

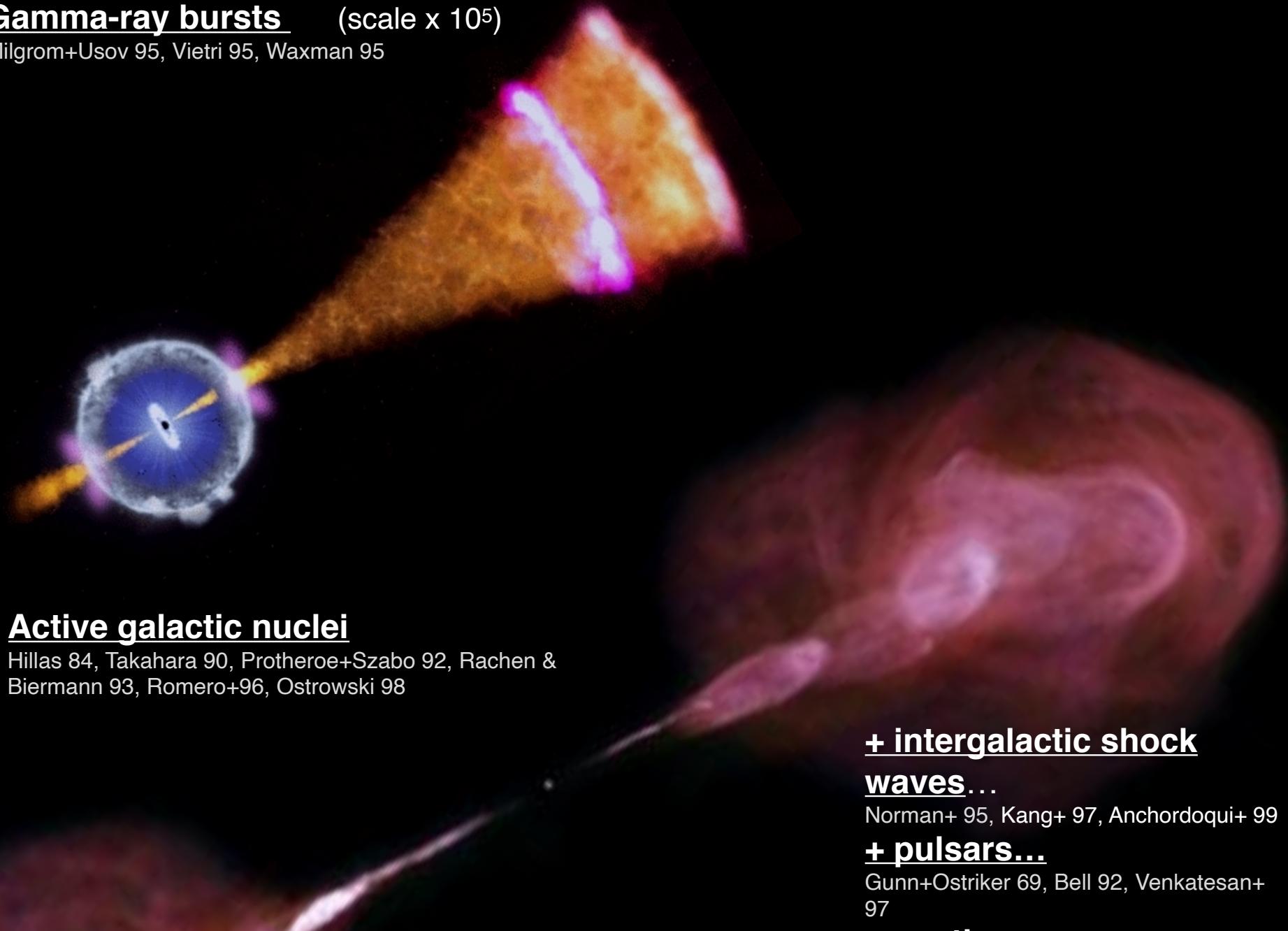
Particle acceleration to UHE

Martin Lemoine
Institut d'Astrophysique de Paris
CNRS – Sorbonne Université

Usual suspects... twenty years ago...

Gamma-ray bursts (scale x 10⁵)

Milgrom+Usov 95, Vietri 95, Waxman 95



Active galactic nuclei

Hillas 84, Takahara 90, Protheroe+Szabo 92, Rachen & Biermann 93, Romero+96, Ostrowski 98

+ intergalactic shock waves...

Norman+ 95, Kang+ 97, Anchordoqui+ 99

+ pulsars...

Gunn+Ostriker 69, Bell 92, Venkatesan+ 97

Usual suspects... now...

Powerful transients: high and low luminosity GRBs, relativistic SNe, fast spinning magnetars/pulsars...

High luminosity GRB:

Milgrom + Usov 95, Vietri 95, Waxman 95, Rachen + Meszaros 96, Gallant + Achterberg 99, Pelletier + Kersale 00, Dermer + Humi 01, Waxman 01, Scully + Stecker 02, Gialis + Pelletier 03, Waxman 04, Rieger + Duffy 06, Asano+ 09, 10, Razzaque+ 10, Giannios 10, Eichler + Pohl 11, Metzger+ 11, Baerwald+ 15, Globus+ 15, Asano + Meszaros 16, Samuelsson+ 19

Low-luminosity GRBs

/ Relativistic SNe:

Murase+ 06, Budnik+ 08, Wang+ 08, Liu+ 11, Chakraborti+ 11 , Liu + Wang 12, Zhang + Murase 18, Zhang+ 18, Boncioli+ 19

Magnetar WNe/ Pulsar WNe :

Bell 92, Blasi+ 00, de Gouveia dal Pino + Lazarian 01, Arons 03, Vietri+ 03, Fang+ 13, Lemoine+ 15, Fang+ 18, Kirk + Giacinti 19

Active galactic nuclei:

AGN Core:

Protheroe + Szabo 92, Boldt + Ghosh 99, Levinson 00, Levinson + Boldt 00, Torres+ 02, Dempsey + Rieger 09, Istomin + Sol 09, Neronov + 09, Rieger + Aharonian 09, Dutan + Caramete 14, Moncada+ 17

Blazar/Radio-galaxy jet+lobe:

Takahara 90, Rachen + Biermann 93, Ostrowski 98, Farrar + Piran 00, Tinyakov + Tkachev 01, Gorbunov+ 02, Lyutikov + Ouyed 07, Atoyan+Dermer 08, Gorbunov + 08, Dermer+ 09, Hardcastle+ 09, O'Sullivan+ 09, Dermer + Razzaque 10, Gopal-Krishna+ 10, Biermann+ deSouza 12, Murase+ 12, Ptuskin+ 13,Caprioli 15, Wang + Loeb17, Eichmann+ 18, Liu+ 17, Resconi+ 17, Kimura+ 18, Matthews+ 18,19,Fang + Murase18,Rieger 19

+ intergalactic shock waves (clusters, starbursts)...

Norman+95, Kang+ 97, Anchordoqui+ 99, + 01, Murase+ 08, Kotera +09, Malkov +11, Anchordoqui 18, Romero+ 18

+ other transients: e.g. tidal disruption

Connections to many (contemporaneous) fields of physics...

High-energy astrophysics:

e.g. the origin, dynamics, energy content and dissipative physics of (relativistic) jets...

e.g. environments of compact objects...

Experimental astroparticle physics:

... a data-starved field of research!

e.g., chemical composition and anisotropies fundamental clues...

Multi-messenger astrophysics:

e.g. neutrinos unambiguous signatures of hadron acceleration...

e.g. photons indirect probes of the physical conditions...

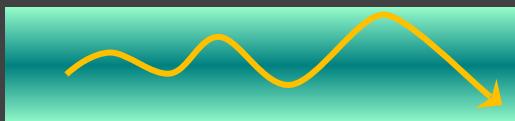
Extreme plasma astrophysics:

e.g. (highly?) magnetized, highly energetic collisionless plasmas with bulk (or random?) velocities $v \sim c$...

e.g. plasmas in strong gravity/radiation fields...

Various (more or less known) acceleration scenarios...

Shear acceleration



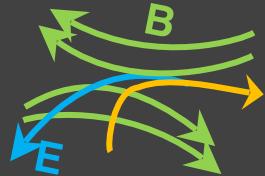
Shock acceleration



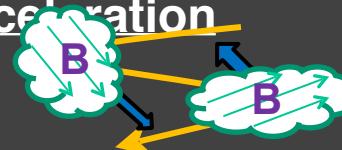
Linear acceleration in gaps



Reconnection

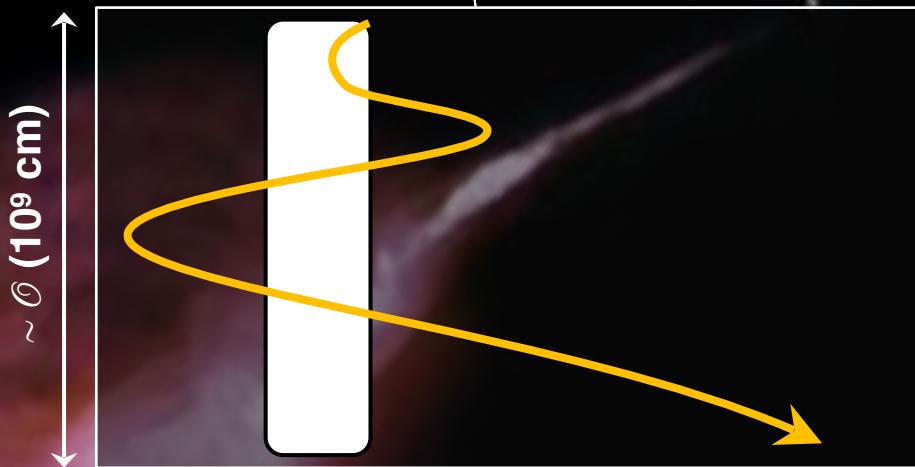
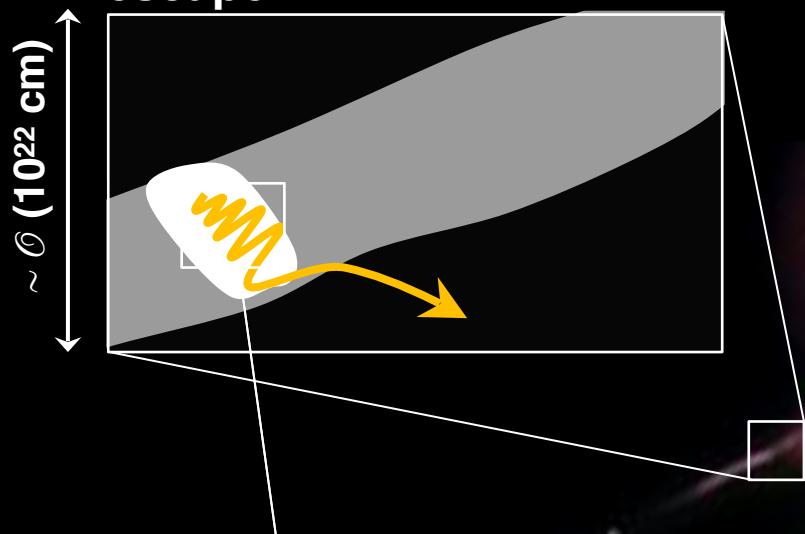


Turbulent acceleration



A challenge of scales...

Large scales: particle escape...
 $E = 10^{20} \text{ eV } ZB_{100\mu\text{G}} L_{\text{kpc}}$



Meso-scales: particle acceleration

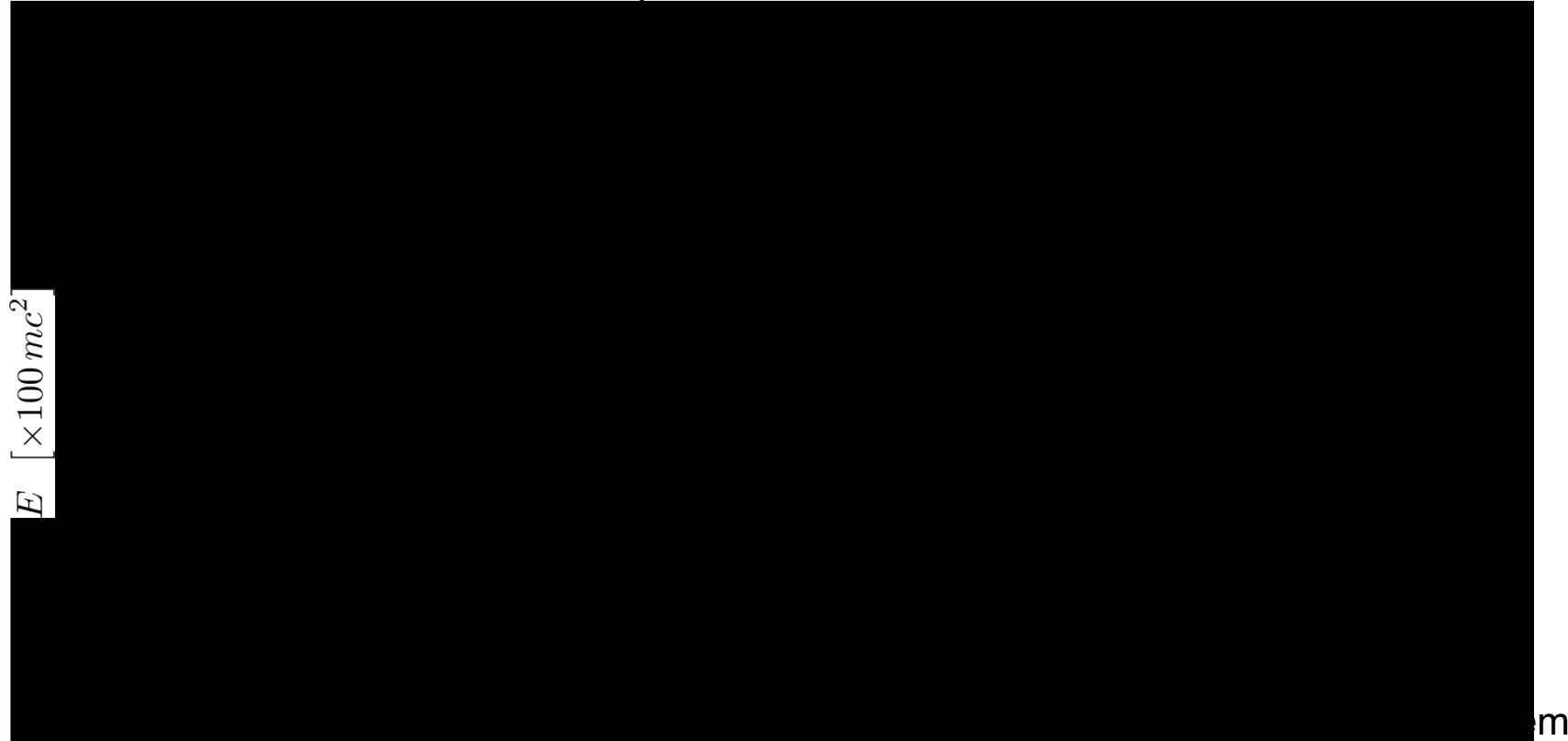
Micro-scales: particle injection...

$c/\omega_c \sim 10^8 \text{ cm } B_{100\mu\text{G}}$

Numerical simulations of particle acceleration...

... HPC numerical simulations (Monte Carlo, particle-in-cell, MHD+PIC, hybrid ...) allow to explore the physics of acceleration to unprecedented details...

→ explore the issue of injection and backreaction of accelerated particles on the acceleration process



... **an important caveat:** ab initio simulations remain limited to small scales...

Top (of the energy/length/time scales) - down...

Two critical properties of a source of UHECRs:

→ **a high output of cosmic rays**: $\dot{E}_{\text{UHECR}} \sim 10^{44} \text{ erg/Mpc}^3/\text{yr}$

Katz+10, ...

... a non-trivial constraint: e.g. $L_{\text{UHE}}/L_\gamma \sim 10$ for HL GRBs...

e.g. $L_{\text{UHE}}/L \sim \mathcal{O}(1\text{-}10\%)$ for radio-galaxies...

→ **a high magnetic luminosity**:

$$L_{\text{tot}} \gtrsim 10^{45} \text{ erg/s} \dots \left(\frac{t_{\text{acc}}}{t_g} \right)^2 \left(\frac{E/Z}{10^{20} \text{ eV}} \right)^2$$

Hillas 84, Blandford 00, Aharonian 02,

Waxman 04, Lyutikov+Ouyed 07,

Lemoine+Waxman 09, ...

... a non-trivial constraint... Chemical composition (Z) controls the phenomena

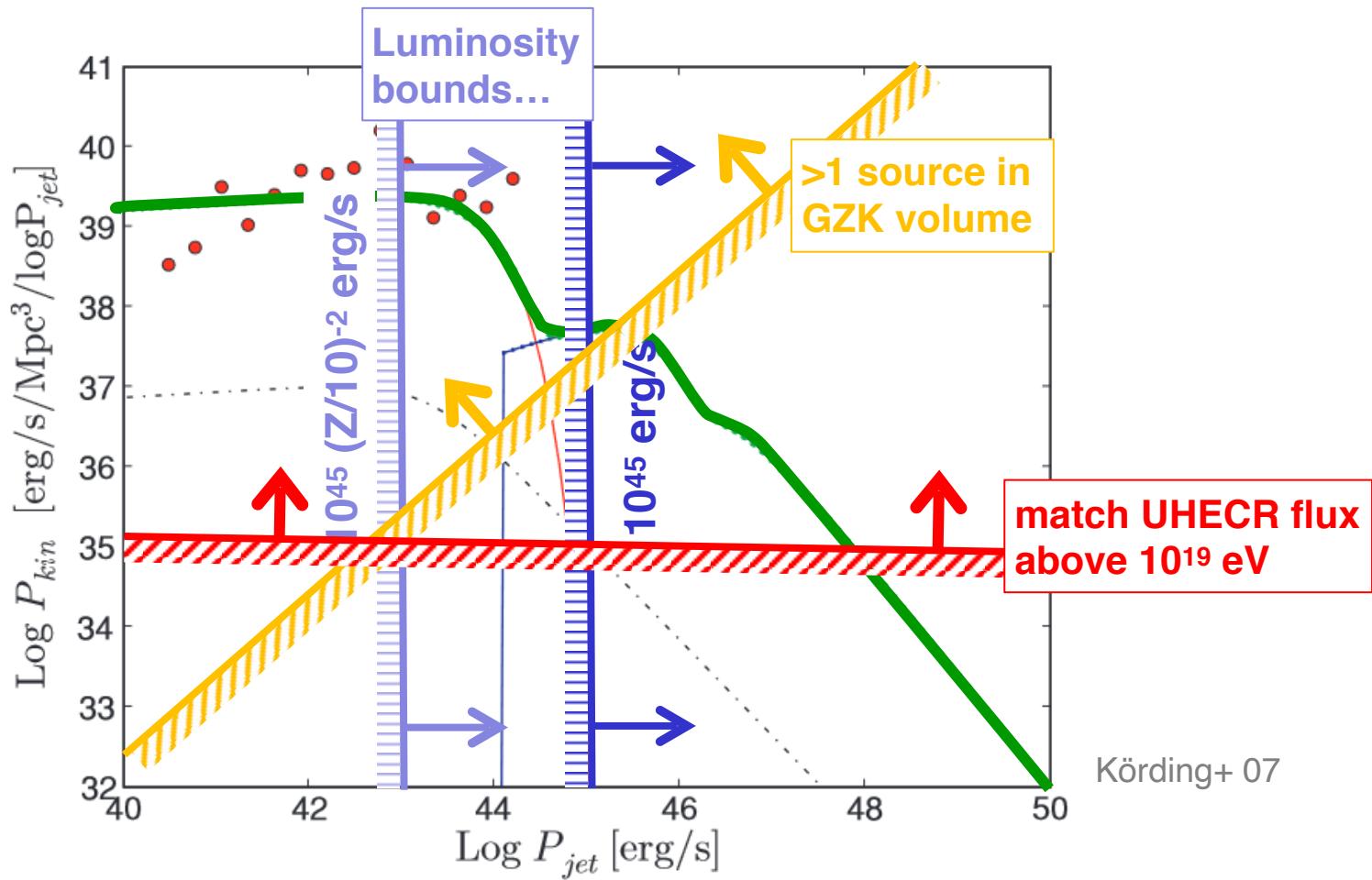
... to go further:

$$\dots \left(\frac{t_{\text{acc}}}{t_g} \right)^2 \propto \frac{t_{\text{scatt}}}{t_g} \begin{cases} u^2 & (u \gg 1) \\ 1/u & (u \ll 1) \end{cases} \quad (\text{u 4-velocity})$$

⇒ favors relativistic (mildly?) sources...

... to reach the confinement energy: $t_{\text{acc}} \sim t_g$!

Top - down... zooming in on the radio-galaxy population...



local radio-galaxies barely satisfy the luminosity bound: accelerate Z ~ 10+ nuclei?

Top - down... some lessons learned...

Two critical properties of a source of UHECRs:

→ **a high output of cosmic rays:** $\dot{E}_{\text{UHECR}} \sim 10^{44} \text{ erg/Mpc}^3/\text{yr}$

→ **a high magnetic luminosity:** $L_{\text{tot}} \gtrsim 10^{45} \text{ erg/s} \dots \left(\frac{t_{\text{acc}}}{t_g} \right)^2 \left(\frac{E/Z}{10^{20} \text{ eV}} \right)^2$

... leading contenders, for accelerating intermediate nuclei ($Z \sim 10$):

→ powerful radio-galaxies, $L \sim 10^{44} \text{ erg/s}$, $u \sim 1 \dots$

→ relativistic supernovae (LLGRB), $L \sim 10^{44} \text{ erg/s}$, $u \sim 1 \dots$

... need extreme sources for accelerating light nuclei ($Z \sim 1$):

→ gamma-ray bursts, fast-spinning magnetar/pulsar wind nebulae...

→ most powerful FRII like radio-galaxies

Bottom – up : the microphysics of acceleration

$$\rightarrow \text{Lorentz force: } \frac{dp}{dt} = q \left(\mathbf{E} + \frac{\mathbf{v}}{c} \times \mathbf{B} \right)$$

... recall however that \mathbf{E} and \mathbf{B} transform into each other in a change of frame...

... rely on Lorentz scalars $\mathbf{E} \cdot \mathbf{B}$ and $\mathbf{E}^2 - \mathbf{B}^2$...

Case 1: highly conducting plasma (ideal MHD) $\mathbf{E} \cdot \mathbf{B} = 0$ and $\mathbf{E}^2 - \mathbf{B}^2 < 0$

→ **generic for UHECR**: corresponds to ideal Ohm's law $\mathbf{E} = -\mathbf{v}_p \times \mathbf{B}/c$...

(on timescales \gg micro, comoving electric field is screened out)

→ **Fermi-type scenarios: magnetized turbulence, shear flows, shock waves**

Case 2: non-MHD behavior...

$\mathbf{E} \cdot \mathbf{B} \neq 0$ or $\mathbf{E}^2 - \mathbf{B}^2 > 0$

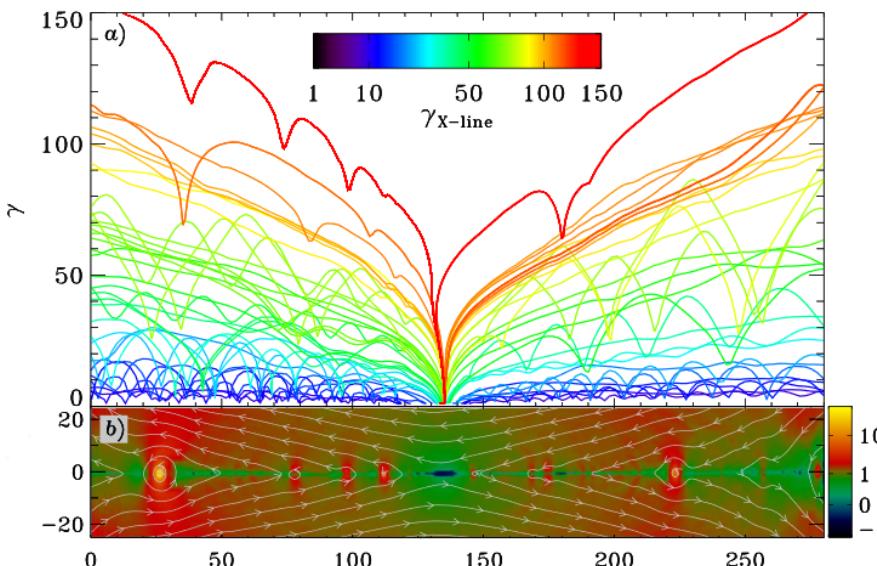
→ acceleration can proceed unbounded along \mathbf{E} (or at least $\mathbf{E}_{||}$)...

→ **gaps in BH magnetospheres**: possible, but many open issues, e.g. Levinson
00, Rieger 19

→ **reconnection**: unlikely, max energy in (relativistic) reconnection $\propto \sigma = u_B/u_p$

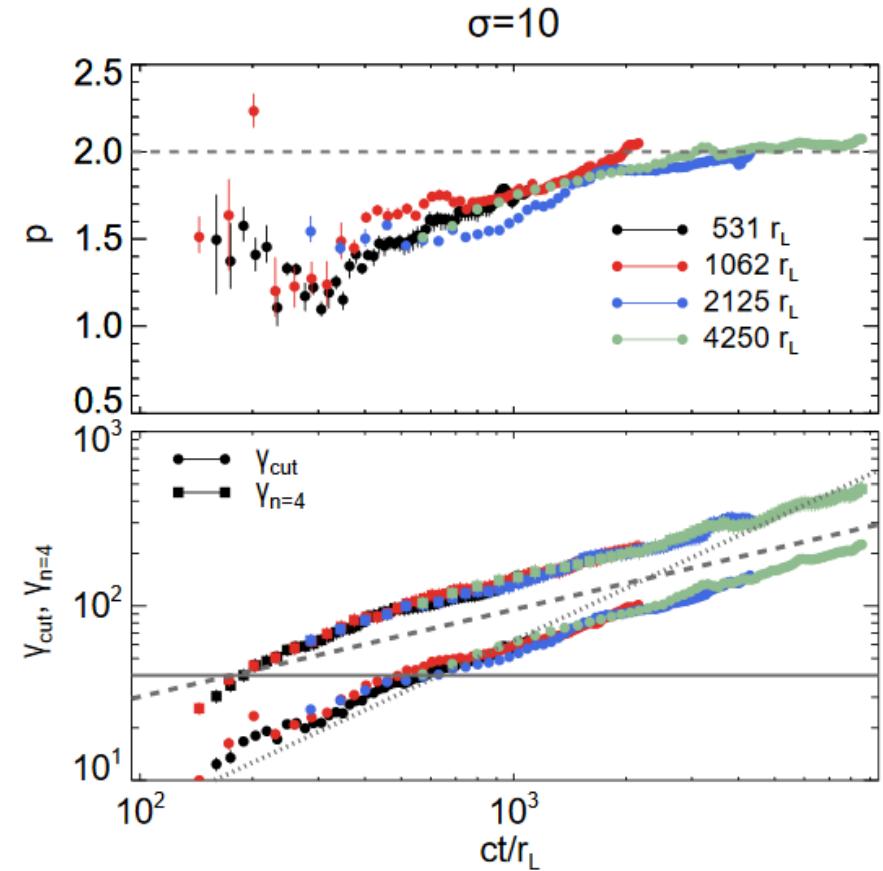
Reconnection vs acceleration to VHE

→ acceleration proceeds in two stages



(1) acceleration at and around
the annihilation regions...

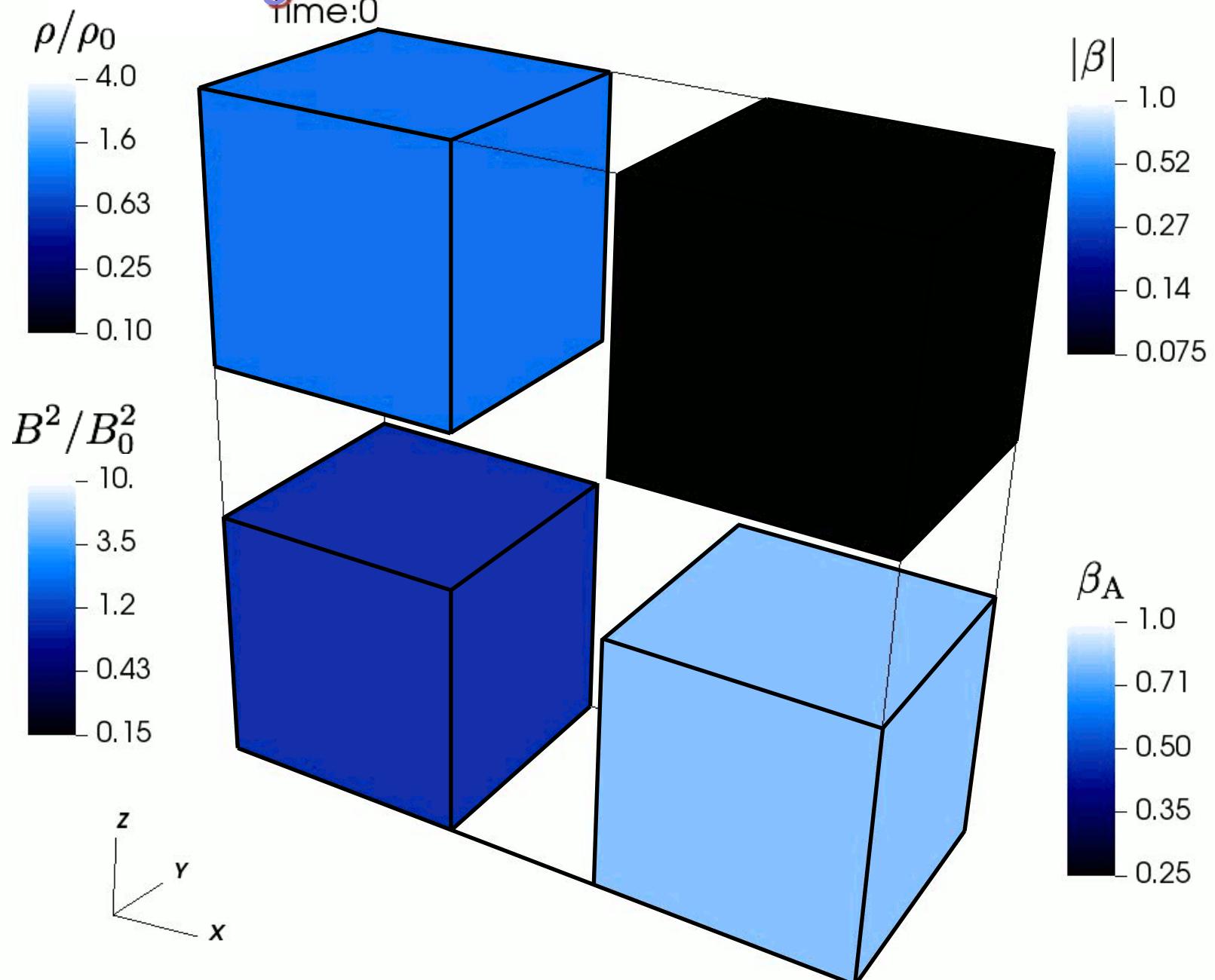
a fast rate: $t_{\text{acc}} \sim 0.1 t_g$
... leading to hard powerlaws with
 $p < 2$ if $\sigma \gtrsim 1$...
but a cut-off at Lorentz factor $\sim 4\sigma$
(Kagan+15, Werner+16)



(2) acceleration in the
resulting large-scale
turbulence...
through Fermi processes!

leading to softer powerlaws...
and slow acceleration: $t_{\text{acc}} \propto t_g^{-2}$

Relativistic magnetized turbulence



Relativistic MHD simulation (256^3 , $\sigma_0=30$, $u_A \sim 5$) by Camilia Demidem (Demidem+ 19, in

Relativistic magnetized turbulence

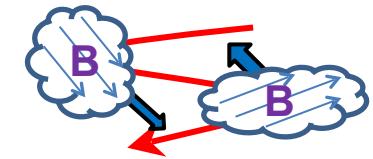
Fermi model for acceleration:

... particle interaction with random moving scattering centers...

... acceleration becomes $t_{\text{acc}} \sim \frac{L}{\beta_u^2}$... what are t_{int} ? β_u ?

coefficient: ... (quasilinear theory of) resonant wave-particle interactions in wave turbulence:

$t_{\text{int}} \sim t_{\text{scatt}}$, $\beta_u \sim \beta_A$... (Kennel + Engelmann 66... e.g. Schlickeiser 89)



... however, in (modern) anisotropic turbulence, resonances disappear...

(Chandran 00, Yan+Lazarian 02, Lynn+ 14, Demidem+ 19)

$$t_{\text{acc}} \sim \frac{L_{\text{max}}}{\beta_A^2}$$

... for particle interacting with structures (~eddies) rather than waves:

(Lemoine 19, Demidem+ 19)

$$t_{\text{acc}} \sim \frac{L_{\text{max}}}{\langle \delta u^2 \rangle} \quad \Rightarrow \text{slow at (low) energies } r_g \ll L_{\text{max}}, \text{ fast at high}$$

... in recent PIC simulations: energy...

(Zhdankin+17, Comisso+ Sironi 18,19, Wong+ 19)

$$t_{\text{acc}} \sim \frac{L_{\text{max}}}{\langle \delta u^2 \rangle}$$

not efficient for leptons, but acceleration to close to confinement possible

Shear acceleration

Fermi shear acceleration:

... the electric field in a sheared velocity flow cannot be boosted away globally: particles gain energy by exploring the shear gradient...

(e.g. Rieger+Duffy 04, 06, 08, Liu+ 17, Rieger 19, Webb+ 18,19, Lemoine 19)

... acceleration timescale:

$$t_{\text{acc}} \sim \frac{\Delta r^2}{t_{\text{scatt}}} \frac{1}{\Delta u^2 / \gamma_u^2}$$

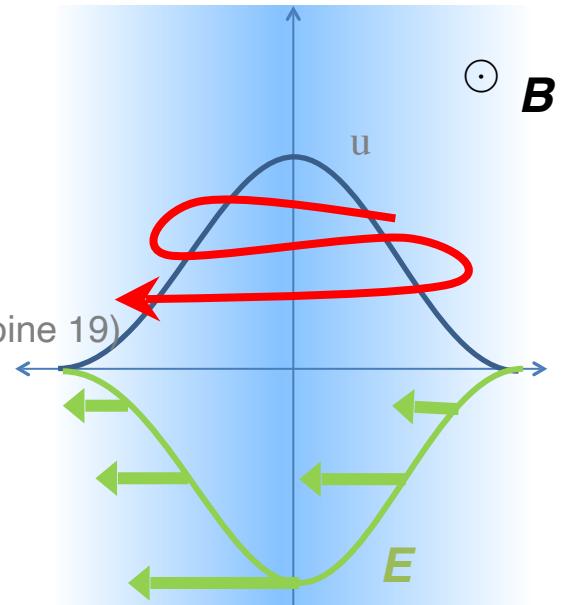
⇒ inefficient at low energies, since $t_{\text{scatt}} \nearrow p$,
requires a seed population of particles

... $t_{\text{acc}} \lesssim t_{\text{esc}} \sim \Delta r^2 / t_{\text{scatt}}$ if $\Delta u \gtrsim 1$: optimal for (mildly?) relativistic shear!

... if $L_{\text{max}} \sim \Delta r$, at confinement energy $r_g \sim \Delta r$, $\Rightarrow t_{\text{scatt}} \sim r_g \sim \Delta r \Rightarrow t_{\text{acc}} \sim r_g$

for $\Delta u \sim u \sim 1$

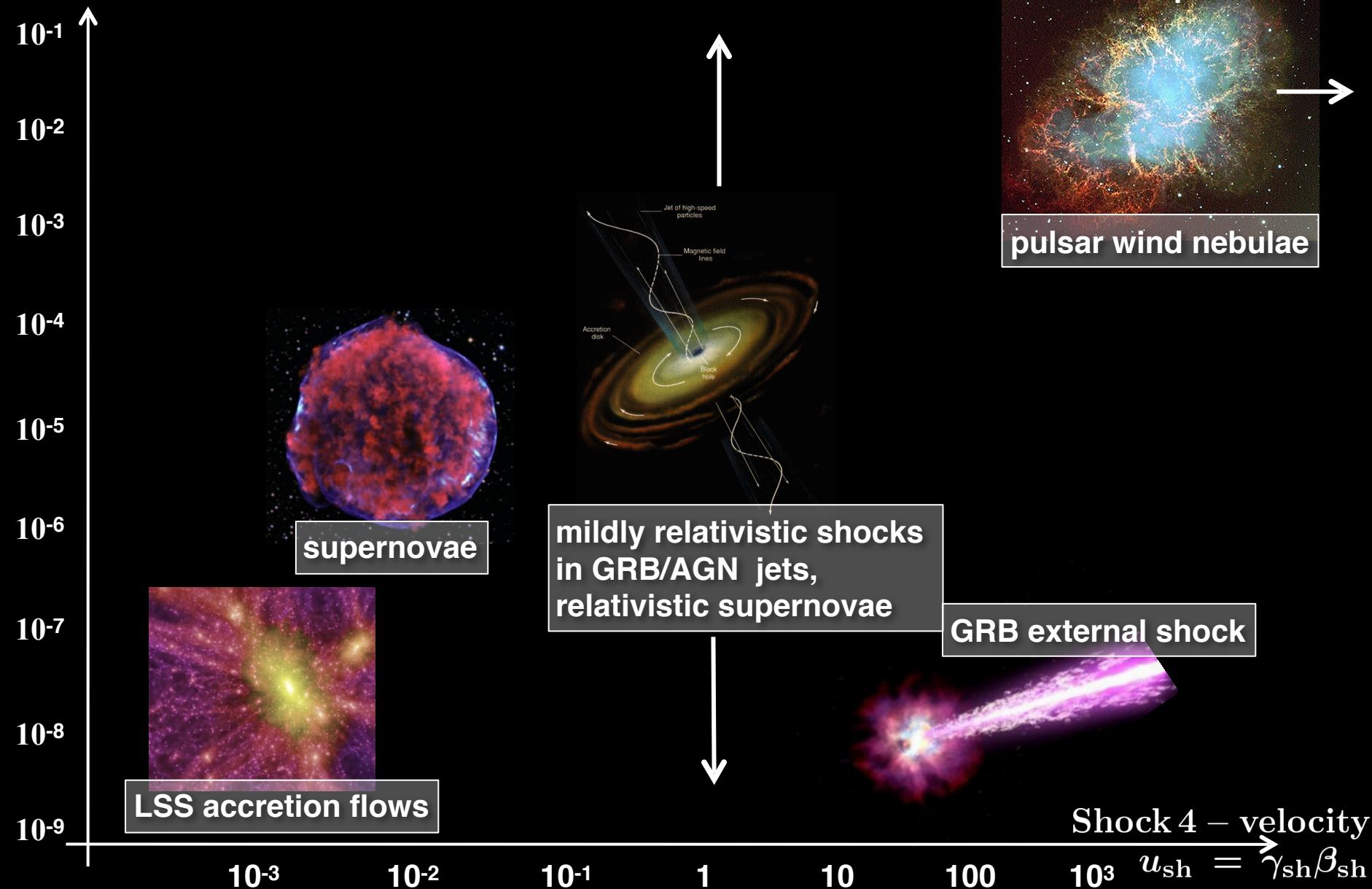
⇒ reacceleration of a population of energetic CRs in mildly relativistic shear may reach confinement energy... (radiative signatures? beyond test-particle?)



particles with larger mean free paths explore larger gradient of $E \Rightarrow$ faster acceleration...

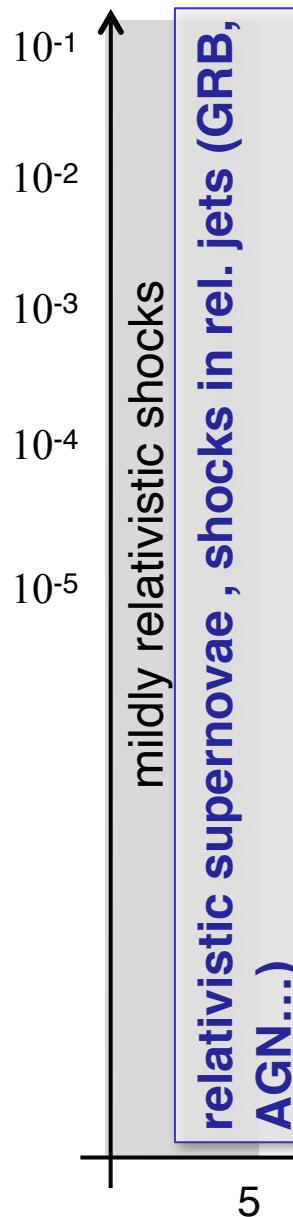
The (HE) astrophysical shock landscape

Magnetization σ $\sigma = u_A^2/v_{sh}^2$



Particle acceleration at relativistic shock waves

$$\sigma = (u_A/c)^2$$



Pulsar
Wind
Nebulae

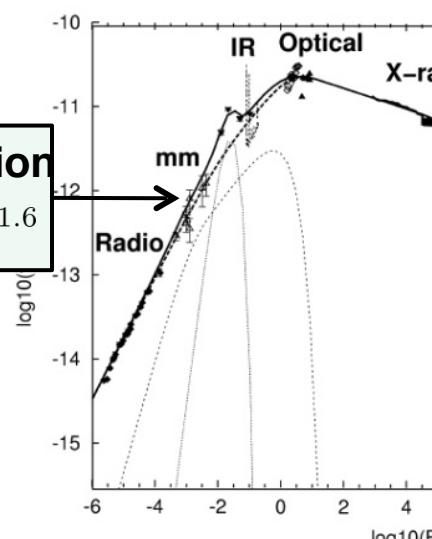
The issue of pulsar wind nebulae...

$$\sigma = (u_{\perp}/c)^2$$

→ **theory may not be complete:** predicts no acceleration at pulsar wind termination shock, while SED suggests **Fermi-type acceleration at Bohm regime:**
(Atoyan & Aharonian 96)

Pulsar
Wind
Nebulae

pair distribution
 $dN_e/dp_e \propto p_e^{-1.6}$



pair distribution

$$dN_e/dp_e \propto p_e^{-2.2}$$

synchrotron burn-off limit:

$$\epsilon_{\text{syn,max}} \sim 100 (t_{\text{acc}}/t_g)^{-1} \text{ MeV}$$
$$\Rightarrow t_{\text{acc}} \sim t_g$$

→ Nature realizes optimal acceleration... how?

Gamma-ray burst afterglows

5

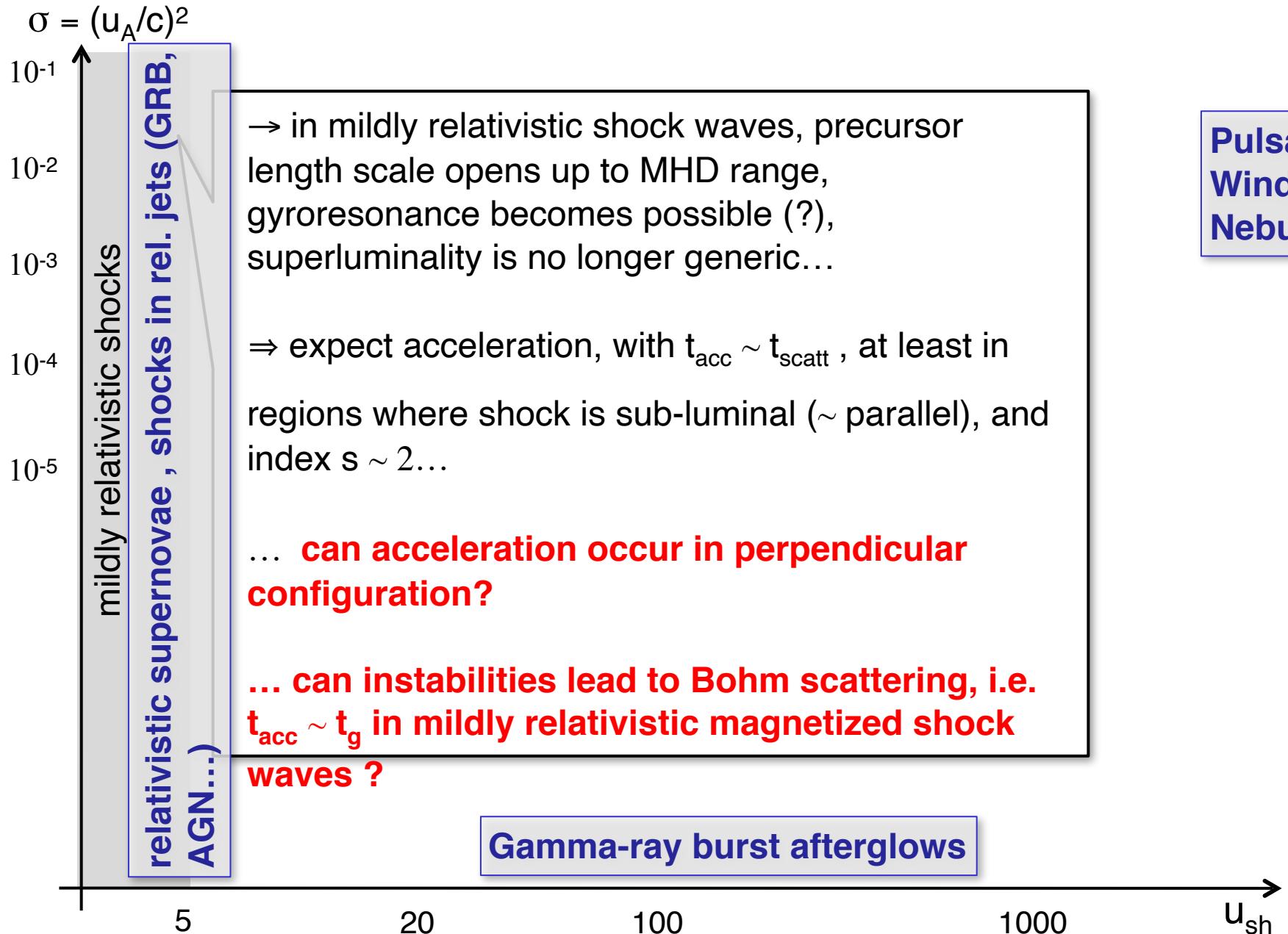
20

100

1000

u_{sh}

Mildly relativistic shocks... poorly explored but promising



Pulsar
Wind
Nebulae

Summary...

→ **a challenge of scales:** a complex problem on microscopic scales, needing extrapolation to macroscopic scales, in the midst of many astrophysical unknowns...

→ **Top-down, insights from phenomenology:**

$$L_{\text{tot}} \gtrsim 10^{45} \text{ erg/s} \dots \left(\frac{t_{\text{acc}}}{t_g} \right)^2 \left(\frac{E/Z}{10^{20} \text{ eV}} \right)^2$$

... powerful sources, likely relativistic, w/ strong dependence on Z...

→ **Bottom-up, from microphysics of acceleration upward:**

... acceleration to UHE favors Fermi-type processes in mildly relativistic flows of

substantial magnetization, e.g. shocks, turbulence or shear acceleration...

→ **an interesting question:**

... are UHECRs accelerated all the way up in one source, or accelerated in one,

then re-accelerated in others?

one source; the perfect accelerator / many sources; radiative

Particle acceleration at relativistic shock waves

$$\sigma = (u_A/c)^2$$

