

STUDY OF DE-CONFINEMENT PHASE TRANSITION IN THE COLOR STRING PERCOLATION APPROACH IN ALICE AT LHC

Aditya Nath Mishra, Eleazar Cuautle and Guy Paic
Instituto de Ciencias Nucleares UNAM, CDMX, Mexico

Aditya Nath Mishra, ICN, UNAM

Introduction: Color Strings

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- Multiparticle production at high energies is currently described in terms of color strings stretched between the projectile and target.
- These strings decay into new ones by production and subsequently hadronize to produce the observed hadrons.
- The number of strings and the string density grows with energy and with the number of participating nucleons.
- As the no. of strings grow, they start to overlap and form clusters.
- At a critical density a macroscopic cluster appears and marks the percolation phase transition.
- Particles are produced by the Schwinger mechanisms.

Introduction: Color Strings

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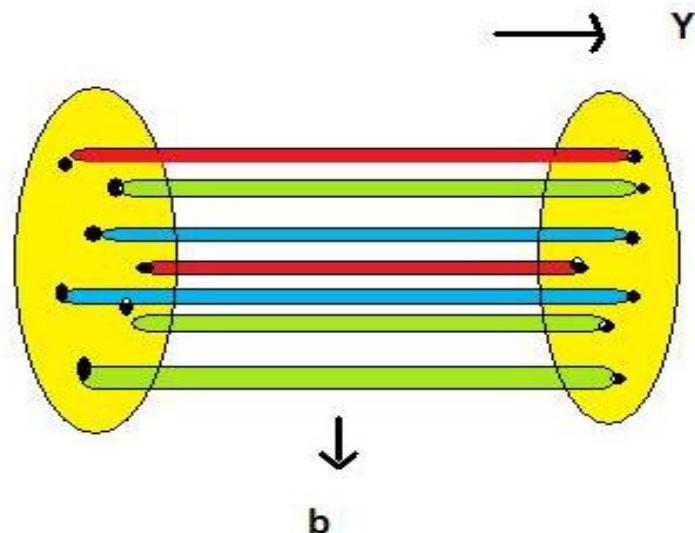
- In the case of a nuclear collisions, the density of disks - elementary strings:

$$\xi = \frac{N^s S_1}{S_N}$$

N^s = Number of strings

S_1 = Single string area

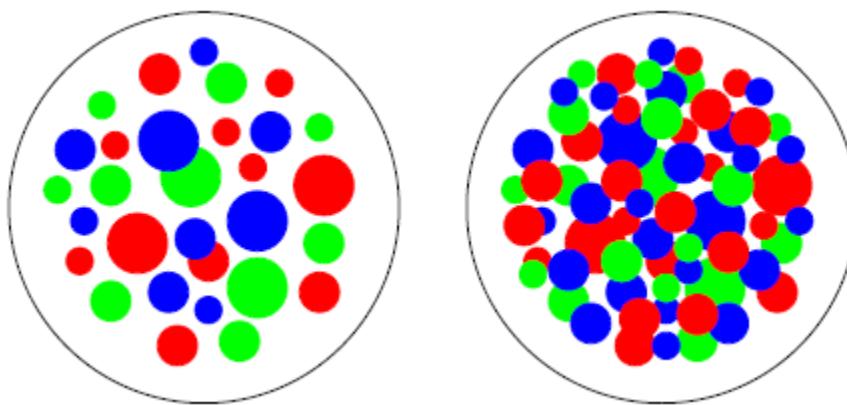
S_N = total nuclear overlap area



Introduction: Clustering of Color Sources

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- De-confinement is expected when the density of quarks and gluons becomes so high that it no longer makes sense to partition them into color-neutral hadrons, since these would overlap strongly.
- We have clusters within which color is not confined : De-confinement is thus related to cluster formation very much similar to cluster formation in percolation theory and hence a connection between percolation and de-confinement seems very likely.



Parton distributions in the transverse plane of nucleus-nucleus collisions

In two dimensions, for uniform string density, the percolation threshold for overlapping discs is:

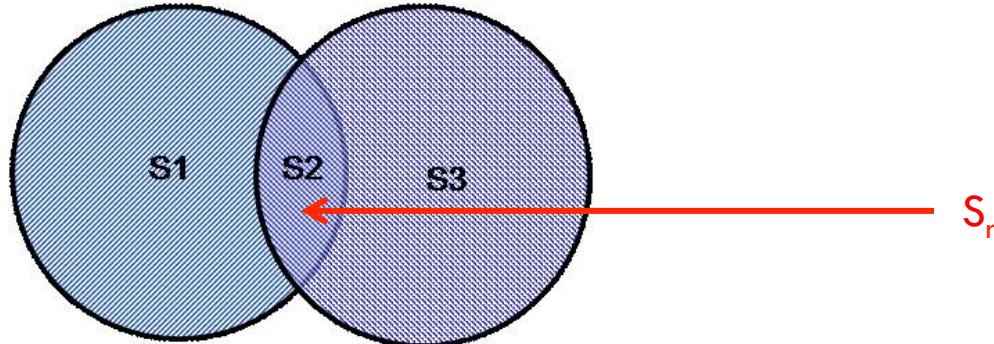
$$\xi_c = 1.18$$

Critical Percolation Density

H. Satz, Rep. Prog. Phys. 63, 1511(2000).
H. Satz, hep-ph/0212046

Introduction: Clustering of Color Sources

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- The transverse space occupied by a cluster of overlapping strings split into a number of areas in which different number of strings overlap, including areas where no overlapping takes place.
- A cluster of n strings that occupies an area S_n behaves as a single color source with a higher color field \vec{Q}_n corresponding to vectorial sum of color charges of each individual string \vec{Q}_1

$$\vec{Q}_n^2 = n \vec{Q}_1^2 \quad \text{If strings are fully overlap}$$

$$\vec{Q}_n^2 = n \frac{S_n}{S_1} \vec{Q}_1^2 \quad \text{Partially overlap}$$

Introduction: Clustering of Color Sources

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Multiplicity and $\langle p_T^2 \rangle$ of particles produced by a cluster of n strings

Multiplicity (μ_n)

$$\mu_n = F(\xi) N^s \mu_1$$

Average Transverse Momentum

$$\langle p_T^2 \rangle_n = \langle p_T^2 \rangle_1 / F(\xi)$$

$$F(\xi) = \sqrt{\frac{1 - e^{-\xi}}{\xi}}$$

Color suppression factor
(due to overlapping of discs).

$$\xi = \frac{N^s S_1}{S_N}$$

N^s = # of strings

S_1 = disc area

S_N = total nuclear overlap area

ξ is the string density parameter

M. A. Braun and C. Pajares, Eur. Phys. J. C16, 349 (2000)

M. A. Braun et al, Phys. Rev. C65, 024907 (2002)

Methodology

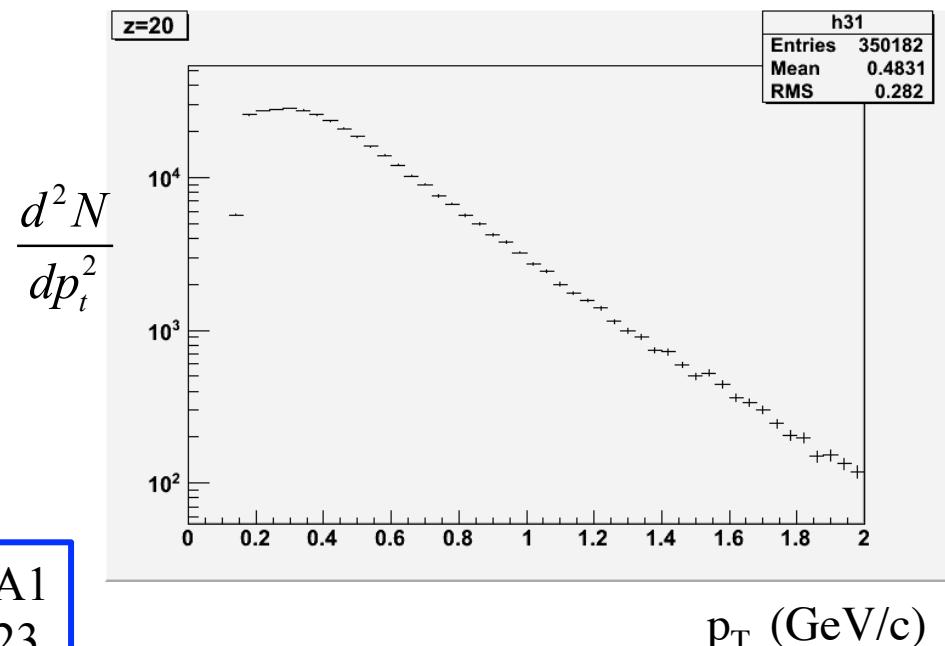
- Using the p_T spectrum to calculate $F(\xi)$ The experimental p_T distribution from pp data is used

$$\frac{d^2N}{dp_T^2} = \frac{a}{(p_0 + p_T)^n}$$

Here, a , p_0 and n are fit parameters to the proton-proton data.

For parameterizations, pp data from UA1 200, 500 and 900 GeV & ISR 53 and 23 GeV are used.

$p_0 = 1.71$ and $n = 12.42$

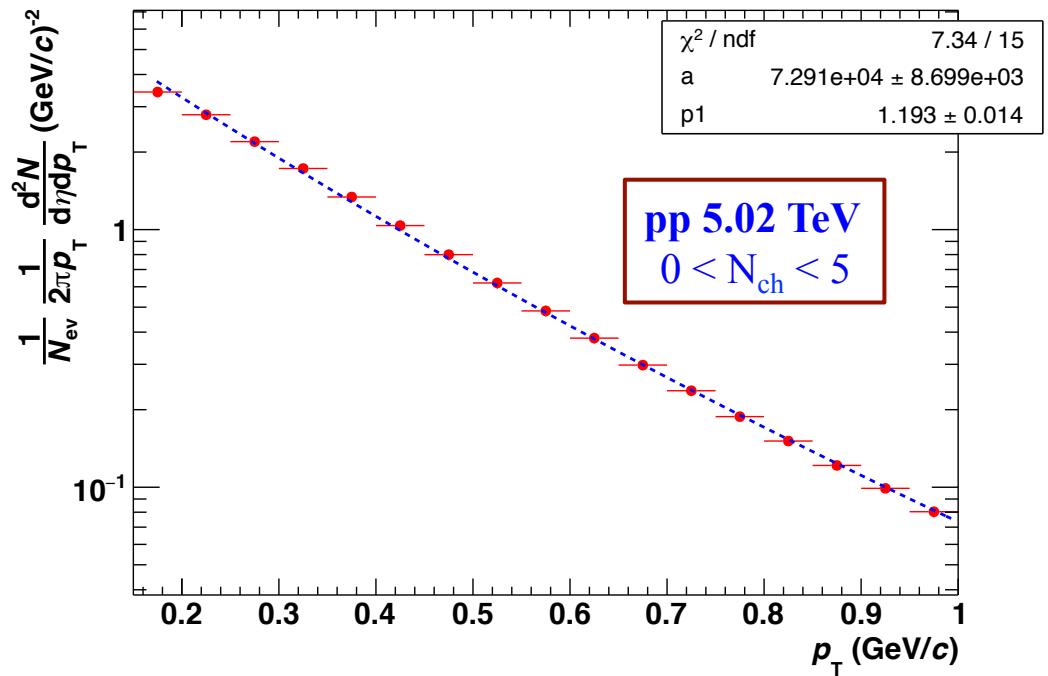


Methodology

- This parameterization can be used for high multiplicity events in pp and Heavy-ion collisions to account for the clustering :

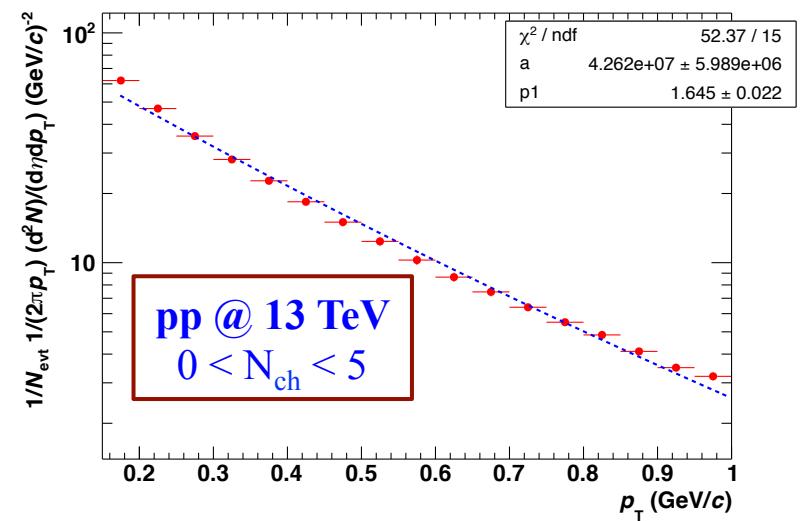
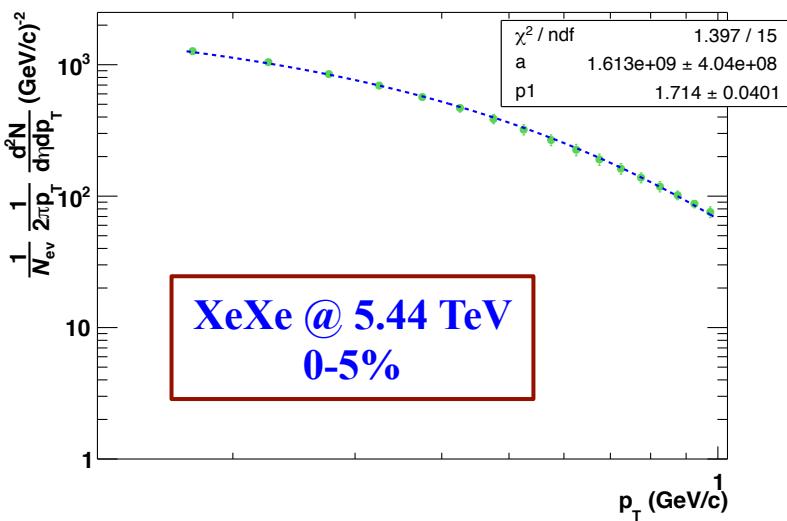
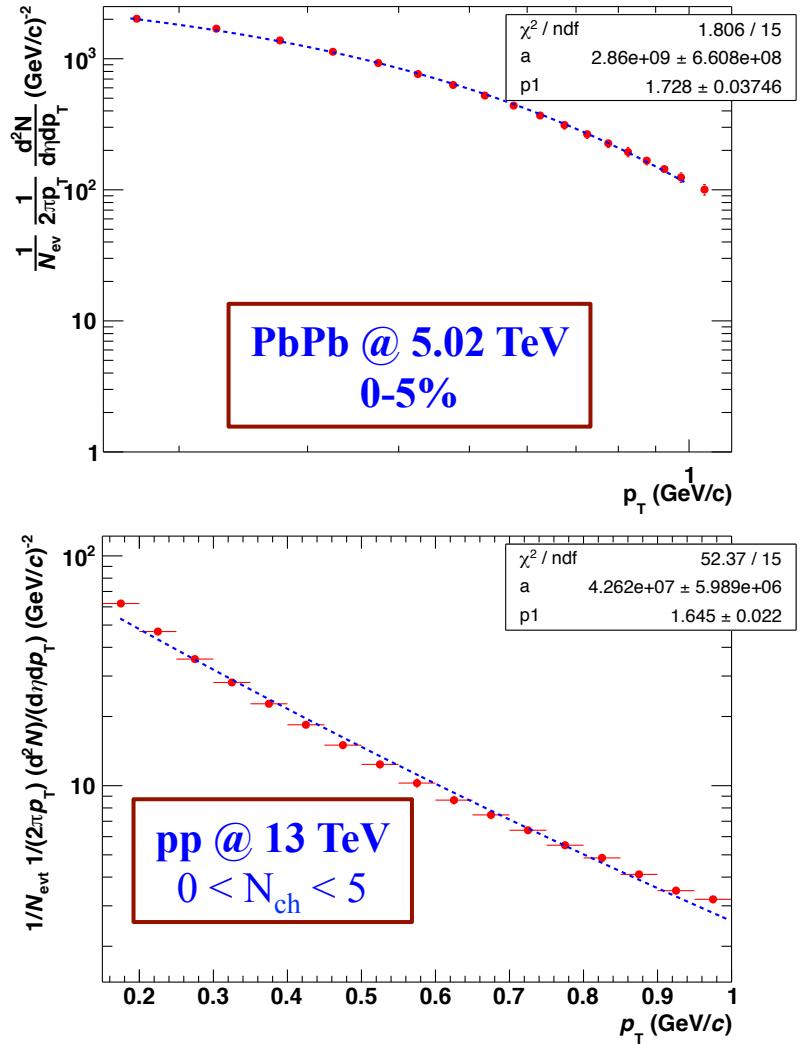
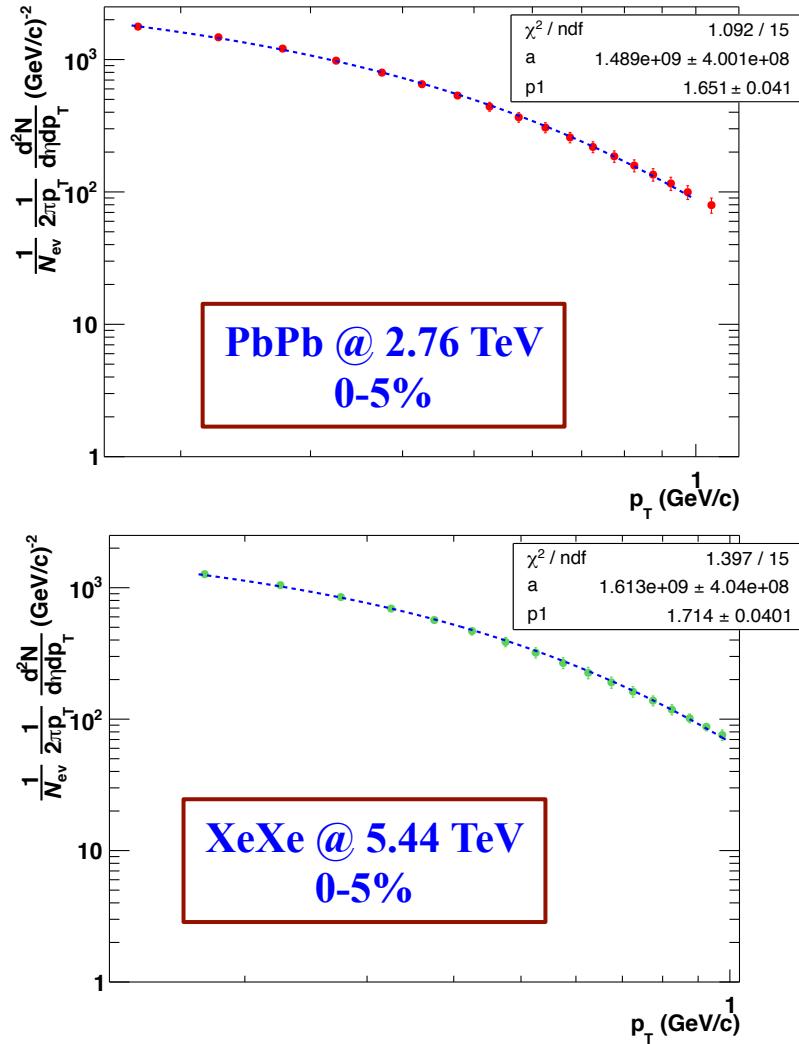
$$\frac{d^2 N}{dp_T^2} = \frac{b}{\left(p_0 \sqrt{\frac{F(\xi_{pp})}{F(\xi_{pp})_{HM}}} + p_T \right)^n}$$

$$F(\xi_{pp}) = 1$$



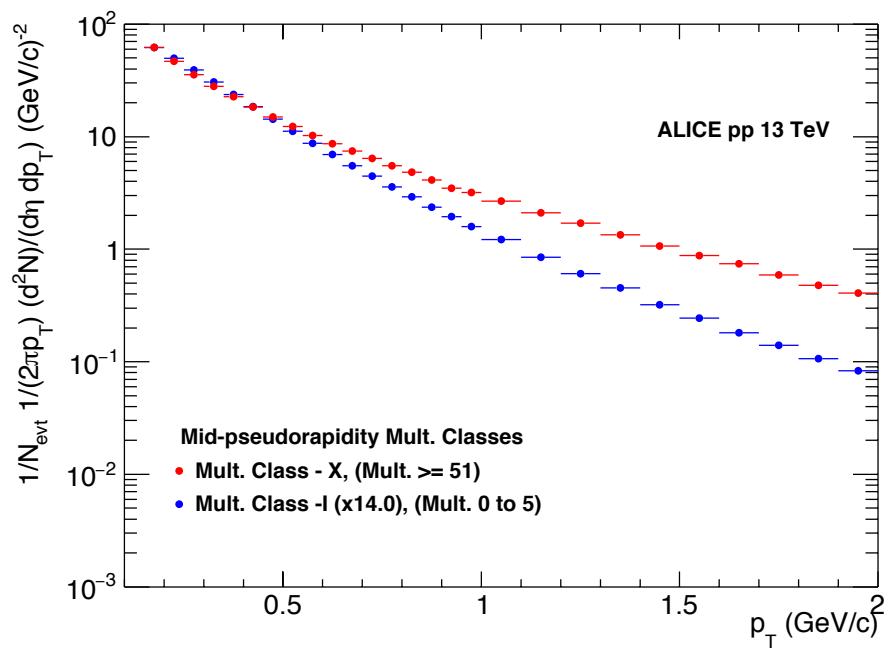
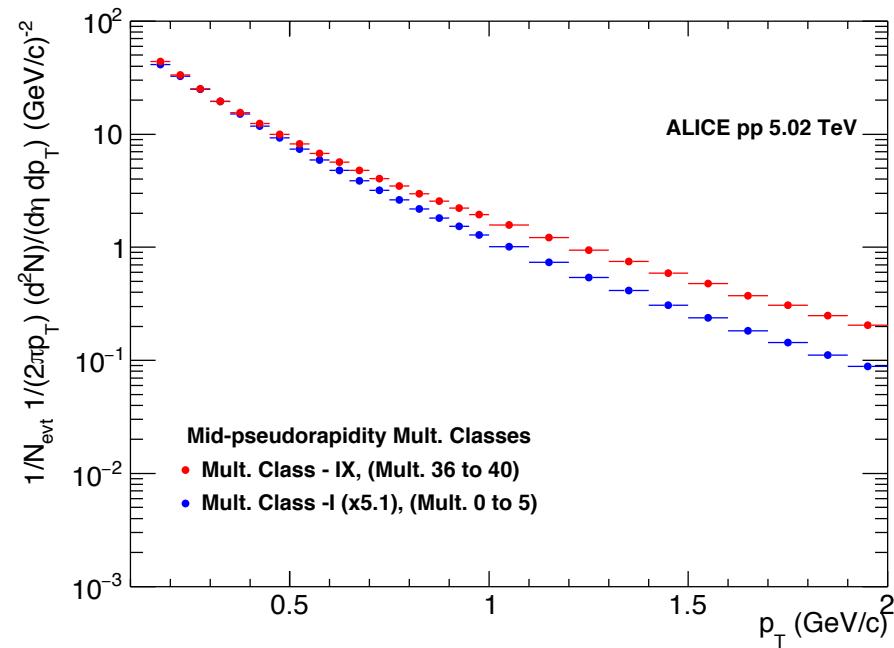
M. A. Braun and C. Pajares, Eur. Phys. J. C16,349 (2000)
M. A. Braun et al, Phys. Rev. C65, 024907 (2002)

Methodology



Analysis: p_T Spectra

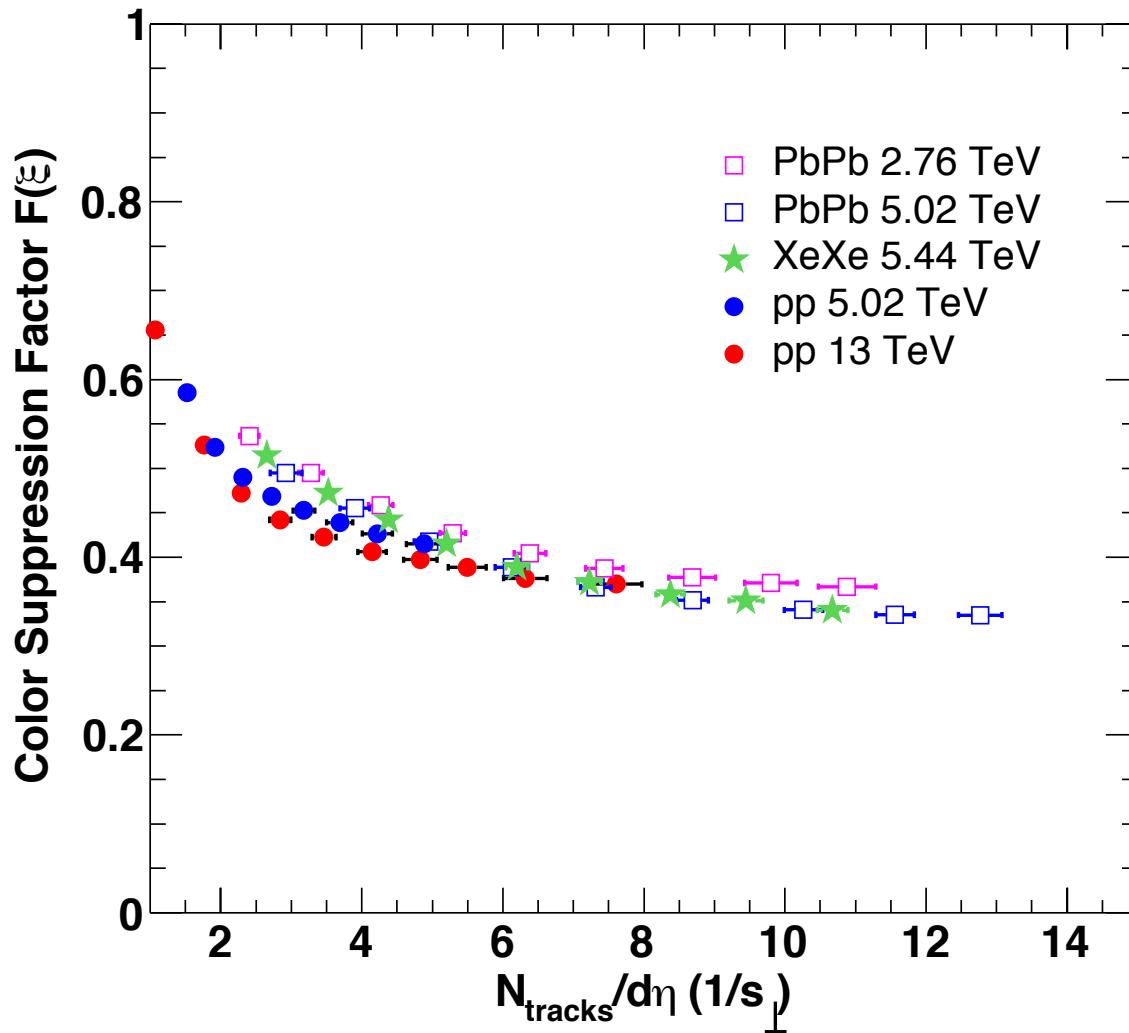
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- p_T spectra for two different multiplicity classes to show that for higher multiplicity p_T spectra gets harder which results a lower value of **Color Supuration Factor “ $F(\xi)$ ”** and higher values of **Temperature**.

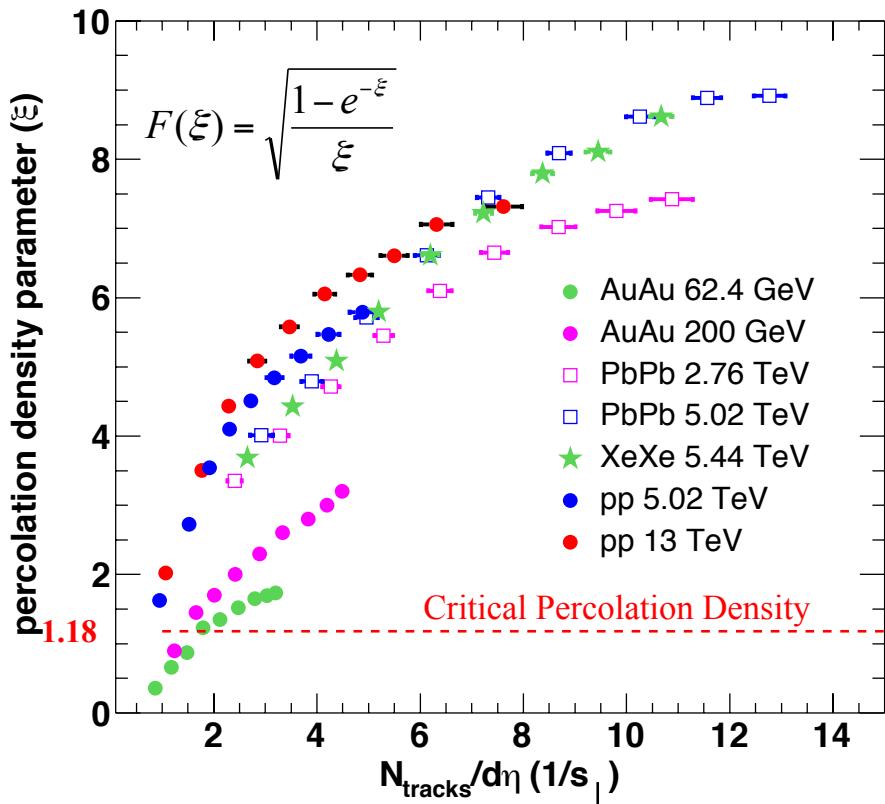
Analysis: Color Suppression Factor

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Results: Percolation Density Parameter

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➤ **ξ value for AuAu 62.4 GeV collisions:**
Almost all centralities, except for the three most peripheral bins, lie above the critical percolation density.

➤ **ξ value for AuAu 200 GeV collisions:**
All centralities, except for the most peripheral bin, lie above the critical percolation threshold.

➤ For all the ALICE data shown here,
 ξ values are well above than the critical percolation threshold.

Results: Cluster/Initial Temperature

$$T = \sqrt{\frac{\langle p_t^2 \rangle_1}{2F(\xi)}}$$

At the critical percolation density
 $\xi = 1.2 \rightarrow T_c = 167$

PHENIX:

For Au+Au @200 GeV
 0-10% centrality

$\xi = 2.88 \rightarrow T \sim 195 \text{ MeV}$

Direct Photon Measurement

$T = 220 \pm 19^{(\text{stat})} \pm 19^{(\text{syst})} \text{ MeV}$
Phys. Rev. Lett. 104, 132301 (2010)

ALICE:

For Pb+Pb @ 2.76 TeV
 0-5% centrality

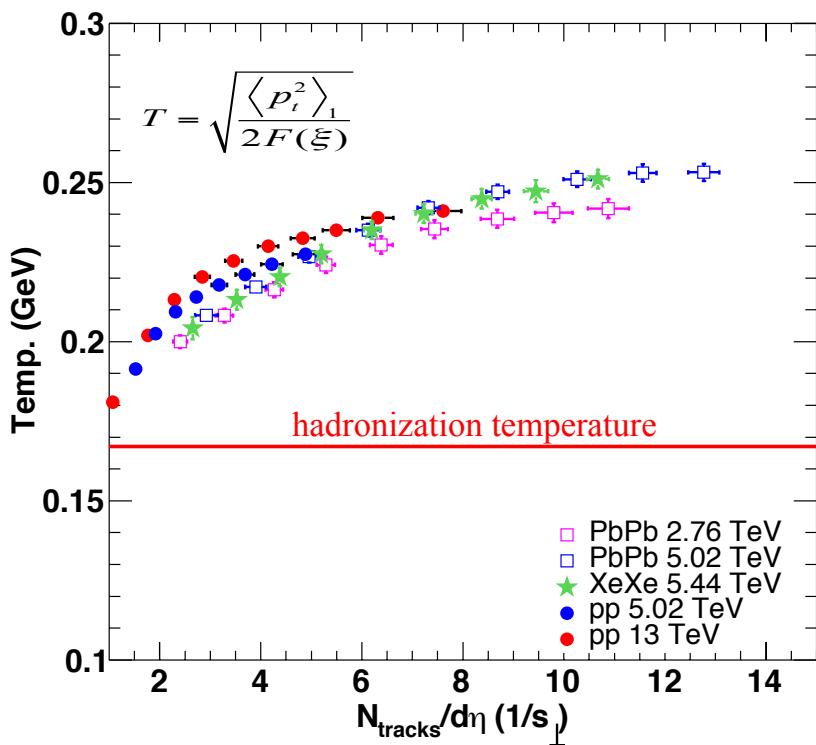
$\xi = 10.56 \rightarrow T \sim 262 \pm 13 \text{ MeV}$

Direct Photon Measurement

$T = 297 \pm 12^{(\text{stat})} \pm 41^{(\text{syst})} \text{ MeV}$
Phys. Lett. B 754 (2016) 235-248

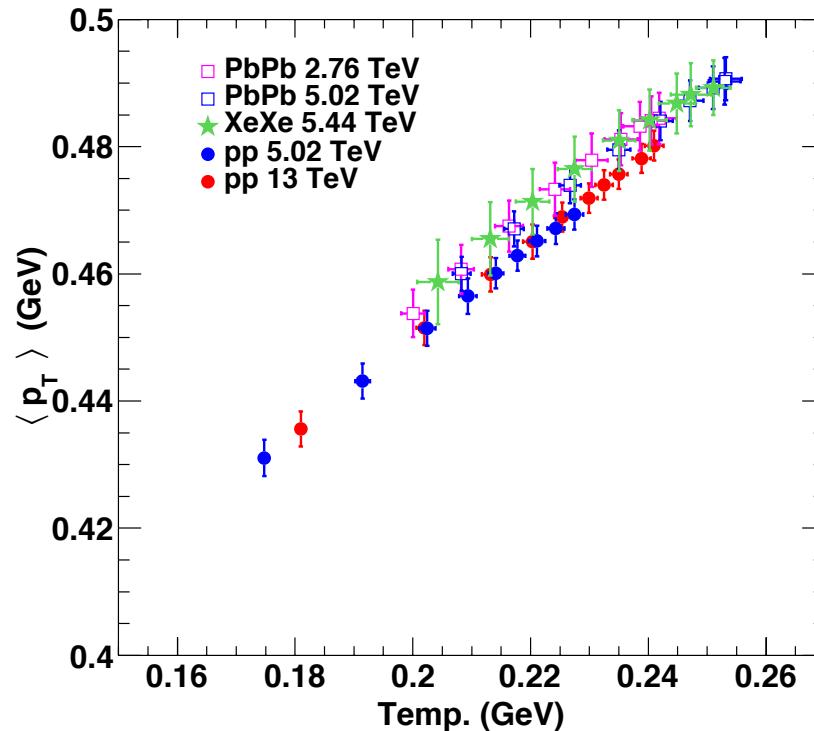
Results: Cluster/Initial Temperature

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- Temperature from hadron-hadron and nucleus-nucleus collisions fall doesn't show a clear scaling multiplicity is scaled by the transverse interaction area.
- All the temperatures are above the universal hadronization temperature (165 MeV)
- The obtained temperatures indicate the creation of de-confined matter in pp collisions at $\sqrt{s} = 5.02$ and 13 TeV even at very low multiplicity.

Results: Mean Transverse Momentum



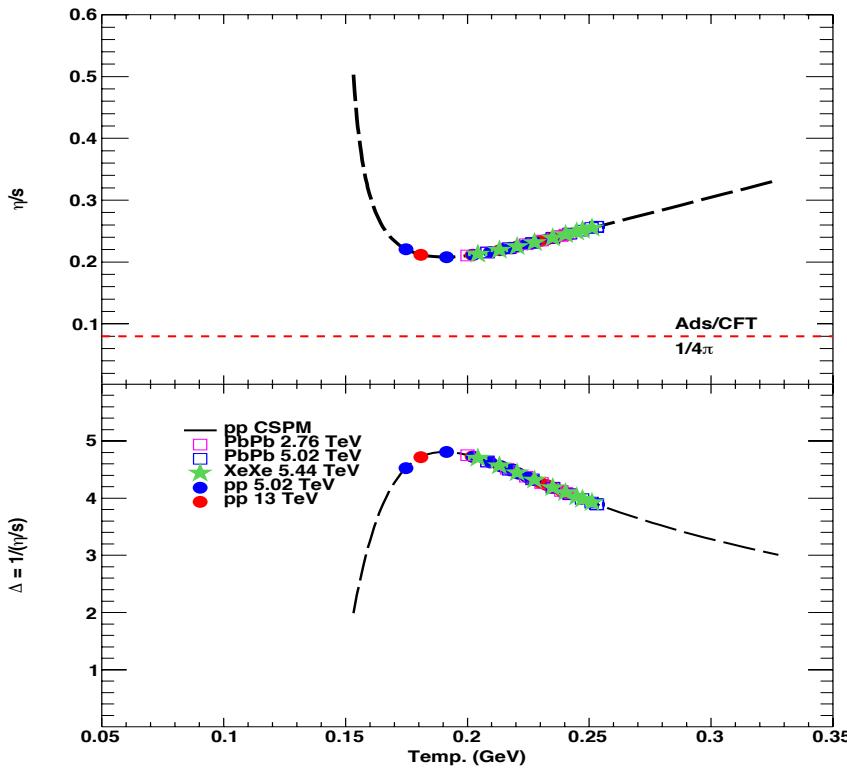
For an ultrarelativistic ideal gas $\langle p_T \rangle \propto T$. When the transverse momentum is exponentially distributed with inverse slope T in a given event, $\langle p_T \rangle = 2T^{**}$. Using this relation one can predict the $\langle p_T \rangle$ for upcoming collision energies.

**H. Heiselberg *Phys. Rep.* 351 (2001) 161.

Results: eta/s

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- η/s was obtained in the framework of kinetic theory and the string percolation with the following expression:



$$\frac{\eta}{s} = \frac{TL}{5(1 - e^{-\xi})}$$

where T is the temperature and L is the longitudinal extension of the source ~ 1 fm.

- The inverse of η/s also measures how strong are the interactions in the medium.
- Both Δ and η/s describe the transition from a strongly coupled QGP to a weakly coupled QGP.

Results: Energy Density

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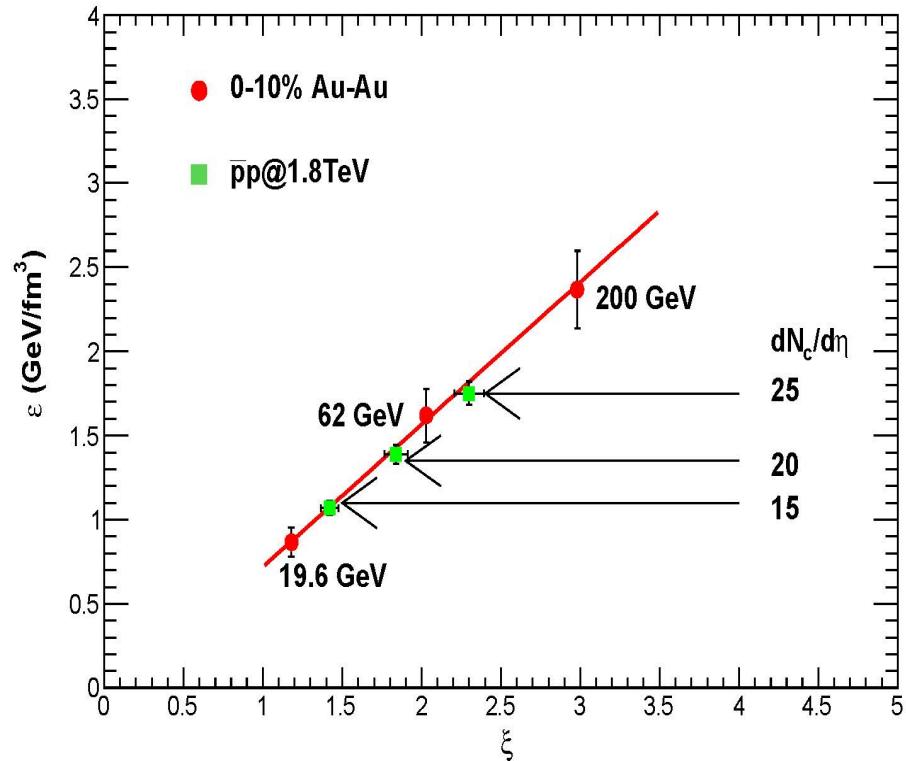
Bjorken 1D expansion

$$\varepsilon = \frac{3}{2} \frac{dN_c}{dy} \frac{\langle m_t \rangle}{A} \frac{1}{\tau_{pro}} \text{GeV / fm}^3$$

Transverse overlap area Proper Time

τ_{pro} is the QED production time for a boson which can be scaled from QED to QCD and is given by

$$\tau_{pro} = \frac{2.405\hbar}{\langle m_t \rangle}$$



Results: Energy Density

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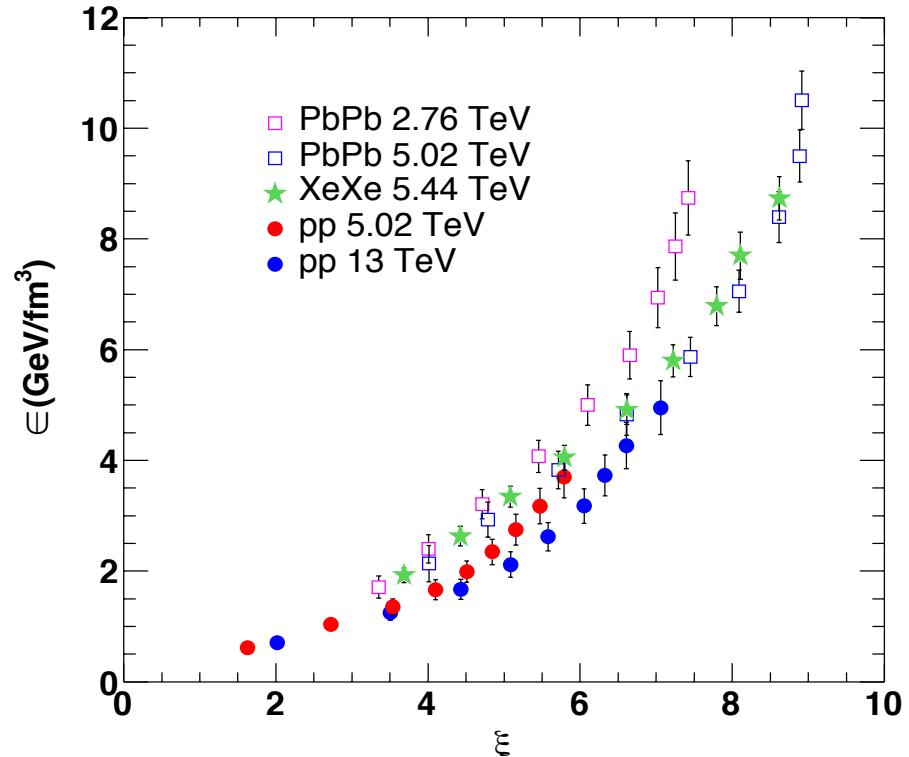
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Transverse overlap area Proper Time

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We find the non-linear relationship between Energy density and percolation density parameter.

Summary

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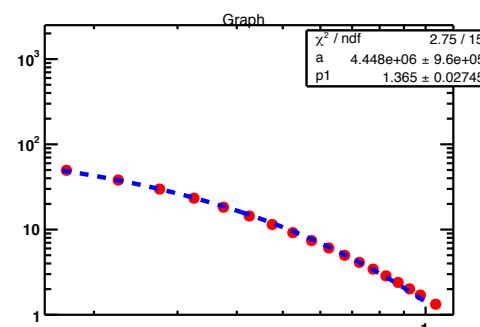
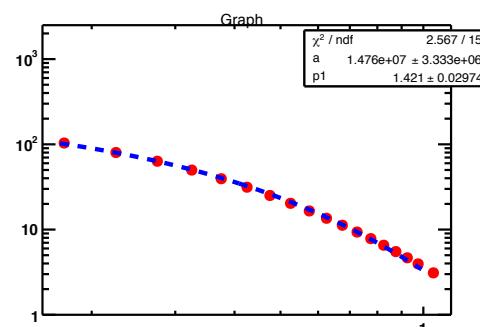
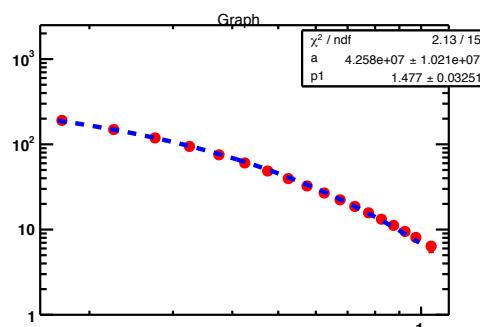
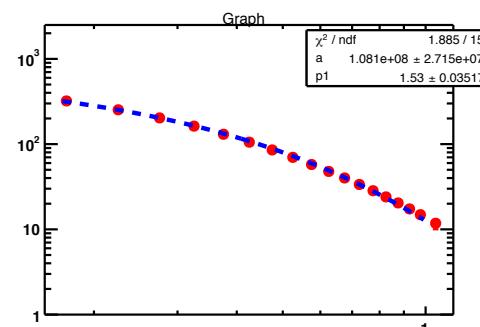
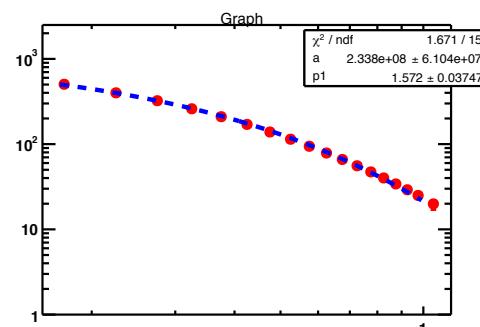
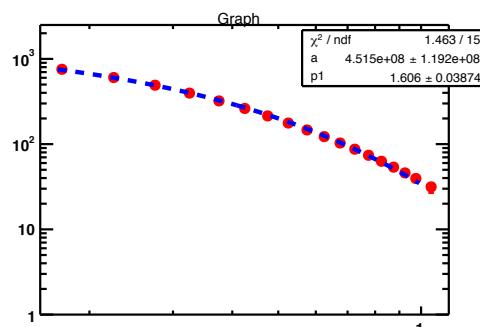
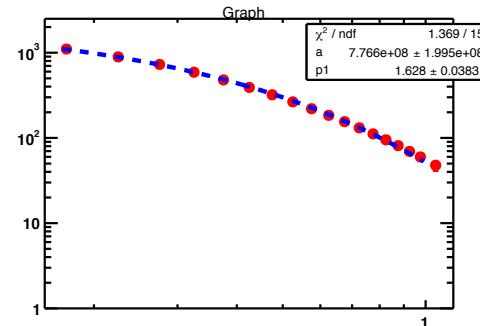
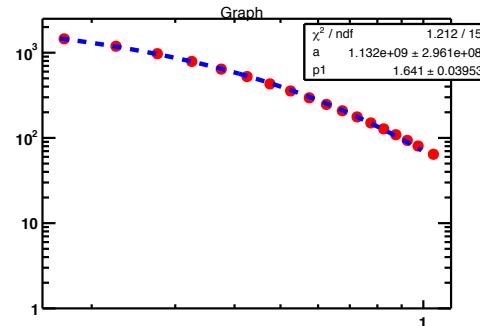
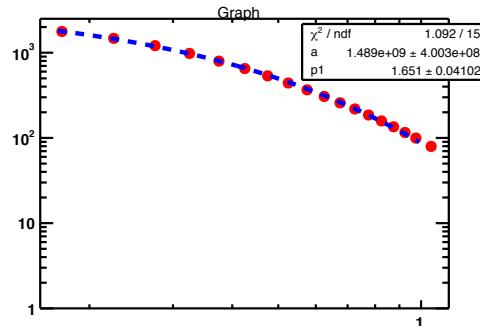
- Color string percolation concept has been explored to study the de-confinement in nuclear collisions
- Energy density obtained at LHC energies show a different behaviour than at RHIC energy.
- We are to understand the non-linear behaviour of the energy density as a function of percolation density parameter.
- Initial temperature obtained using percolation for all considered LHC energies are above the universal hadronization temperature (165 MeV) .

Thank
you

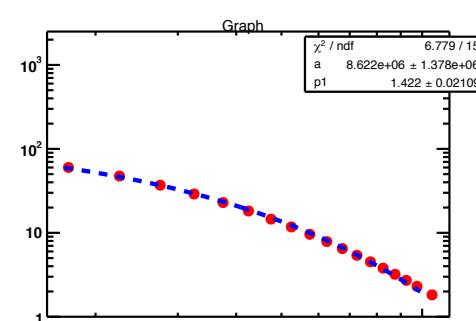
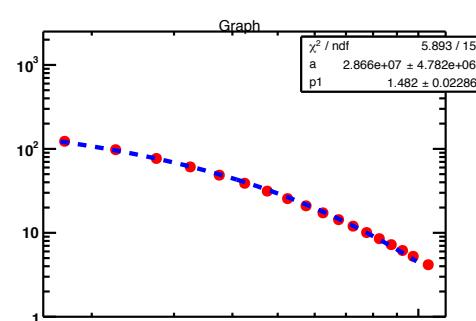
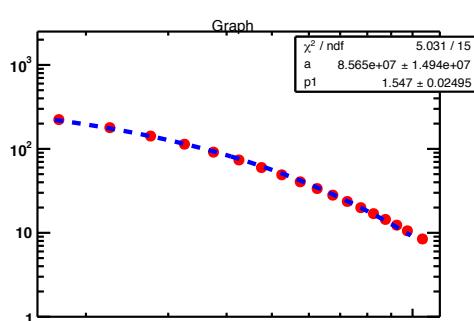
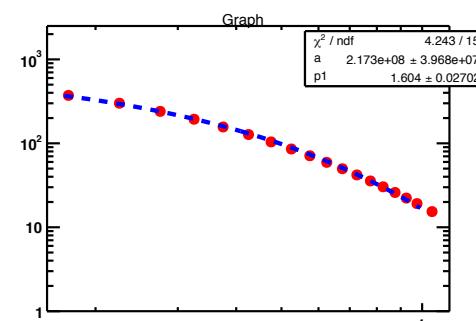
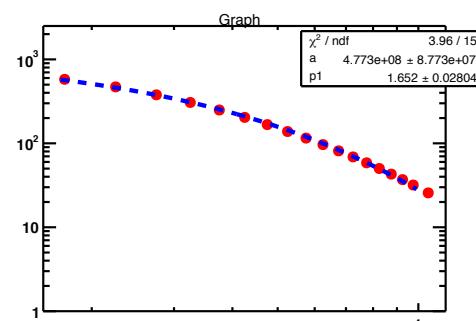
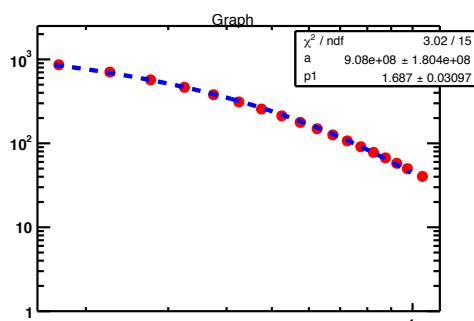
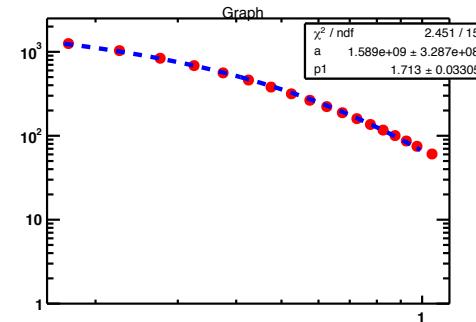
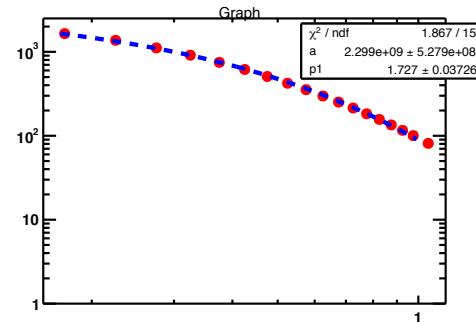
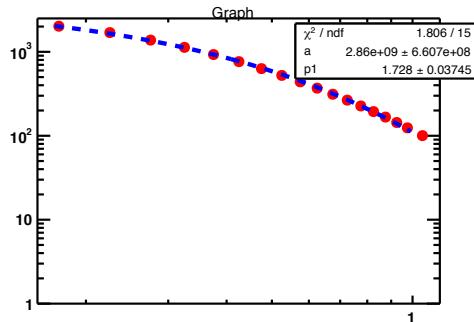


Backup Slides

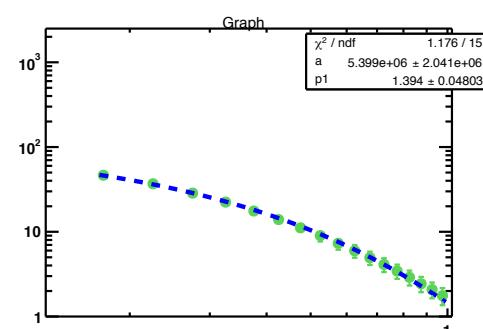
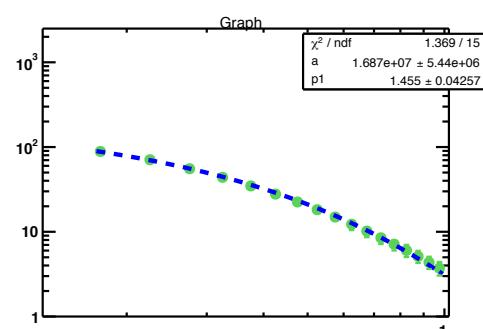
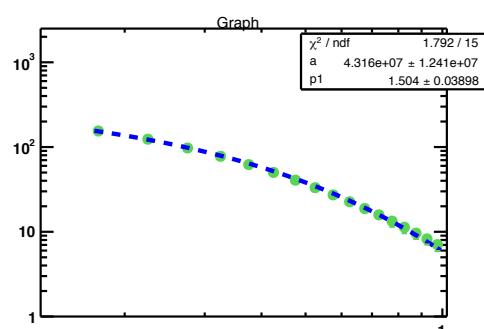
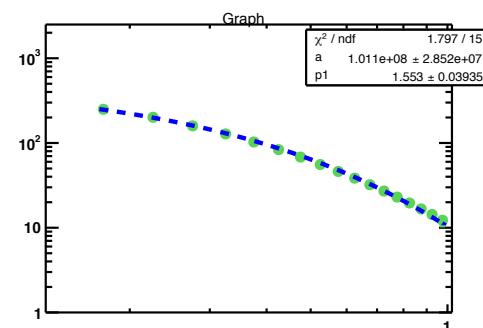
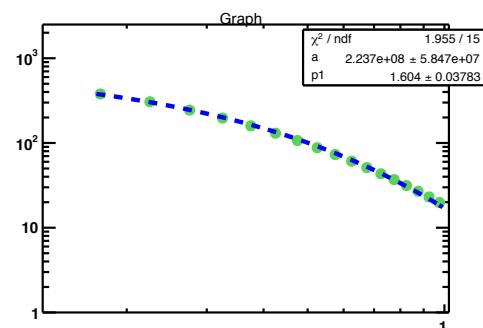
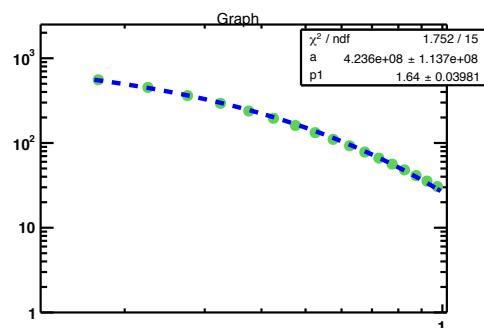
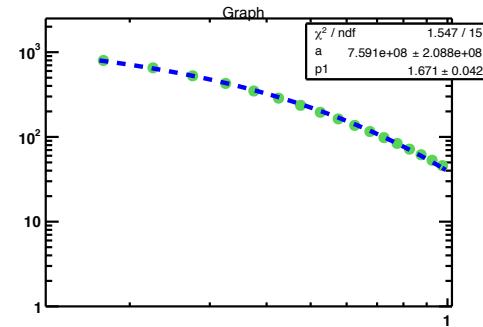
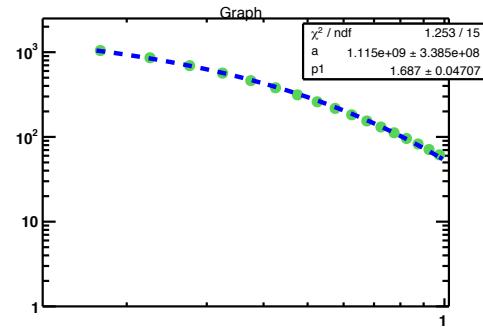
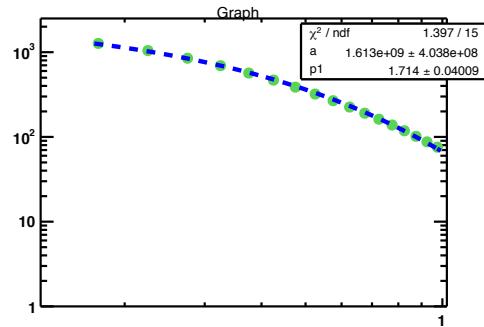
PbPb @ 2.76 TeV



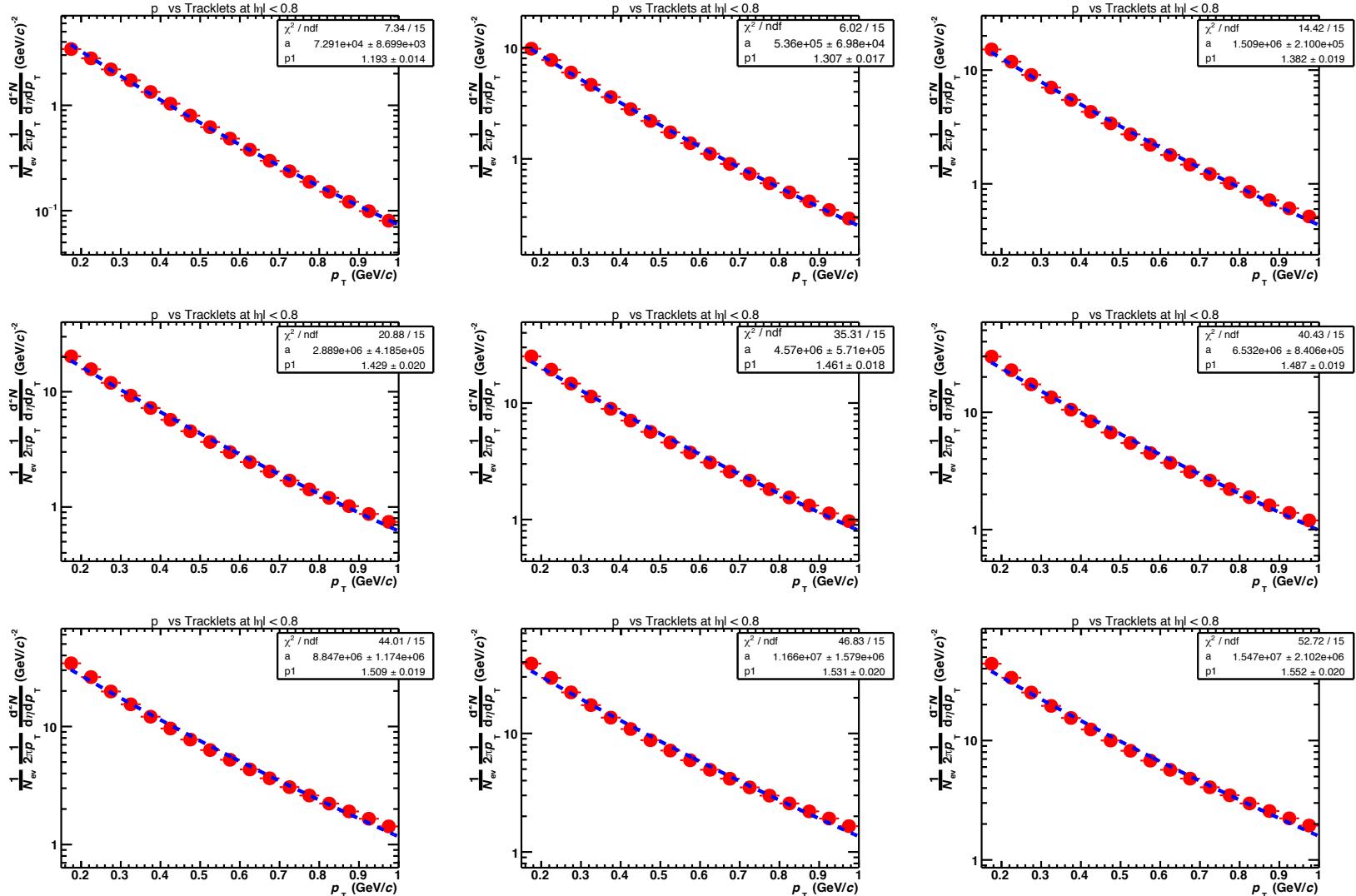
PbPb @ 5.02 TeV



XeXe @ 5.44 TeV



pp @ 5.02 TeV



pp @ 13 TeV

