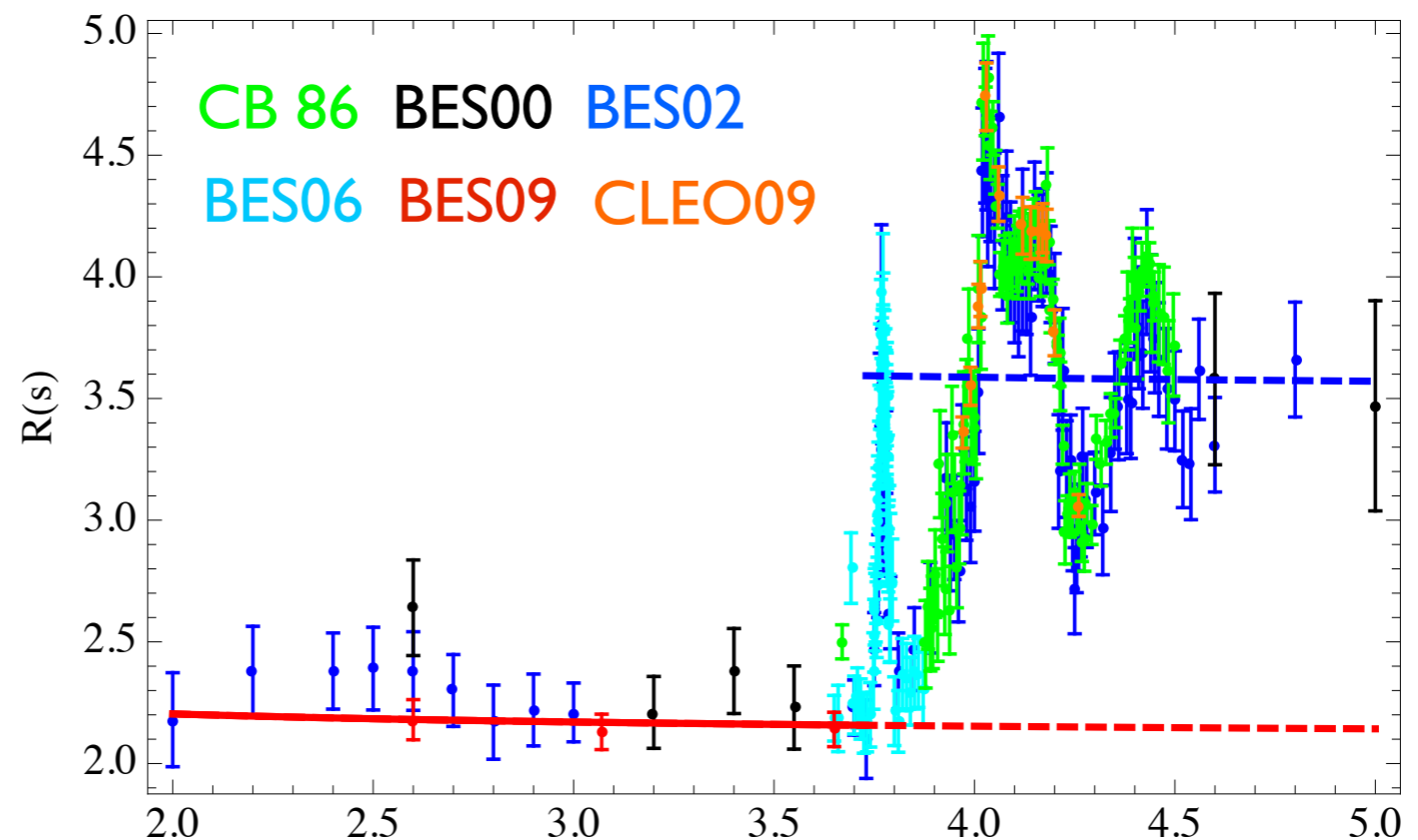
Light flavor
contribution in
charm region
+
secondary
production
+
singlet contribution



Background

$$R_{\text{background}} = R_{\text{uds}} + R_{\text{uds}(\text{cb})} + R_{\text{sing}} + R_{\text{QED}}$$

Light flavor
contribution in
charm region
+
secondary
production
+
singlet contribution
+
2loop QED



Non-perturbative effects

Non-perturbative effects due to gluon condensates to the moments are:

[Chetyrkin et al '12]

$$\mathcal{M}_n^{\text{nonp}}(\mu^2) = \frac{12\pi^2 Q_q^2}{(4\hat{m}_q^2)^{n+2}} \text{Cond} a_n \left(1 + \frac{\alpha_s(\hat{m}_q^2)}{\pi} b_n \right)$$

a_n, b_n are numbers, and $\text{Cond} = \langle \frac{\alpha_s}{\pi} G^2 \rangle = (5 \pm 5) \cdot 10^{-3} \text{GeV}^4$ [Dominguez et al '14]

↙ from fits to tau data

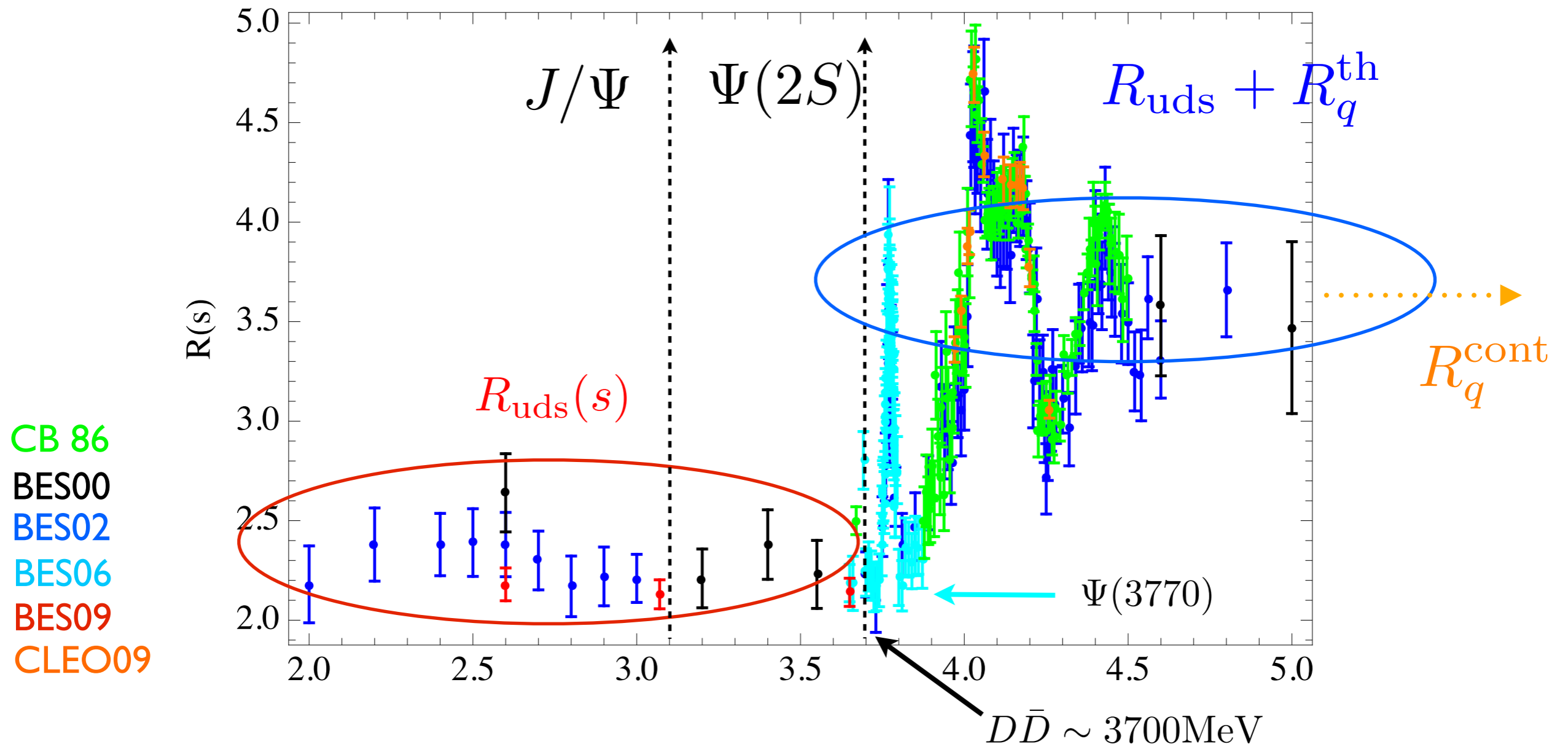
$$\frac{\mathcal{M}_n^{\text{nonp}}(\hat{m}_c)}{\mathcal{M}_n^{\text{th}}} \sim 0.5\% - 2\% \longrightarrow \Delta\hat{m}_c(\hat{m}_c) \sim 2\text{MeV} - 8\text{MeV}$$

QCD Sum Rules

$$R(s) = R_{\text{uds}}(s) + R_q(s)$$

$$R_q(s) = R_q^{\text{Res}}(s) + R_q^{\text{th}}(s) + R_q^{\text{cont}}(s)$$

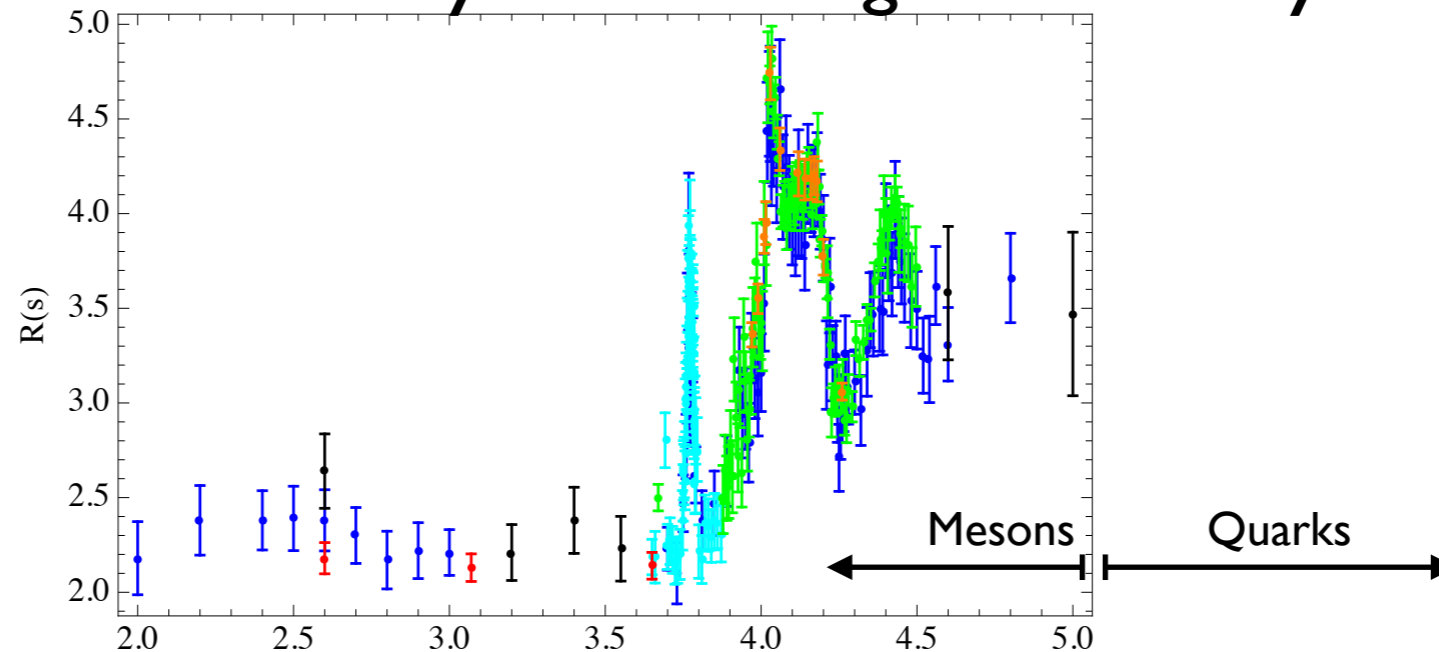
[Erlar, P.M., Spiesberger '17]



QCD Sum Rules

Our approach is *different*

- We try to avoid *local* duality: consider *global* duality



Standard procedure:
$$\mathcal{M}_n^{\text{pQCD}} = R_q^{\text{Res}} + \int_{D\bar{D}}^{s_0} \frac{ds}{s^{n+1}} R_q(s) + p\text{QCD}(\mu)$$

We really want:
$$\mathcal{M}_n^{\text{pQCD}} = R_q^{\text{Res}} + \int_{D\bar{D}}^{\infty} \frac{ds}{s^{n+1}} R_q(s)$$

QCD Sum Rules

Our approach is *different*

- We try to avoid *local* duality: consider *global* duality
- Then, we do *not* use *experimental data* on threshold region, only resonances below threshold
 - Experimental data in threshold used for error estimation
- How you do it then? Use *two different moment* equations to determine the continuum requiring self-consistency:

- extract the quark mass

$$\mathcal{M}_n^{\text{pQCD}} = R_q^{\text{Res}} + \int_{D\bar{D}}^{\infty} \frac{ds}{s^{n+1}} R_q^{\text{cont}}(s)$$

Charm

QCD Sum Rules

Our approach

For a global duality:

$\hat{\Pi}_q(s)$ in \overline{MS}

$$12\pi^2 \frac{\hat{\Pi}_q(0) - \hat{\Pi}_q(-t)}{t} = \int_{4m_q^2}^{\infty} \frac{ds}{s} \frac{R_q(s)}{s+t}$$

$t \rightarrow \infty$ define the \mathcal{M}_0

[Erler, Luo '03]

QCD Sum Rules

Our approach

For a global duality:

$\hat{\Pi}_q(s)$ in \overline{MS}

$$12\pi^2 \frac{\hat{\Pi}_q(0) - \hat{\Pi}_q(-t)}{t} = \int_{4m_q^2}^{\infty} \frac{ds}{s} \frac{R_q(s)}{s+t}$$

$t \rightarrow \infty$ define the \mathcal{M}_0 (but has a divergent part)

[Erler, Luo '03]

$$\lim_{t \rightarrow \infty} \hat{\Pi}_q(-t) \sim \log(t) \longleftrightarrow \int_{4m_q^2}^{\infty} \frac{ds}{s} R_q(s) \sim \log(\infty)$$

Fortunately, divergence given by the zero-mass limit of $R(s)$, can be easily subtracted

[Chetyrkin, Harlander, Kühn, '00]

QCD Sum Rules

zero-mass limit of R(s)

Our approach

$$\begin{aligned} \lambda_1^q(s) = & 1 + \frac{\alpha_s(s)}{\pi} && \text{[Chetyrkin, Harlander, Kühn, '00]} \\ & + \left[\frac{\alpha_s(s)}{\pi} \right]^2 \left[\frac{365}{24} - 11\zeta(3) + n_q \left(\frac{2}{3}\zeta(3) - \frac{11}{12} \right) \right] \\ & + \left[\frac{\alpha_s(s)}{\pi} \right]^3 \left[\frac{87029}{288} - \frac{121}{8}\zeta(2) - \frac{1103}{4}\zeta(3) + \frac{275}{6}\zeta(5) \right. \\ & \quad \left. + n_q \left(-\frac{7847}{216} + \frac{11}{6}\zeta(2) + \frac{262}{9}\zeta(3) - \frac{25}{9}\zeta(5) \right) \right. \\ & \quad \left. + n_q^2 \left(\frac{151}{162} - \frac{1}{18}\zeta(2) - \frac{19}{27}\zeta(3) \right) \right] \\ & \quad \quad \quad n_q \text{ active flavors} \end{aligned}$$

QCD Sum Rules

Our approach

Zeroth Sum Rule:

$$\begin{aligned}
 & \sum_{\text{resonances}} \frac{9\pi\Gamma_R^e}{3Q_q^2 M_R \hat{\alpha}_{em}^2(M_R)} + \int_{4M^2}^{\infty} \frac{ds}{s} \frac{R_q^{\text{cont}}}{3Q_q^2} - \int_{\hat{m}_q^2}^{\infty} \frac{ds}{s} \lambda_1^q(s) \\
 &= -\frac{5}{3} + \frac{\hat{\alpha}_s}{\pi} \left[4\zeta(3) - \frac{7}{2} \right] \qquad \hat{\alpha}_s = \alpha_s(\hat{m}_q^2) \\
 & \quad + \left(\frac{\hat{\alpha}_s}{\pi} \right)^2 \left[\frac{2429}{48} \zeta(3) - \frac{25}{3} \zeta(5) - \frac{2543}{48} + n_q \left(\frac{677}{216} - \frac{19}{9} \zeta(3) \right) \right] \\
 & \quad + \left(\frac{\hat{\alpha}_s}{\pi} \right)^3 \left[-9.86 + 0.40 n_q - 0.01 n_q^2 \right] \\
 &= -1.667 + 1.308 \frac{\hat{\alpha}_s}{\pi} + 1.595 \left(\frac{\hat{\alpha}_s}{\pi} \right)^2 - 8.427 \left(\frac{\hat{\alpha}_s}{\pi} \right)^3
 \end{aligned}$$

QCD Sum Rules

Our approach

Zeroth Sum Rule:

$$\sum_{\text{resonances}} \frac{9\pi\Gamma_R^e}{3Q_q^2 M_R \hat{\alpha}_{em}^2(M_R)} + \int_{4M^2}^{\infty} \frac{ds}{s} \frac{R_q^{\text{cont}}}{3Q_q^2} - \int_{\hat{m}_q^2}^{\infty} \frac{ds}{s} \lambda_1^q(s)$$

$\hat{\alpha}_s = \alpha_s(\hat{m}_q^2)$

[PDG]

| R | M_R [GeV] | Γ_R^e [keV] |
|------------|-------------|--------------------|
| J/Ψ | 3.096916 | 5.55(14) |
| $\Psi(2S)$ | 3.686109 | 2.36(4) |

QCD Sum Rules

Our approach

Zeroth Sum Rule:

$$\sum_{\text{resonances}} \frac{9\pi\Gamma_R^e}{3Q_q^2 M_R \hat{\alpha}_{em}^2(M_R)} + \int_{4M^2}^{\infty} \frac{ds}{s} \frac{R_q^{\text{cont}}}{3Q_q^2} - \int_{\hat{m}_q^2}^{\infty} \frac{ds}{s} \lambda_1^q(s)$$

$$\hat{\alpha}_s = \alpha_s(\hat{m}_q^2)$$

$$\hat{\alpha}_{em}(0) \sim 0.98\hat{\alpha}_{em}(M_{J/\Psi})$$

$$\Delta\hat{\alpha}_{em} \rightarrow \Delta m_c \sim 12\text{MeV}$$

QCD Sum Rules

Our approach: ansatz

Zeroth Sum Rule: invoke global quark-hadron duality

[Erler, Luo '03]

$$R_q^{\text{cont}}(s) = 3Q_q^2 \lambda_1^q(s) \sqrt{1 - \frac{4\hat{m}_q^2(2M)}{s'}} \left[1 + \lambda_3^q \frac{2\hat{m}_q^2(2M)}{s'} \right]$$

Simpler version of analytic reconstruction [Greynat, PM, Peris'12]

$$s' = s + 4(\hat{m}_q^2(2M) - M^2)$$

Two parameters to determine: m_q , λ_3^q

QCD Sum Rules

Our approach: ansatz

Zeroth Sum Rule: invoke global quark-hadron duality

[Erler, Luo '03]

$$R_q^{\text{cont}}(s) = 3Q_q^2 \lambda_1^q(s) \sqrt{1 - \frac{4\hat{m}_q^2(2M)}{s'}} \left[1 + \lambda_3^q \frac{2\hat{m}_q^2(2M)}{s'} \right]$$

Simpler version of analytic reconstruction [Greynat, PM, Peris'12]

$$s' = s + 4(\hat{m}_q^2(2M) - M^2)$$

Two parameters to determine: m_q , λ_3^q

We need two equations: **zeroth moment** + **nth moment**

$$\frac{9}{4} Q_q^2 \left(\frac{1}{2\hat{m}_q(\hat{m}_q)} \right)^{2n} \bar{C}_n = \sum_{\text{resonances}} \frac{9\pi\Gamma_R^e}{M_R^{2n+1} \hat{\alpha}_{em}^2(M_R)} + \int_{4M^2}^{\infty} \frac{ds}{s^{n+1}} R_q(s)$$

$n \geq 1$

QCD Sum Rules

Our approach: ansatz

Zeroth Sum Rule: invoke global quark-hadron duality

[Erler, Luo '03]

$$R_q^{\text{cont}}(s) = 3Q_q^2 \lambda_1^q(s) \sqrt{1 - \frac{4\hat{m}_q^2(2M)}{s'}} \left[1 + \lambda_3^q \frac{2\hat{m}_q^2(2M)}{s'} \right]$$

Simpler version of analytic reconstruction [Greynat, PM, Peris'12]

$$s' = s + 4(\hat{m}_q^2(2M) - M^2)$$

Two parameters to determine: m_q , λ_3^q

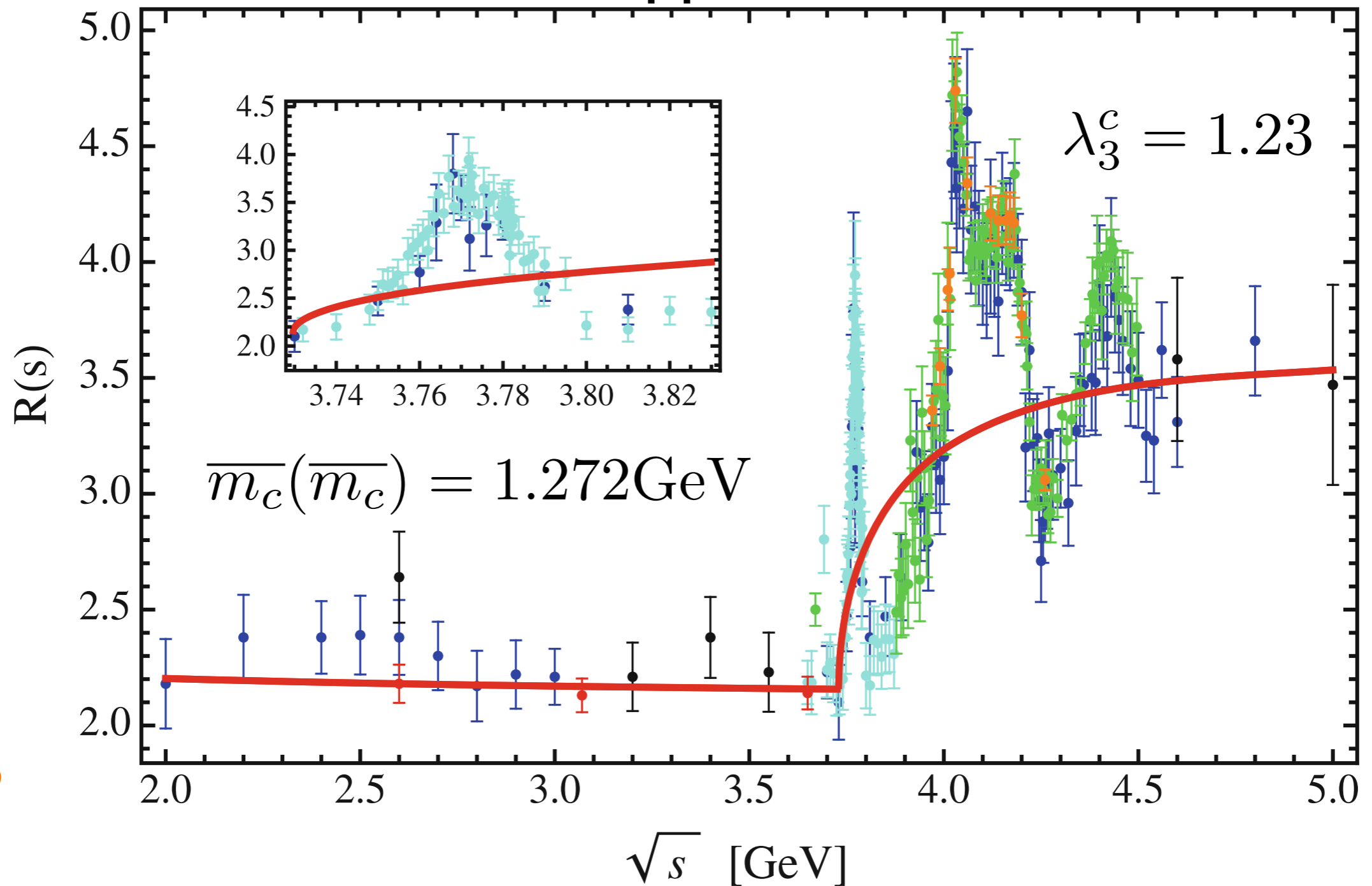
We use **Zeroth** + **2nd** moments
(no experimental data on R(s) so far)

we require self-consistency among the 2 moments

| n | Resonances | Continuum | Total | Theory |
|-----|------------|--------------------|------------|-------------------|
| 0 | 1.231 (24) | -3.229(+28)(43)(1) | -1.999(56) | Input (11) |
| 1 | 1.184 (24) | 0.966(+11)(17)(4) | 2.150(33) | 2.169(16) |
| 2 | 1.161 (25) | 0.336(+5)(8)(9) | 1.497(28) | Input (25) |
| 3 | 1.157 (26) | 0.165(+3)(4)(16) | 1.322(31) | 1.301(39) |
| 4 | 1.167 (27) | 0.103(+2)(2)(26) | 1.270(38) | 1.220(60) |
| 5 | 1.188 (28) | 0.080(+1)(1)(38) | 1.268(47) | 1.175(95) |

QCD Sum Rules

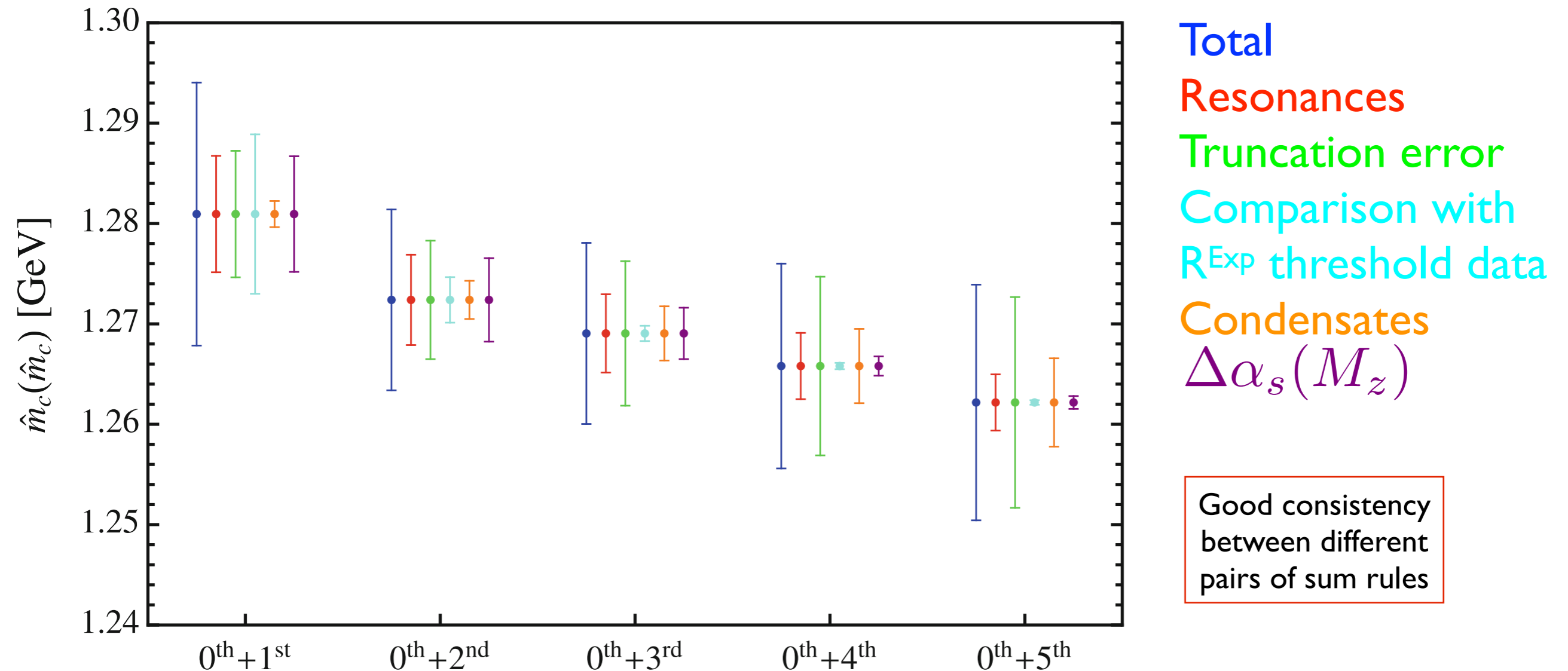
Our approach



QCD Sum Rules

Our approach

Repeat for each pair Zeroth+nth moment



QCD Sum Rules


Our approach: **error budget**

Resonances:

$$\frac{9}{4} Q_q^2 \left(\frac{1}{2\hat{m}_q(\hat{m}_q)} \right)^{2n} \bar{C}_n = \sum_{\text{resonances}} \frac{9\pi\Gamma_R^e}{M_R^{2n+1} \hat{\alpha}_{em}^2(M_R)} + \int_{4M^2}^{\infty} \frac{ds}{s^{n+1}} R_q(s)$$

from 6 MeV to 3 MeV
(0th+1st) (0th+5th)

(completely dominated by J/Ψ)



| R | M_R [GeV] | Γ_R^e [keV] |
|------------|-------------|--------------------|
| J/Ψ | 3.096916 | 5.55(14) |
| $\Psi(2S)$ | 3.686109 | 2.36(4) |

QCD Sum Rules

Our approach: **error budget**

Truncation Error (theory error):

$$\mathcal{M}_n^{\text{pQCD}} = \frac{9}{4} Q_q^2 \left(\frac{1}{2\hat{m}_q(\hat{m}_q)} \right)^{2n} \bar{C}_n$$

$$\bar{C}_n = \bar{C}_n^{(0)} + \left(\frac{\hat{\alpha}}{\pi} \right) \bar{C}_n^{(1)} + \left(\frac{\hat{\alpha}}{\pi} \right)^2 \bar{C}_n^{(2)} + \left(\frac{\hat{\alpha}}{\pi} \right)^3 \bar{C}_n^{(3)} + \mathcal{O} \left(\frac{\hat{\alpha}}{\pi} \right)^4$$

(use the largest group th. factor in the next uncalculated pert. order)

[Erler, Luo '03]

$$\Delta \mathcal{M}_n^{(4)} = \pm N_C C_F C_A^3 Q_q^2 \left[\frac{\hat{\alpha}_s(\hat{m}_q)}{\pi} \right]^4 \left(\frac{1}{2\hat{m}_q(\hat{m}_q)} \right)^{2n}$$

Example known orders

| n | $\frac{\Delta \mathcal{M}_n^{(2)}}{ \mathcal{M}_n^{(2)} }$ | $\frac{\Delta \mathcal{M}_n^{(3)}}{ \mathcal{M}_n^{(3)} }$ |
|-----|--|--|
| 0 | 1.88 | 3.03 |
| 1 | 2.14 | 2.84 |
| 2 | 1.92 | 4.58 |
| 3 | 3.25 | 5.63 |
| 4 | 6.70 | 4.30 |
| 5 | 19.18 | 3.62 |

from 5 MeV to 10 MeV
(0th+1st) (0th+5th)

More conservative than varying the renorm. scale within a factor of 4

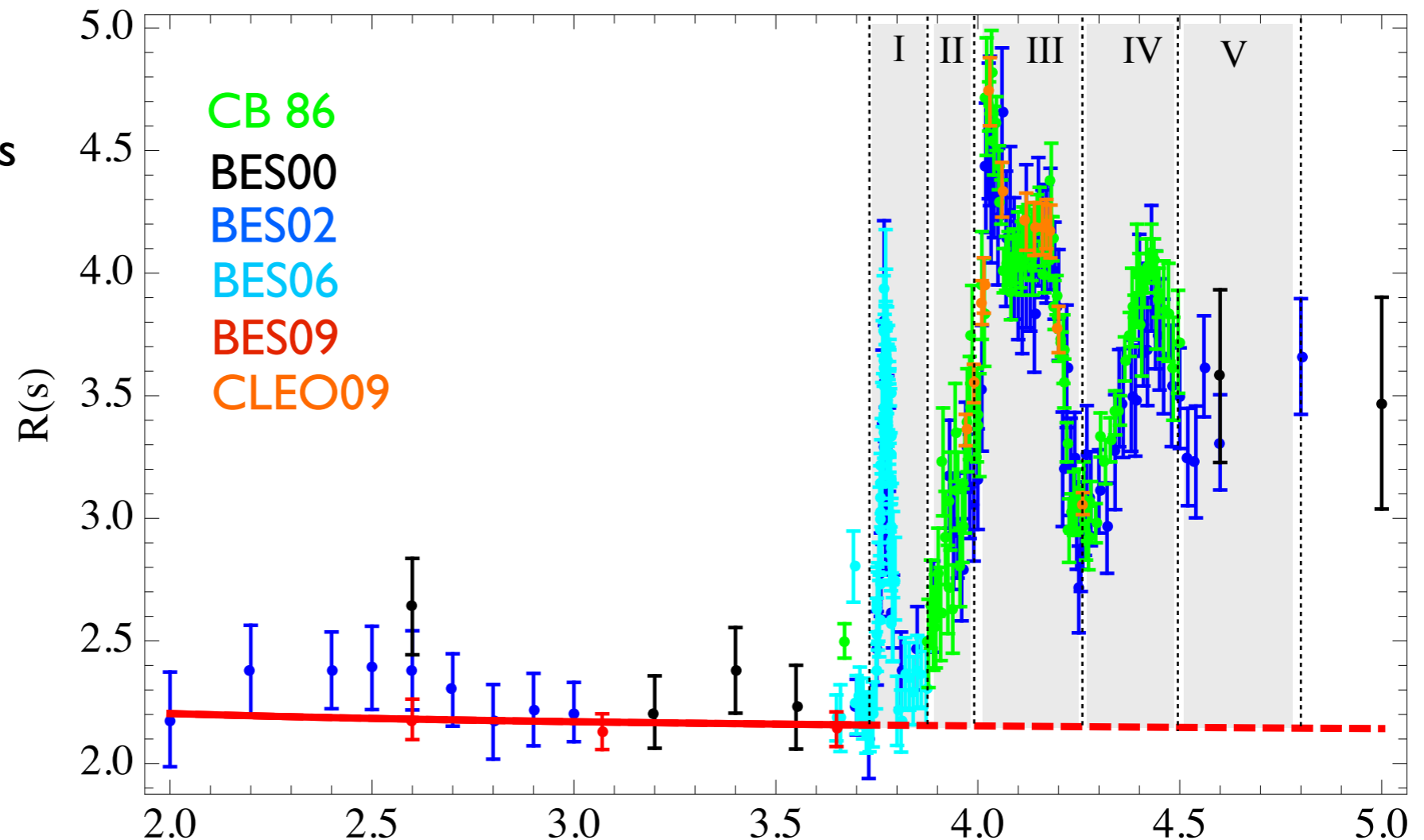
QCD Sum Rules

Our approach: **error budget**

Comparison with R^{Exp} threshold data:

$$(2M_D \leq \sqrt{s} \leq 4.8\text{GeV})$$

Calculate Exp moments



QCD Sum Rules

Our approach: **error budget**

Comparison with R^{Exp} threshold data:

| Collab. | n | $[2M_{D^0}, 3.872]$ | $[3.872, 3.97]$ | $[3.97, 4.26]$ | $[4.26, 4.496]$ | $[4.496, 4.8]$ |
|---------|-----|---------------------|-----------------|-----------------|-----------------|----------------|
| CB86 | 0 | – | 0.0339(22)(24) | 0.2456(25)(172) | 0.1543(27)(108) | – |
| | 1 | – | 0.0220(14)(15) | 0.1459(16)(102) | 0.0801(14)(56) | – |
| | 2 | – | 0.0143(9)(10) | 0.0868 (9)(61) | 0.0416(7)(29) | – |
| BES02 | 0 | 0.0334(24)(17) | 0.0362(29)(18) | 0.2362(41)(118) | 0.1399(38)(70) | 0.1705(63)(85) |
| | 1 | 0.0232(17)(12) | 0.0235(19)(12) | 0.1401(24)(70) | 0.0726(20)(36) | 0.0788(30)(39) |
| | 2 | 0.0161(12)(8) | 0.0152(13)(8) | 0.0832(15)(42) | 0.0378(10)(19) | 0.0365(14)(18) |
| BES06 | 0 | 0.0311(16)(15) | – | – | – | – |
| | 1 | 0.0217(11)(11) | – | – | – | – |
| | 2 | 0.0151(8)(7) | – | – | – | – |
| CLEO09 | 0 | – | – | 0.2591(22)(52) | – | – |
| | 1 | – | – | 0.1539(13)(31) | – | – |
| | 2 | – | – | 0.0915(8)(18) | – | – |
| Total | 0 | 0.0319(14)(11) | 0.0350(18)(15) | 0.2545(18)(46) | 0.1448(27)(59) | 0.1705(63)(85) |
| | 1 | 0.0222(9)(8) | 0.0227(12)(10) | 0.1511(11)(27) | 0.0752(14)(31) | 0.0788(30)(39) |
| | 2 | 0.0155(6)(6) | 0.0147(8)(6) | 0.0899(6)(16) | 0.0391(7)(16) | 0.0365(14)(18) |

QCD Sum Rules

Our approach: **error budget**

Comparison with R^{Exp} threshold data:

$$\int_{(2M_{D0})^2}^{(4.8 \text{ GeV})^2} \frac{ds}{s} R_c^{\text{cont}}(s) \Big|_{\hat{m}_c=1.272 \text{ GeV}} = \mathcal{M}_0^{\text{Data}} = 0.6367(195) \longrightarrow \lambda_3^{c,\text{exp}} = 1.34(17)$$

$(2M_D \leq \sqrt{s} \leq 4.8 \text{ GeV})$

Error induced to Quark mass:

I) $\lambda_3^c = 1.23 \rightarrow \lambda_3^{c,\text{exp}} = 1.34$

from + 6.4 MeV to + 0.2 MeV

II) $\Delta\lambda_3^{c,\text{exp}} = 0.17$

from 4.7 MeV to 0.1 MeV

| n | Data | $\lambda_3^c = 1.34(17)$ | $\lambda_3^c = 1.23$ |
|-----|-------------|--------------------------|----------------------|
| 0 | 0.6367(195) | 0.6367(195) | 0.6239 |
| 1 | 0.3500(102) | 0.3509(111) | 0.3436 |
| 2 | 0.1957(54) | 0.1970(65) | 0.1928 |
| 3 | 0.1111(29) | 0.1127(38) | 0.1102 |
| 4 | 0.0641(16) | 0.0657(23) | 0.0642 |
| 5 | 0.0375(9) | 0.0389(14) | 0.0380 |

QCD Sum Rules

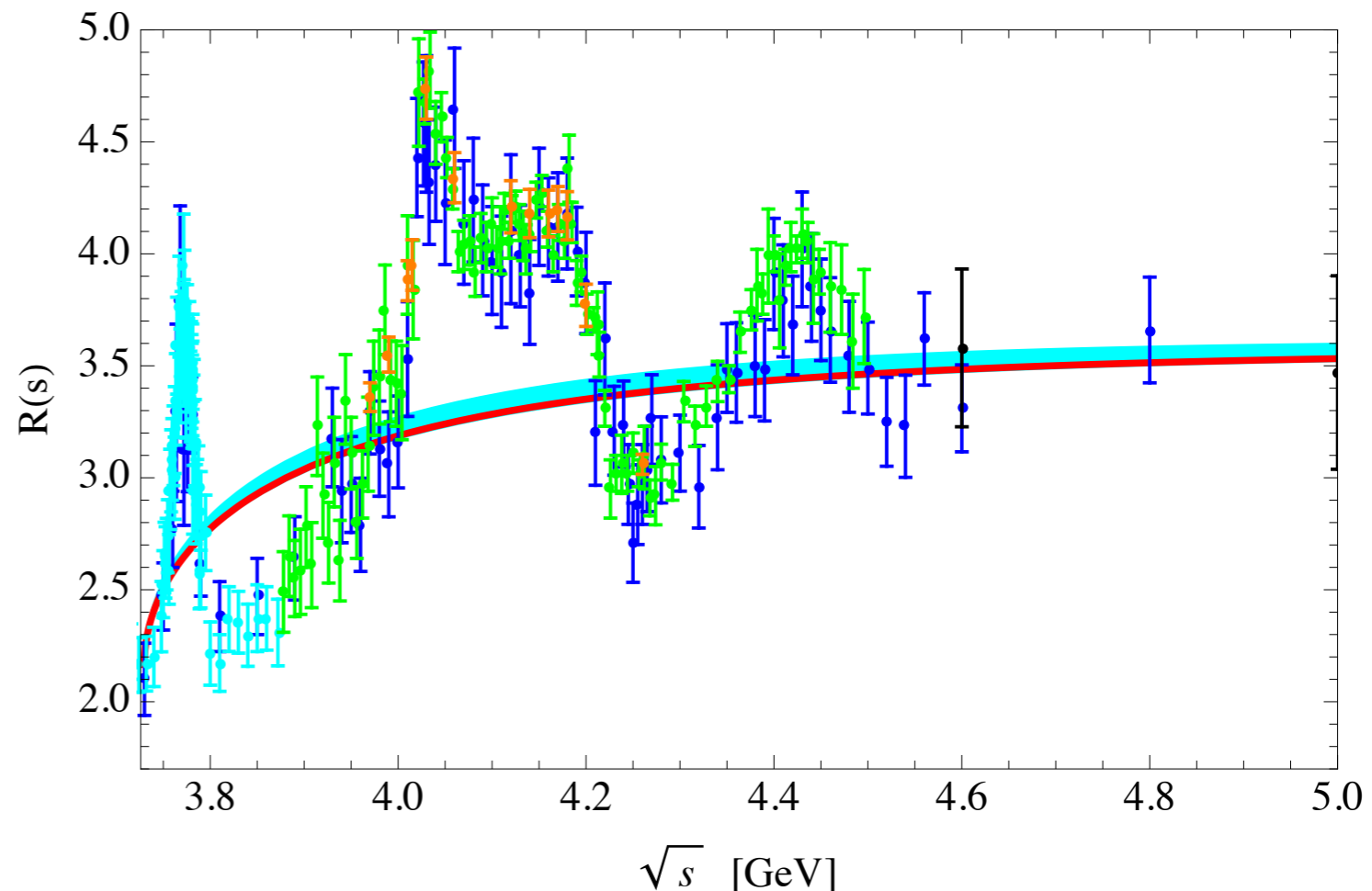
Our approach: **error budget**

Comparison with R^{Exp} threshold data:

$$\int_{(2M_{D^0})^2}^{(4.8 \text{ GeV})^2} \frac{ds}{s} R_c^{\text{cont}}(s) \Big|_{\hat{m}_c = 1.272 \text{ GeV}} = \mathcal{M}_0^{\text{Data}} = 0.6367(195) \longrightarrow \lambda_3^{c,\text{exp}} = 1.34(17)$$

Error induced to Quark mass:

- I) $\lambda_3^c = 1.23 \rightarrow \lambda_3^{c,\text{exp}} = 1.34$
from +6.4 MeV to +0.2 MeV
- II) $\Delta\lambda_3^{c,\text{exp}} = 0.17$
from 4.7 MeV to 0.1 MeV



QCD Sum Rules

Our approach: **error budget**

Condensates:

Non-perturbative effects due to gluon condensates to the moments are: [Chetyrkin et al '12]

$$\mathcal{M}_n^{\text{nonp}}(\mu^2) = \frac{12\pi^2 Q_q^2}{(4\hat{m}_q^2)^{n+2}} \text{Cond} a_n \left(1 + \frac{\alpha_s(\hat{m}_q^2)}{\pi} b_n \right)$$

a_n, b_n are numbers, and $\text{Cond} = \langle \frac{\alpha_s}{\pi} G^2 \rangle = (5 \pm 5) \cdot 10^{-3} \text{GeV}^4$ [Dominguez et al '14]

$$\Delta \langle \frac{\alpha_s}{\pi} G^2 \rangle = 5 \cdot 10^{-3} \text{GeV}^4 \quad \longrightarrow \quad \begin{array}{cc} \text{from 1 MeV to 4 MeV} \\ \text{(0th+1st)} & \text{(0th+5th)} \end{array}$$

Parametric error:

$$\Delta \overline{m}_c(\overline{m}_c) [\text{MeV}] = -0.5 \cdot 10^3 \frac{\text{MeV}}{\text{GeV}^4} \Delta \langle \frac{\alpha_s}{\pi} G^2 \rangle$$

(but this is only the first condensate)

QCD Sum Rules

Our approach: **error budget**

$$\Delta\alpha_s(M_z) \quad \alpha_s(M_z) = 0.1182(16) \quad \text{from PDG16}$$

$$\Delta\alpha_s(M_z) = 0.0016 \quad \longrightarrow \quad \text{from 6 MeV to 1 MeV}$$

Parametric error:

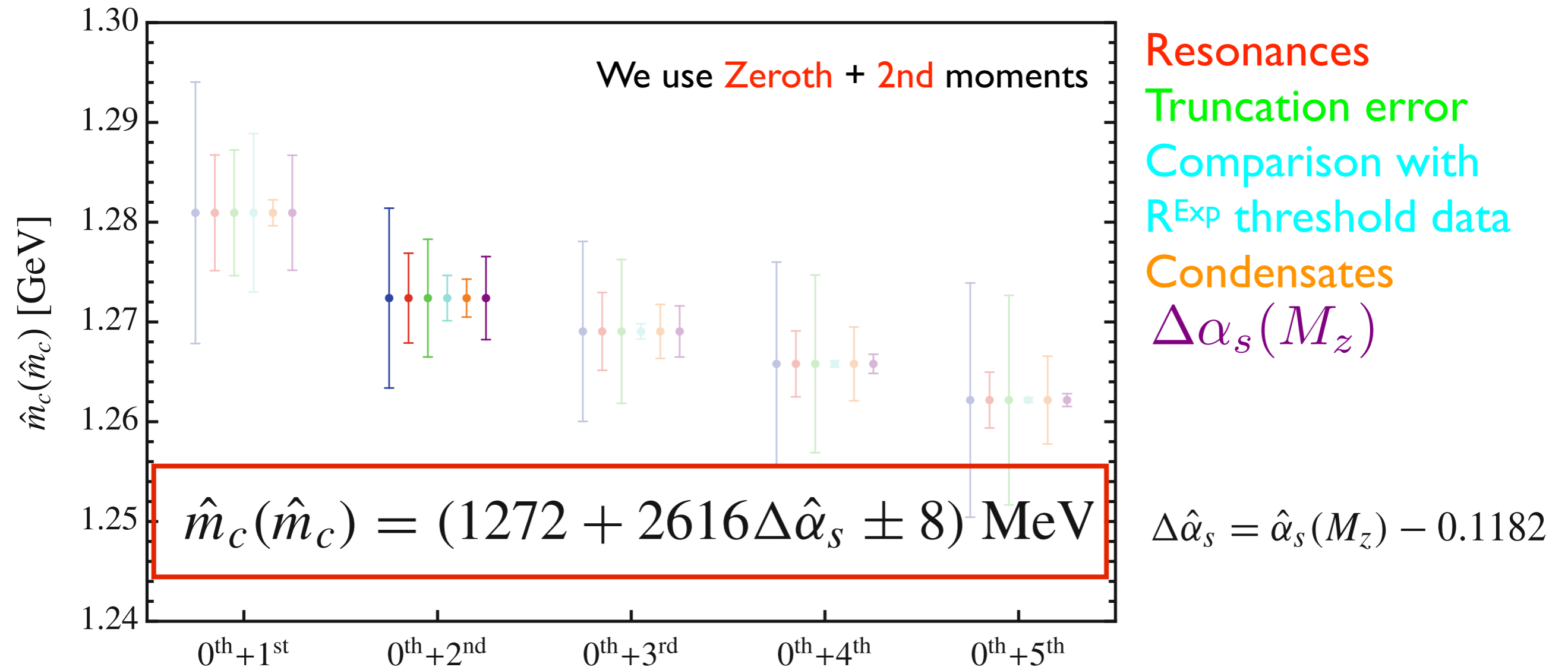
$$(0\text{th}+1\text{st}) \quad \Delta\overline{m}_c(\overline{m}_c) [\text{MeV}] = 3.6 \cdot 10^3 \Delta\alpha_s(M_z)$$

$$(0\text{th}+5\text{th}) \quad \Delta\overline{m}_c(\overline{m}_c) [\text{MeV}] = -0.4 \cdot 10^3 \Delta\alpha_s(M_z)$$

QCD Sum Rules

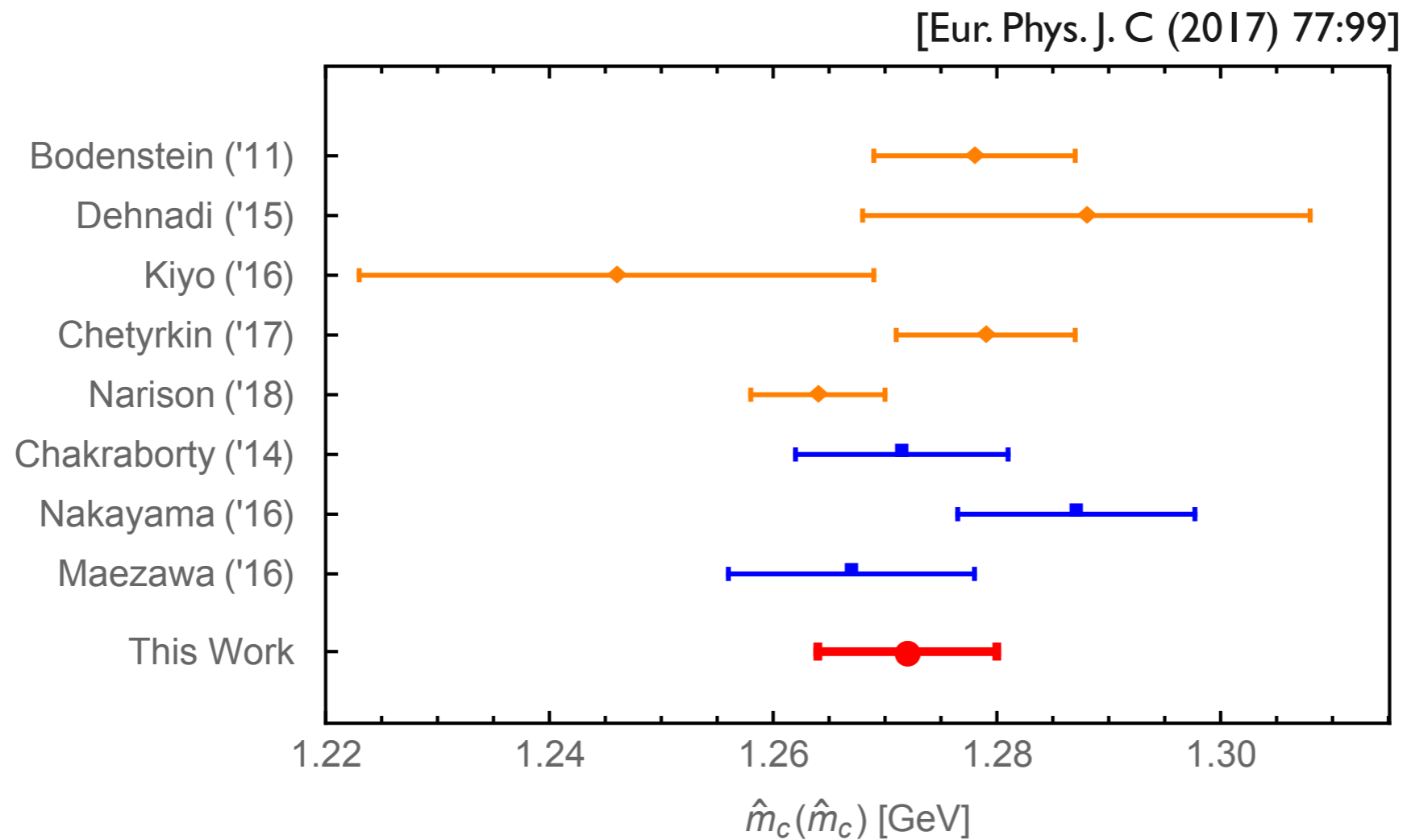
Our approach: **final result**

[J.Erler, P.M., H. Spiesberger'17]



QCD Sum Rules

results for the charm quark mass



Bottom

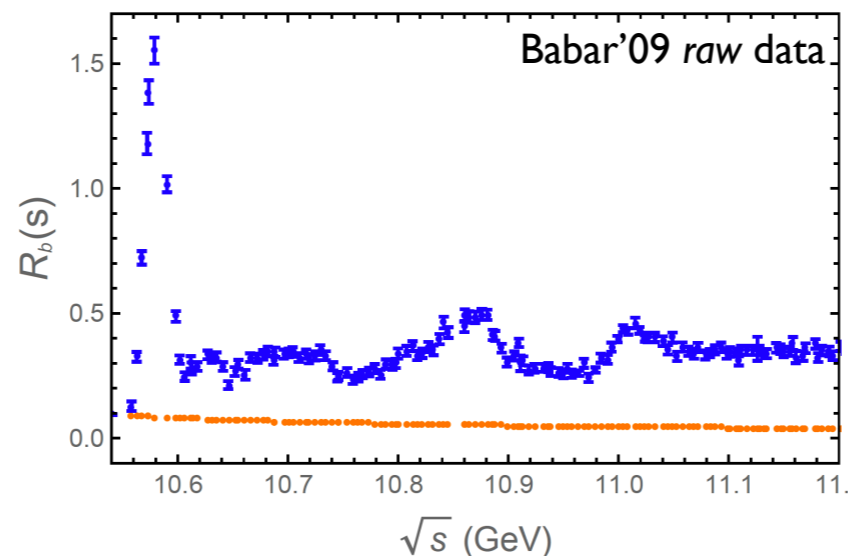
QCD Sum Rules

Bottom case

Procedure: the same as in the charm case

Main differences:

- Data from Babar '09 and Belle '15 for $R_b(s) = \sigma_b(s)/\sigma_{\mu\mu}^0$
- Condensates negligible
- Add systematically the $\Upsilon(4S)$, $\Upsilon(5S)$, $\Upsilon(6S)$



QCD Sum Rules

Bottom case

Procedure: the same as in the charm case

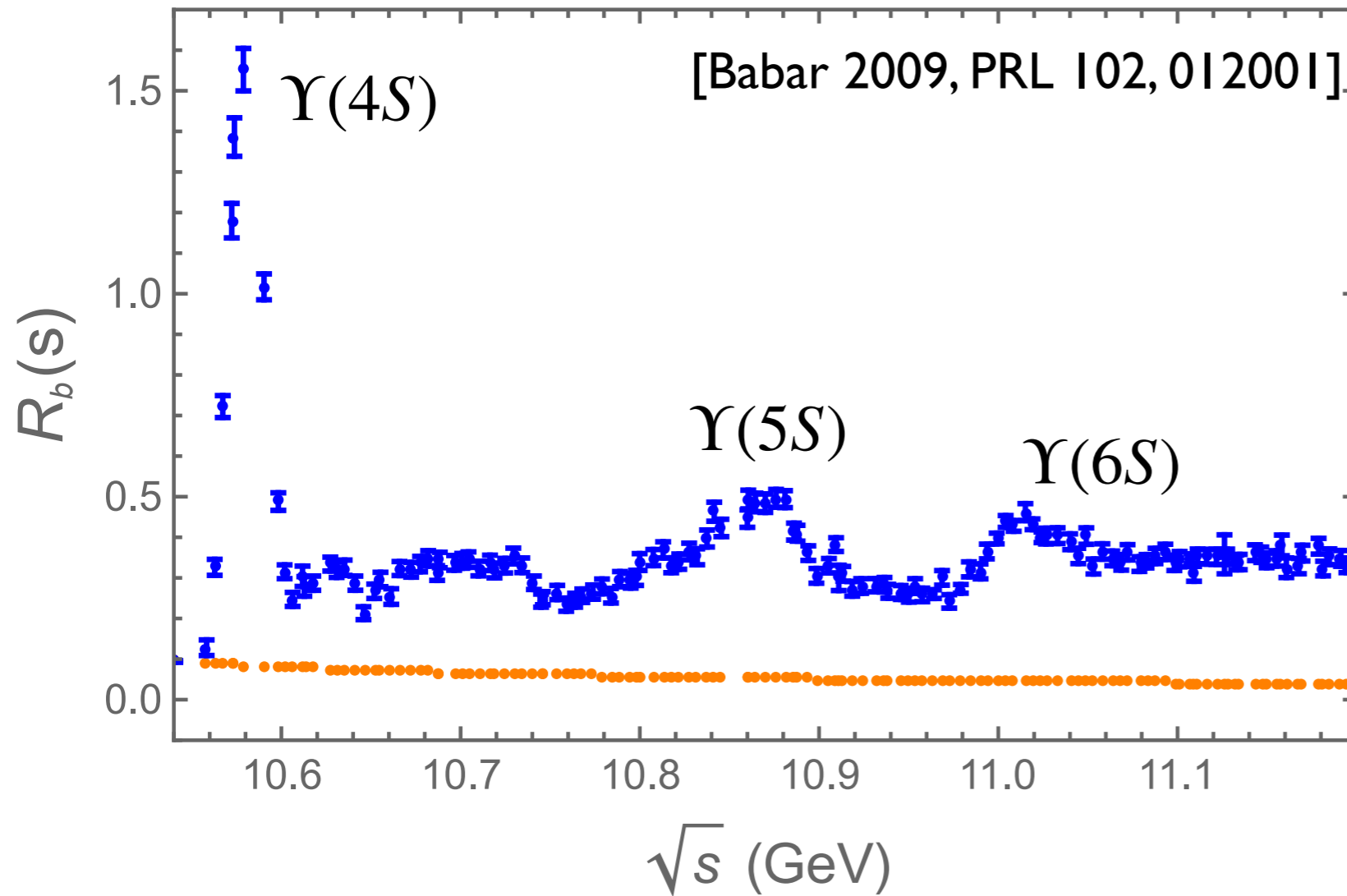
$$R_b^{\text{res}}(s) = \sum_{R=\Upsilon(1S),\Upsilon(2S),\Upsilon(3S)} \frac{9\pi}{\alpha_{\text{em}}^2(M_R)} M_R \Gamma_R^e \delta(s - M_R^2)$$

$$R_q(s) = R_q^{\text{res}}(s) + R_q^{\text{cont}}(s)$$

$$R_q^{\text{cont}}(s) = 3Q_q^2 \lambda_1^q(s) \sqrt{1 - \frac{4\hat{m}_q^2(2M)}{s'}} \left[1 + \lambda_3^q \left(\frac{2\hat{m}_q^2(2M)}{s'} \right) \right]$$

| R | M_R [GeV] | Γ_R | Γ_R^e [keV] | $\alpha_{\text{em}}^2(0)/\alpha_{\text{em}}^2(M_R)$ |
|----------------|-------------|-----------------|--------------------|---|
| $\Upsilon(1S)$ | 9.46030 | 54.02(1.25) keV | 1.340(18) | 0.931308 |
| $\Upsilon(2S)$ | 10.02326 | 31.98(2.63) keV | 0.612(11) | 0.930113 |
| $\Upsilon(3S)$ | 10.3552 | 20.32(1.85) keV | 0.443(8) | 0.929450 |
| $\Upsilon(4S)$ | 10.5794 | 20.5(2.5) MeV | 0.272(29) | 0.929009 |
| $\Upsilon(5S)$ | 10.8852 | 37 (4) MeV | 0.31(7) | 0.928415 |
| $\Upsilon(6S)$ | 11.000 | 24 (7) MeV | 0.130(30) | 0.928195 |

QCD Sum Rules



Vacuum polarization

$$(\alpha(0)/\alpha(M_R))^2 \equiv 0.93$$

Radiative tails

ISR corrections

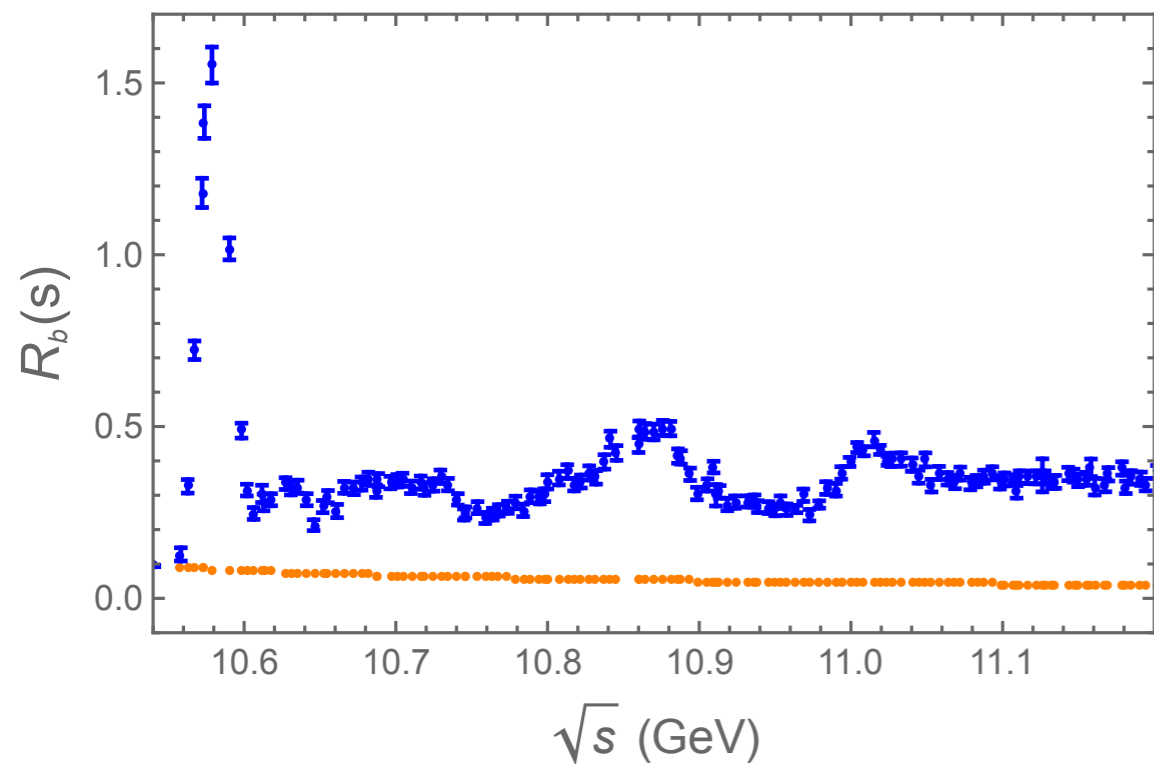
$$\hat{R}(s) = \int_{z_0}^1 \frac{dz}{z} G(z, s) R(zs)$$

$$z_0 = 10.6^2/s$$

QCD Sum Rules

Experimental moments

[Babar 2009, PRL 102, 012001]

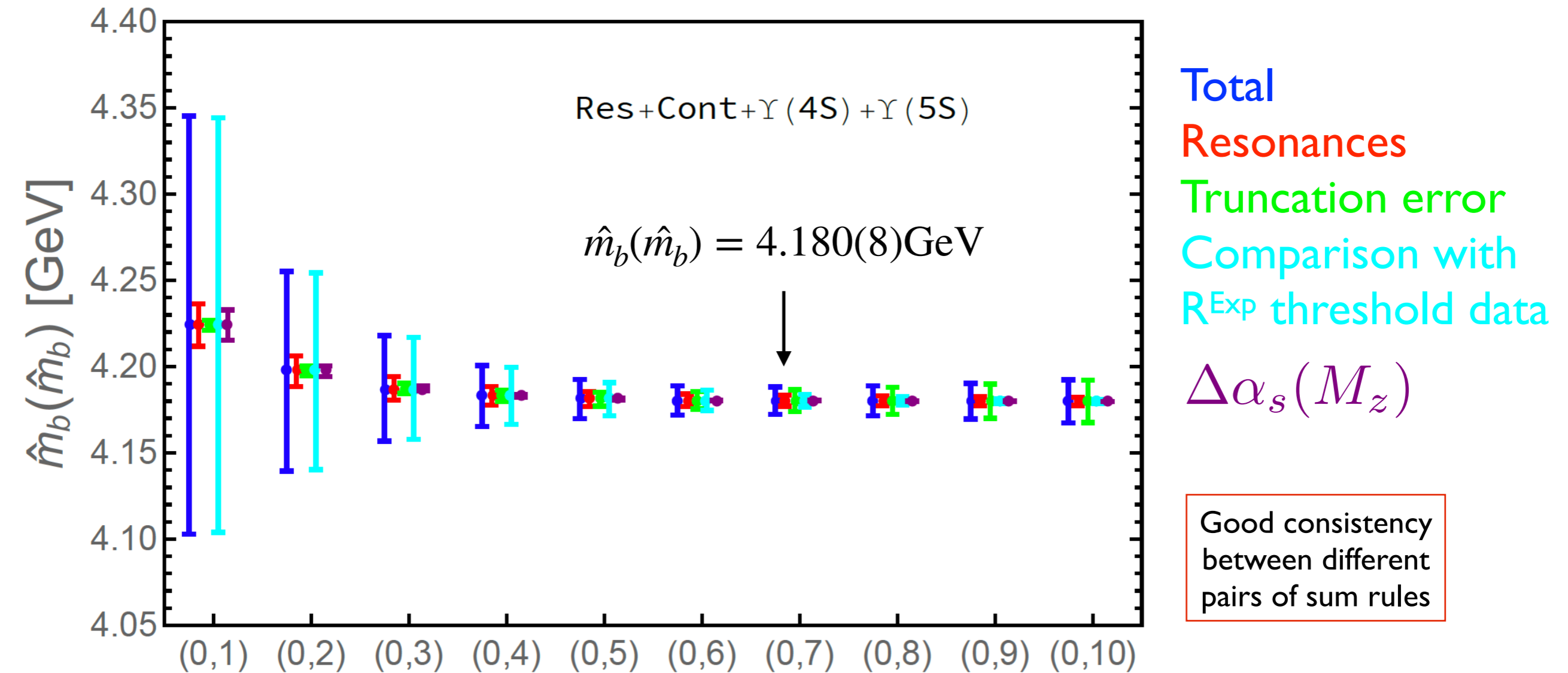


(Belle '15 data used as a crosscheck)

| n | $\mathcal{M}_n^{\text{exp}}$ | $\lambda_3^{b,\text{exp}} = 0.82(20)$ | $\lambda_3^b = 1.53$ | $\mathcal{M}_n^{\text{exp,no corr.}}$ |
|-----|------------------------------|---------------------------------------|----------------------|---------------------------------------|
| 0 | 0.446(2)(11) | 0.446(11) | 0.487 | 0.453(12) |
| 1 | 0.380(2)(9) | 0.381(9) | 0.416 | 0.384(10) |
| 2 | 0.324(1)(8) | 0.327(8) | 0.355 | 0.328(9) |
| 3 | 0.277(1)(7) | 0.280(7) | 0.304 | 0.279(7) |
| 4 | 0.237(1)(6) | 0.240(6) | 0.261 | 0.238(6) |
| 5 | 0.203(1)(5) | 0.207(5) | 0.224 | 0.204(5) |
| 6 | 0.174(1)(4) | 0.178(4) | 0.192 | 0.174(5) |
| 7 | 0.149(1)(4) | 0.153(3) | 0.165 | 0.149(4) |
| 8 | 0.128(1)(3) | 0.132(3) | 0.142 | 0.128(3) |
| 9 | 0.111(0)(3) | 0.114(2) | 0.123 | 0.110(3) |
| 10 | 0.095(0)(2) | 0.099(2) | 0.106 | 0.094(2) |

QCD Sum Rules

Our approach



QCD Sum Rules

Our approach

Explore systematically $R_b(s) = R_b^{\text{res}}(s) + R_b^{\text{cont}}(s) + R_b^{\text{res,Gamma}}(s)$

| | $\hat{m}_b(\hat{m}_b)$ [MeV] | Pair of moments |
|---|---|--|
| Only resonances below threshold + $\Upsilon(4S)$ | 4186.7 – 39.5 $\Delta\hat{\alpha}_s \pm 12.7$ 4183.8 – 68.0 $\Delta\hat{\alpha}_s \pm 9.7$ | $(\mathcal{M}_0, \mathcal{M}_9)$ $(\mathcal{M}_0, \mathcal{M}_8)$ |
| + $\Upsilon(4S) + \Upsilon(5S)$ | 4180.2 – 108.5 $\Delta\hat{\alpha}_s \pm 7.9$ | $(\mathcal{M}_0, \mathcal{M}_7)$ |
| + $\Upsilon(4S) + \Upsilon(5S) + \Upsilon(6S)$ | 4178.9 – 64.0 $\Delta\hat{\alpha}_s \pm 9.7$ | $(\mathcal{M}_0, \mathcal{M}_8)$ |

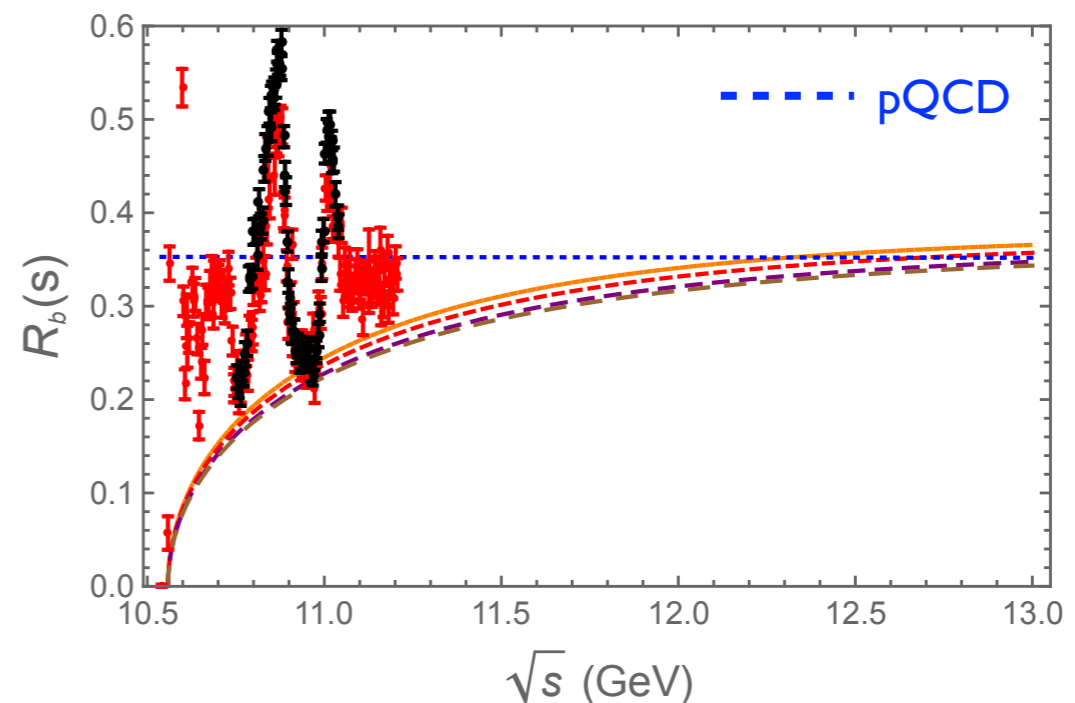
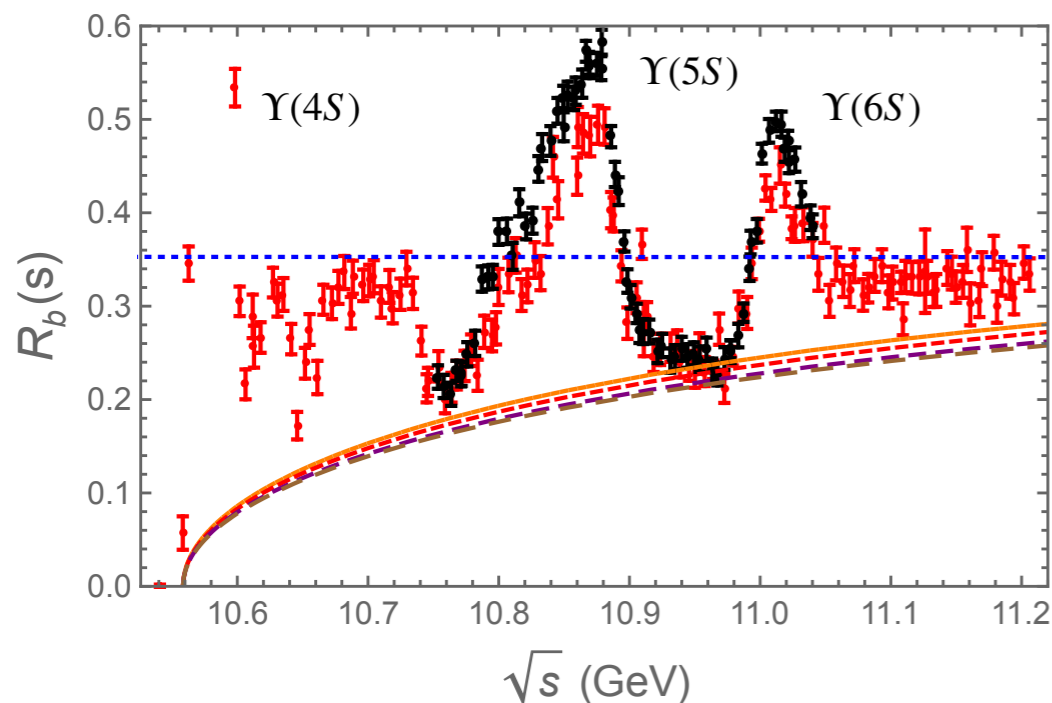
$$R_b^{\text{res,Gamma}}(s) = \sum_{R=\Upsilon(4S),\Upsilon(5S)} \frac{9\pi}{\alpha_{\text{em}}^2(M_R)} \frac{\Gamma_R^e}{M_R} \text{Gamma}(s - 4M_B^2 | \alpha, \beta)$$

$$\alpha = 1 + \frac{2}{\sqrt[3]{\pi}} \frac{(M_R^2 - 4M_B^2)^2}{\Gamma_R^2 M_R^2} \quad \beta = \frac{\alpha - 1}{M_R^2 - 4M_B^2}$$

QCD Sum Rules

Our approach

Data beyond 11.2 GeV will help reducing error: pQCD reaching at 13 GeV

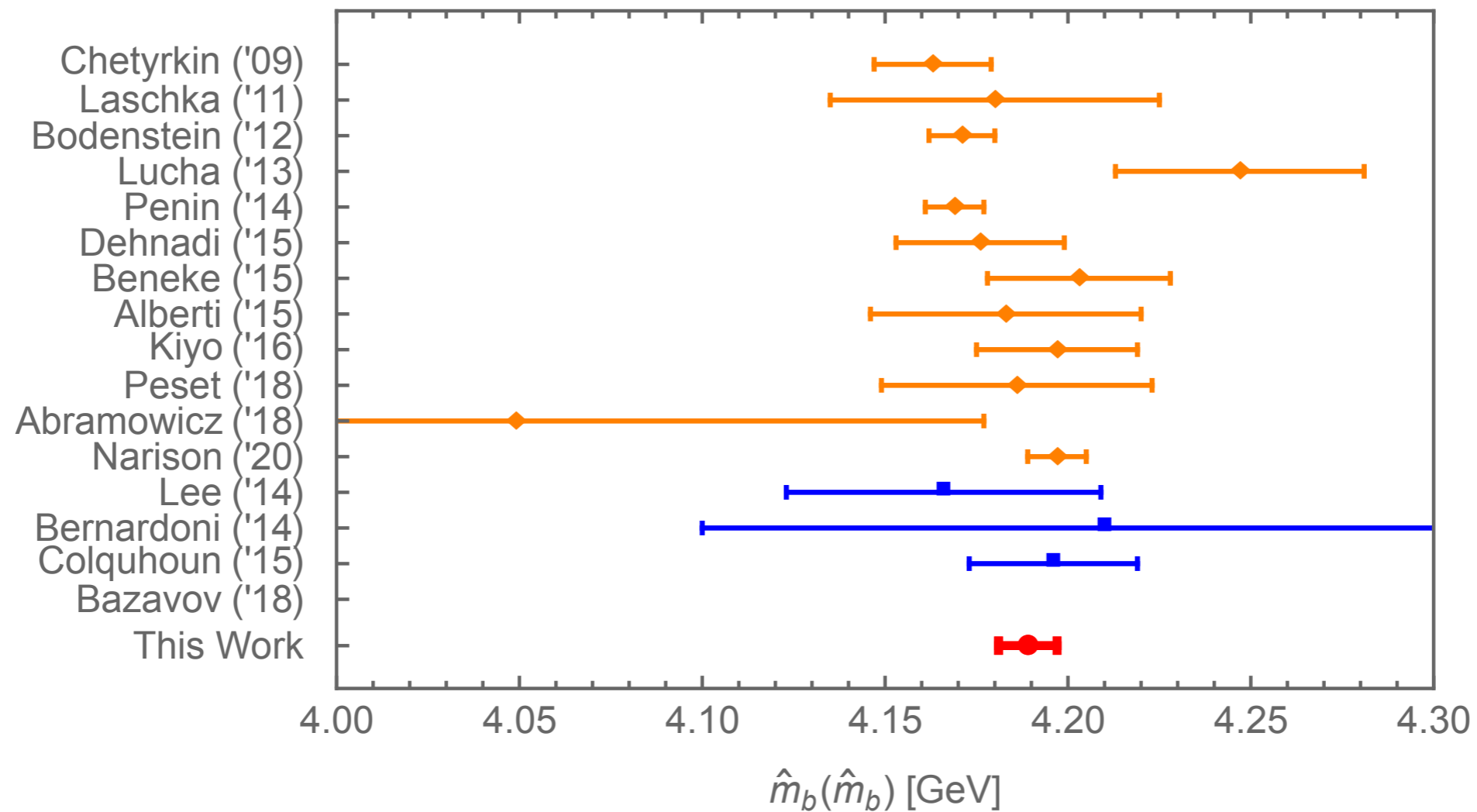


$$R_b(s) = R_b^{\text{res}}(s) + R_b^{\text{cont}}(s) + R_b^{\text{res, Gamma}}(s)$$

QCD Sum Rules

Our approach

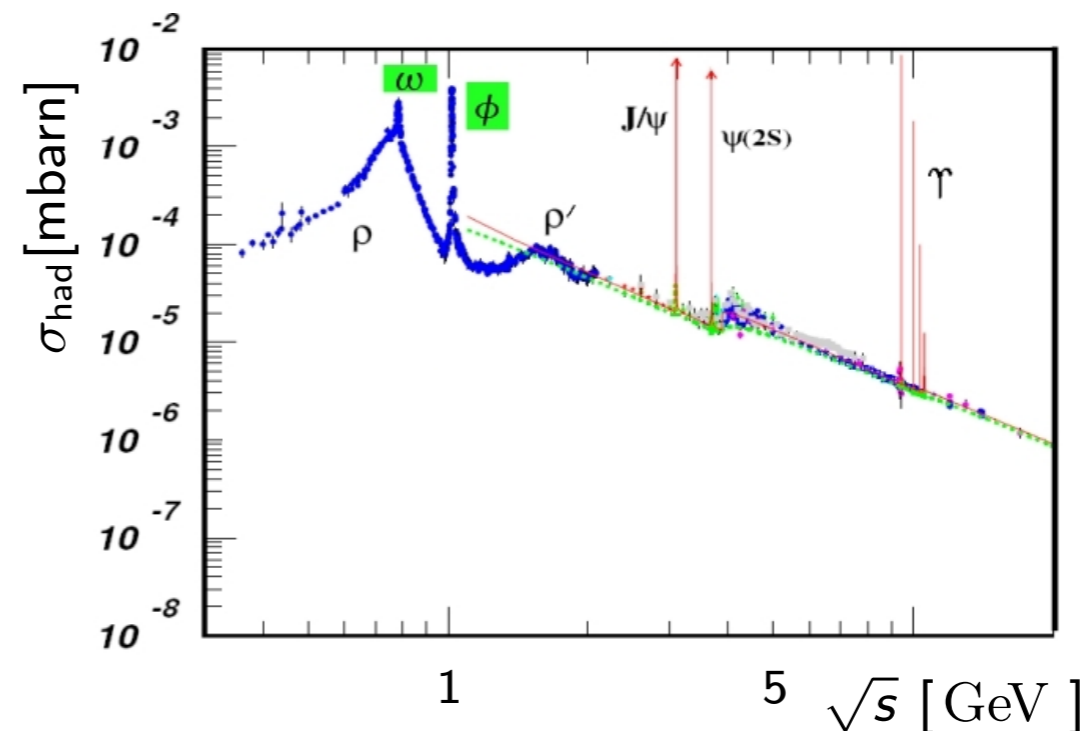
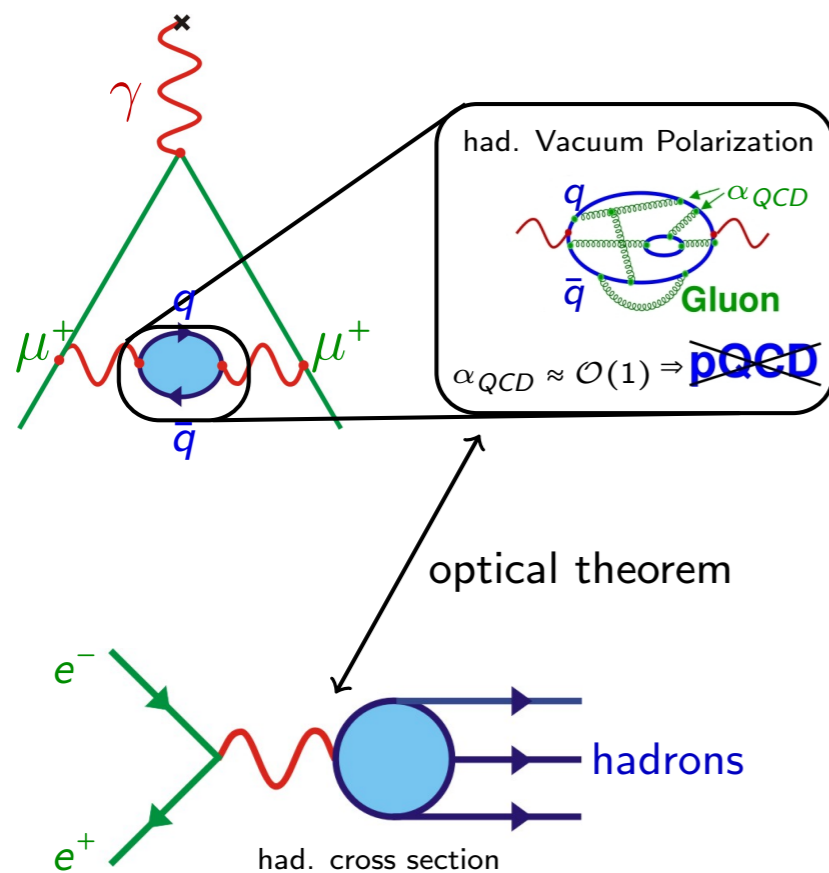
Repeat for each pair Zeroth+nth moment



Heavy-quark contribution to $(g-2)_\mu$

Hadronic Vacuum Polarization: largest source of uncertainty in $(g-2)_\mu$

Flavor decomposition may help, specially to compare with lattice QCD estimates

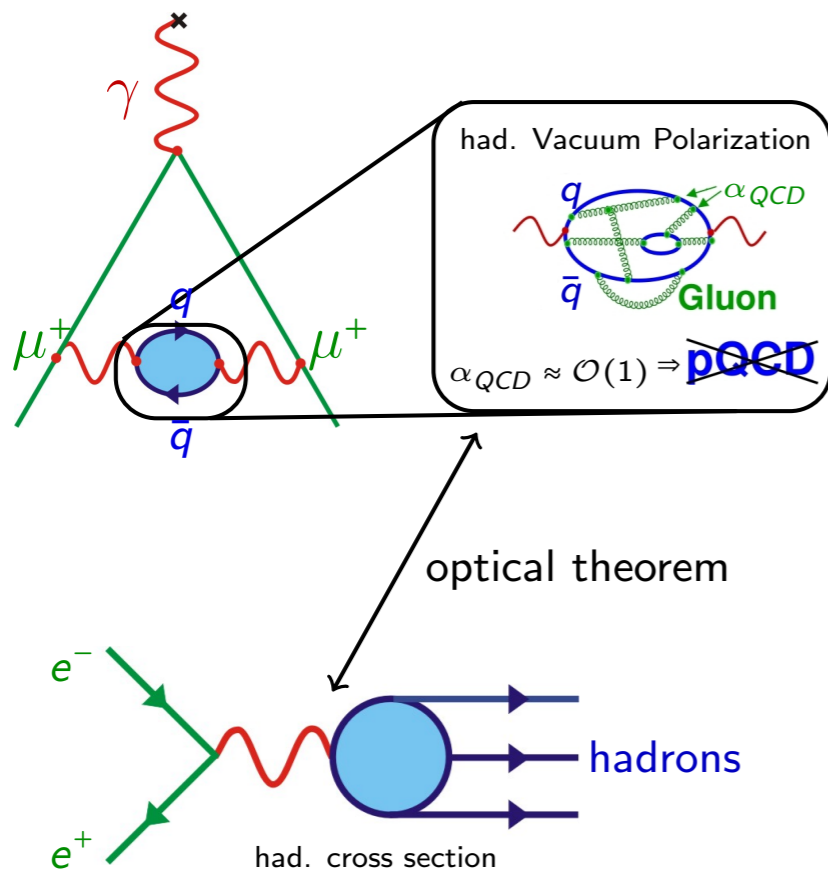


$$a_{\mu, LO}^{\text{had}} = \frac{1}{4\pi^3} \int_{m_{\pi^0}^2}^{\infty} ds K(s) \underline{\sigma_{\text{had}}(s)}$$

Heavy-quark contribution to $(g-2)_\mu$

Hadronic Vacuum Polarization: largest source of uncertainty in $(g-2)_\mu$

Flavor decomposition may help, specially to compare with lattice QCD estimates



$$a_{\mu, LO}^{\text{had}} = \frac{1}{4\pi^3} \int_{m_{\pi^0}^2}^{\infty} ds K(s) \sigma_{\text{had}}(s)$$

$$a_{\mu}^{\text{charm}} = 14.36(23) \times 10^{-10} \quad a_{\mu}^{\text{bottom}} = 0.30(2) \times 10^{-10}$$

$$a_{\mu}^{\text{charm-lattice}} = 14.6(1) \times 10^{-10} \quad a_{\mu}^{\text{bottom-lattice}} = 0.27(4) \times 10^{-10}$$

from Borsanyi et al, *Nature* **593**, 51–55 (2021)

| | central value | total error | resonances | $\Delta\lambda_3$ | $\Delta\alpha_s$ | Condensates | Truncation |
|---------------------------|---------------|-------------|------------|-------------------|------------------|-------------|------------|
| a_{μ}^{charm} | 1.436 | 0.023 | 0.012 | 0.018 | 0.005 | 0.001 | 0.004 |
| a_{μ}^{bottom} | 2.978 | 0.171 | 0.012 | 0.170 | 0.005 | — | 0.004 |

Conclusions and Outlook

- Using SR technique + *zeroth moment* (very sensitive to the continuum) + data on charm resonances below threshold + continuum exploiting self-consistency among different moments:

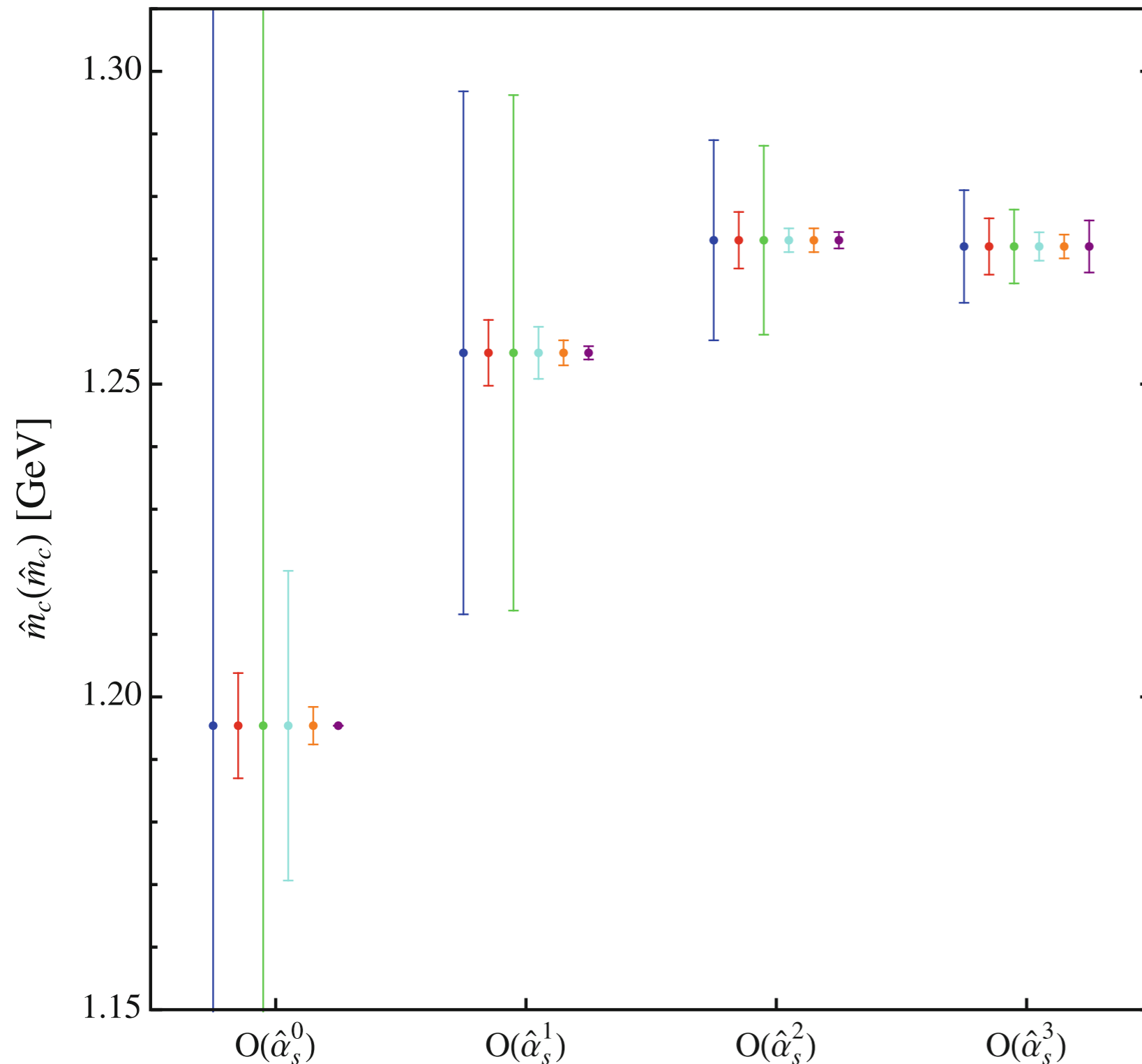
$$\hat{m}_c(\hat{m}_c) = 1.272(9)\text{GeV}$$

$$\hat{m}_b(\hat{m}_b) = 4.180(8)\text{GeV}$$

- Error sources are understood: seems a clear roadmap for improvements
- Impact on $(g-2)_\mu$ from heavy quarks: $a_\mu^{\text{charm+bottom}} = 14.66(23) \times 10^{-10}$

Thanks!

QCD Sum Rules



α_s expansion

for the **Zerth** + **2nd** moments

Total

Resonances

Truncation error

Comparison with
 R^{Exp} threshold data

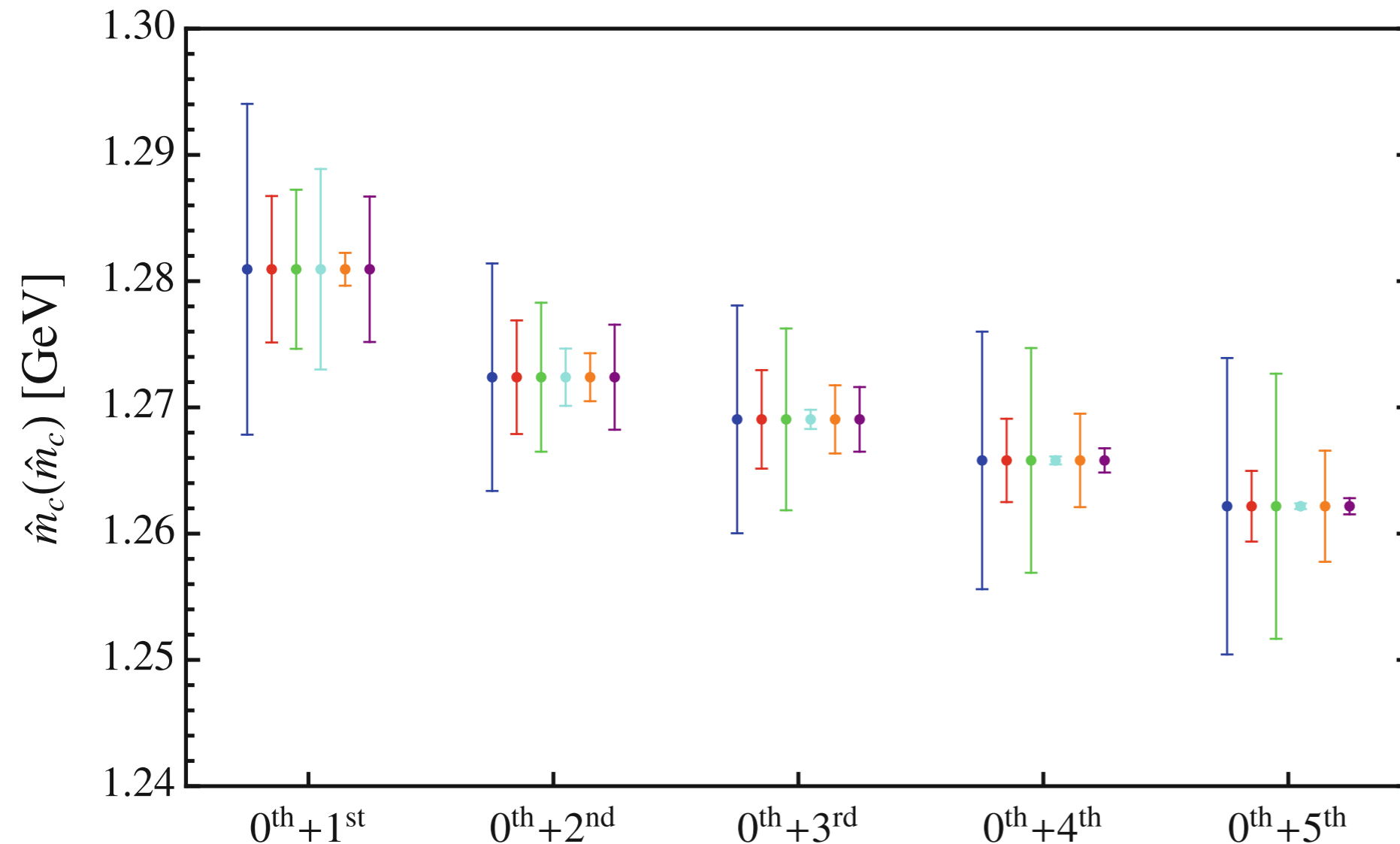
Condensates

$\Delta\alpha_s(M_z)$

QCD Sum Rules

Our approach

What pair/result to choose?



Resonances

Truncation error

Comparison with R^{Exp} threshold data

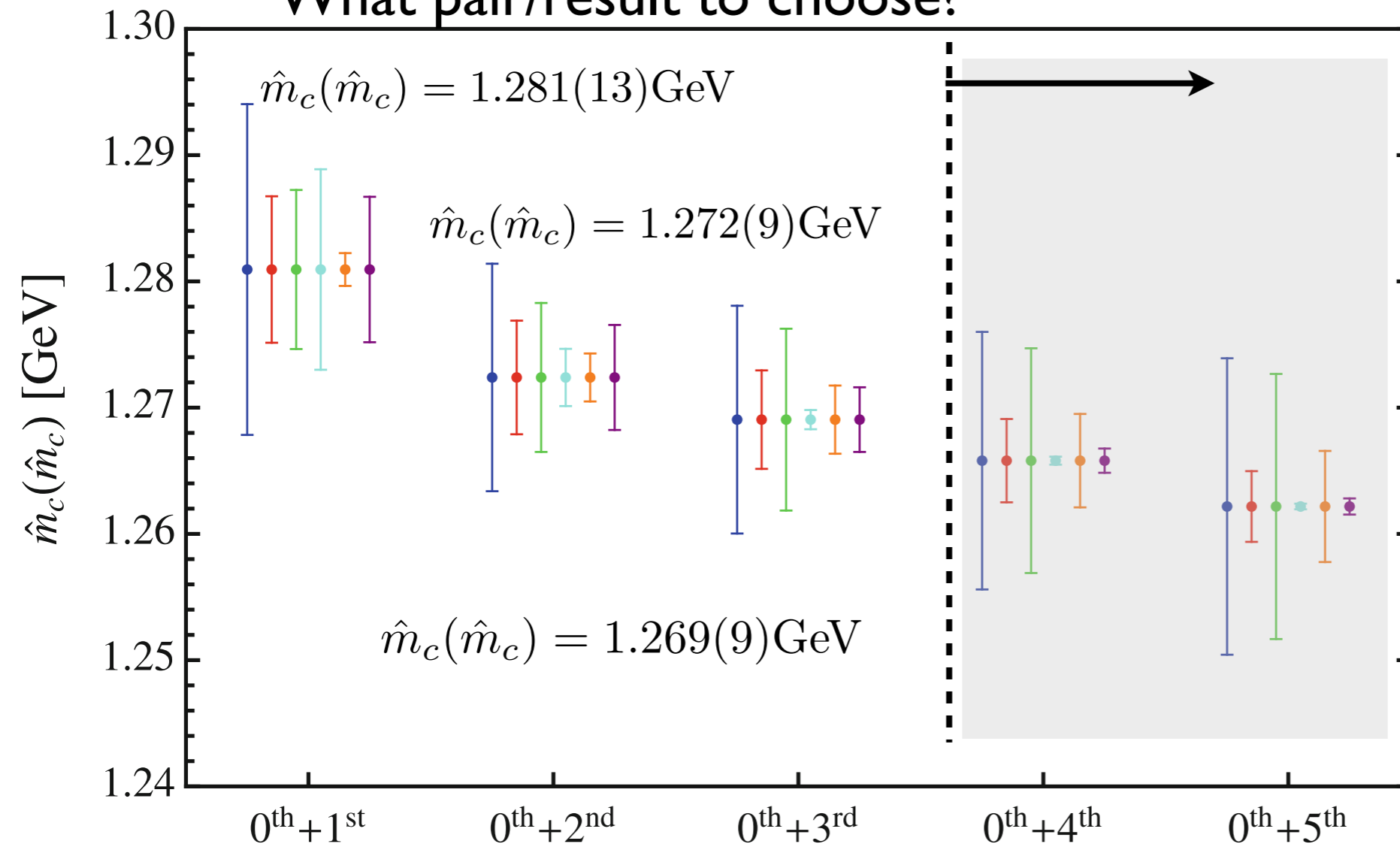
Condensates

$\Delta\alpha_s(M_z)$

QCD Sum Rules

Our approach

What pair/result to choose?



Resonances
 Truncation error
 Comparison with R^{Exp} threshold data
 Condensates
 $\Delta\alpha_s(M_z)$

Large condensate effects
 +
 new condensates will matter

QCD Sum Rules

Our approach: **more than two moments?**

Define a χ^2 function:

$$\chi^2 = \frac{1}{2} \sum_{n,m} (\mathcal{M}_n - \mathcal{M}_n^{\text{pQCD}}) (\mathcal{C}^{-1})^{nm} (\mathcal{M}_m - \mathcal{M}_m^{\text{pQCD}}) + \chi_c^2$$

$$\mathcal{C} = \frac{1}{2} \sum_{n,m} \rho^{\text{Abs}(n-m)} \Delta \mathcal{M}_n^{(4)} \Delta \mathcal{M}_m^{(4)} \quad \rho \text{ a correlation parameter}$$

$$\chi_c^2 = \left(\frac{\Gamma_{J/\Psi(1S)}^e - \Gamma_{J/\Psi(1S)}^{e,\text{exp}}}{\Delta \Gamma_{J/\Psi(1S)}^e} \right)^2 + \left(\frac{\Gamma_{\Psi(2S)}^e - \Gamma_{\Psi(2S)}^{e,\text{exp}}}{\Delta \Gamma_{\Psi(2S)}^e} \right)^2 +$$

$$\left(\frac{\hat{\alpha}_s(M_z) - \hat{\alpha}_s(M_z)^{\text{exp}}}{\Delta \hat{\alpha}_s(M_z)} \right)^2 + \left(\frac{\langle \frac{\alpha_s}{\pi} G^2 \rangle - \langle \frac{\alpha_s}{\pi} G^2 \rangle^{\text{exp}}}{\Delta \langle \frac{\alpha_s}{\pi} G^2 \rangle} \right)^2$$

QCD Sum Rules

Our approach: **more than two moments?**

Define a χ^2 function:

| | Constraints | $(\mathcal{M}_0, \mathcal{M}_1, \mathcal{M}_2)_\rho$ | $\mathcal{M}_0, (\mathcal{M}_1, \mathcal{M}_2)_\rho$ | $\mathcal{M}_0, (\mathcal{M}_1, \mathcal{M}_2, \mathcal{M}_3)_\rho$ |
|------------------------------|-------------|--|--|---|
| ρ | | -0.06 | -0.05 | 0.32 |
| $\hat{m}_c(\hat{m}_c)$ [GeV] | | 1.275(8) | 1.275(8) | 1.271(7) |
| λ_3^c | | 1.19(8) | 1.19(8) | 1.19(7) |
| $\Gamma_{J/\psi}^e$ [keV] | 5.55(14) | 5.57(14) | 5.57(14) | 5.59(14) |
| $\Gamma_{\psi(2S)}^e$ [keV] | 2.36(4) | 2.36(4) | 2.36(4) | 2.36(4) |
| C_G [GeV ⁴] | 0.005(5) | 0.005(5) | 0.005(5) | 0.004(5) |
| $\hat{\alpha}_s(M_z)$ | 0.1182(16) | 0.1178(15) | 0.1178(15) | 0.1173(15) |

QCD Sum Rules

Our approach: **more than two moments?**

Preferred scenario:

| | 0th + (1st + 2nd) _{ρ} $\Delta\hat{m}_c(\hat{m}_c)$ [MeV] | (0th + 2nd) $\Delta\hat{m}_c(\hat{m}_c)$ [MeV] |
|-----------------------------|--|---|
| Central value | 1274.5 | 1272.4 |
| $\Delta\Gamma_{J/\psi}^e$ | 5.9 | 4.5 |
| $\Delta\Gamma_{\Psi(2S)}^e$ | 1.4 | 0.4 |
| Truncation | — | 5.9 |
| $\Delta\lambda_3^c$ | 3.0 | 2.3 |
| Condensates | 1.1 | 1.9 |
| $\Delta\hat{\alpha}_s(M_Z)$ | 5.4 | 4.2 |
| Total | 8.7 | 9.0 |

QCD Sum Rules

