



Universidad Nacional
Autónoma de México



LAr Energy Response to Nuclear Recoils

XXXII Reunión Anual de la División de Partículas y
Campos de la SMF

A. Míchell Martínez Mendoza

Universidad Nacional Autónoma de México
Instituto de Física & Facultad de Ciencias

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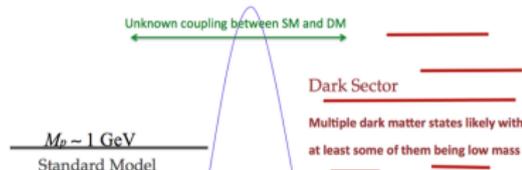
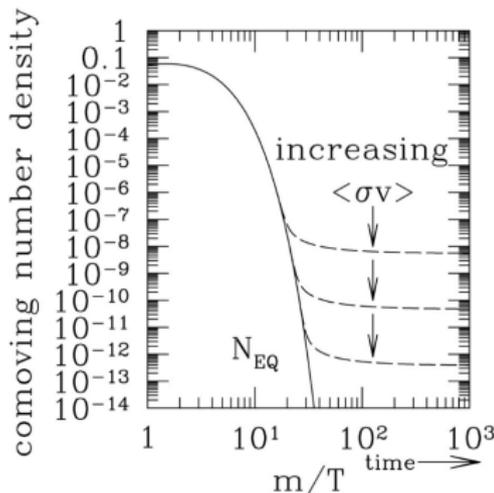
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- Most discussed candidate: **Weakly interacting Massive Particle**
- Produced during Big Bang, in thermal equilibrium in the early Universe
- Decouples from ordinary matter as the Universe expands and cools
- Still around today with densities about a few per liter

Dark sector could be as complicated as the SM. Searches not limited by expectations from SUSY models



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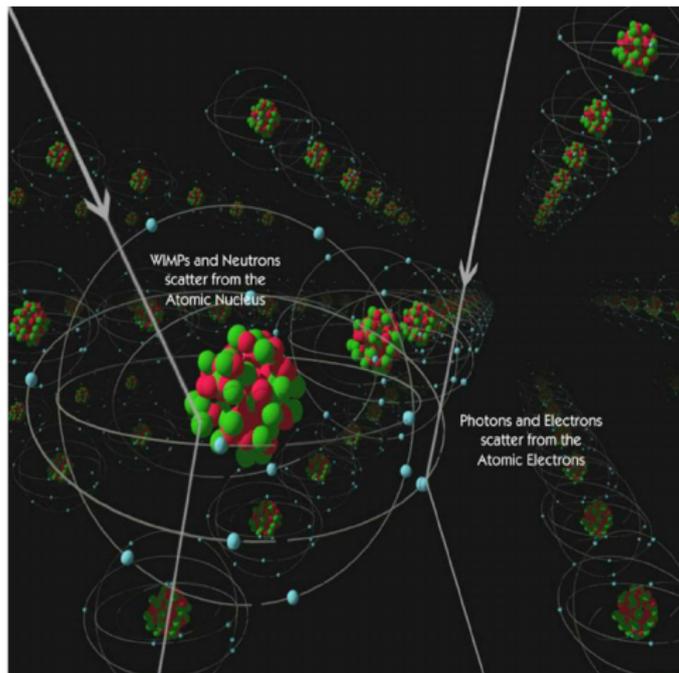
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WIMPs can scatter elastically with nuclei and the recoil can be detected.

Where are the WIMPs?

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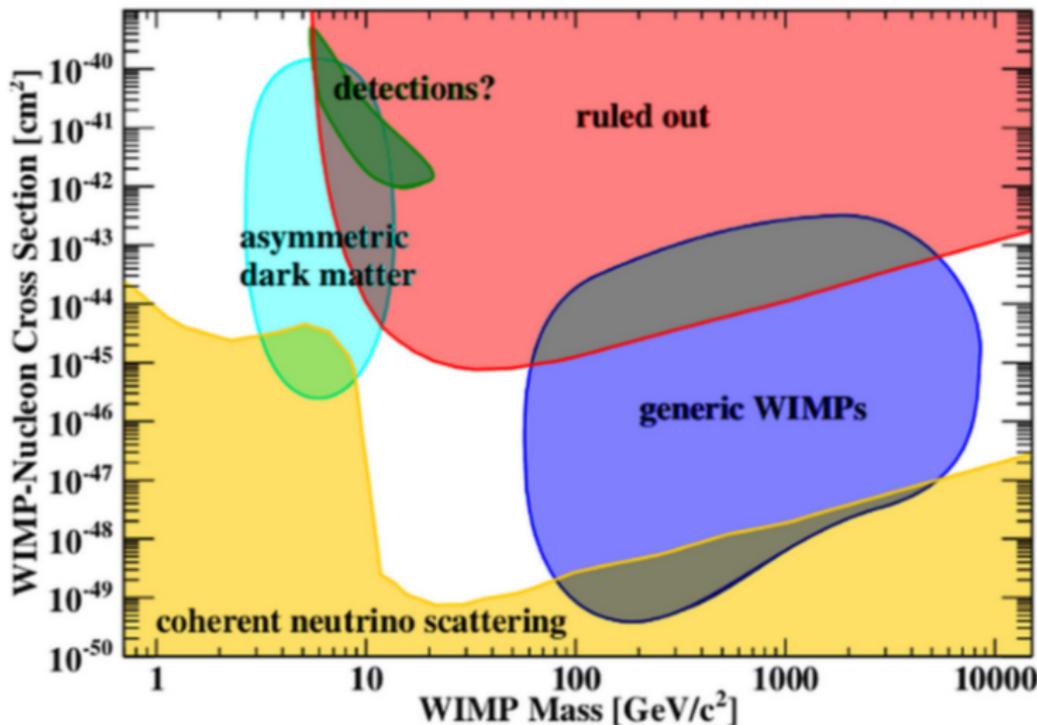
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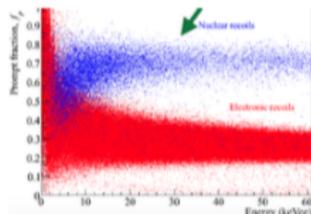
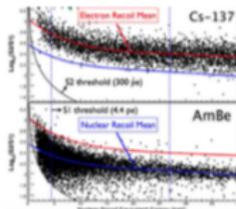
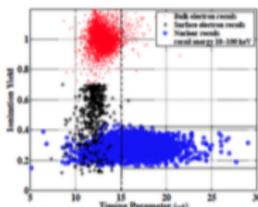
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- Detect tiny energy deposits, energy of recoils is tens of keV
- Background suppression:
 - Deep sites to reduce cosmic ray flux
 - Passive/active shielding
 - Careful choice and preparation of material
- Background discrimination (electronics recoils vs nuclear recoils)



- Large target mass, scalability to ton-scale targets

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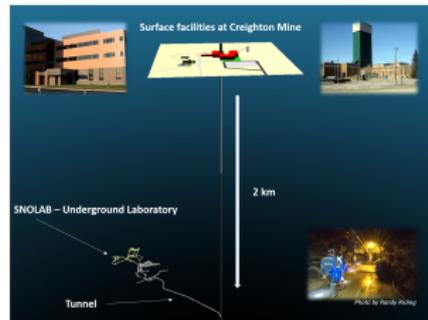
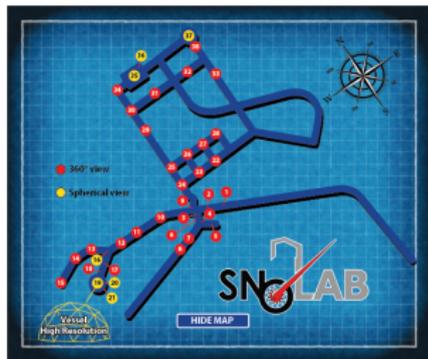
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Deepest and cleanest large-space international facility in the world (UNAM is collaborating with SNOLAB in several experiments since 2015)

- 2 km underground near Sudbury, Ontario
- ultra-low radioactivity background environment Class 2000
- Physics programme focused on neutrino physics and direct dark matter searches

Home of the SNO experiment
 2015 Noble prize in Physics



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Dark Matter Experiment with Argon and Pulse-shape Discrimination:

- Scattered nucleus detected via scintillation
- Pulse shape discrimination for suppression of β/γ events
- LAr advantages
 - is easily purified and high light yield
 - is well understood (ha!)
 - has an easily accessible temperature (85 K)
 - allows a very large detector mass with uniform response
- Detectors:
 - DEAP-1: prototype, 7 kg LAr, 2 PMTs
 - DEAP-3600: 3600 kg LAr, 255 8" PMTs

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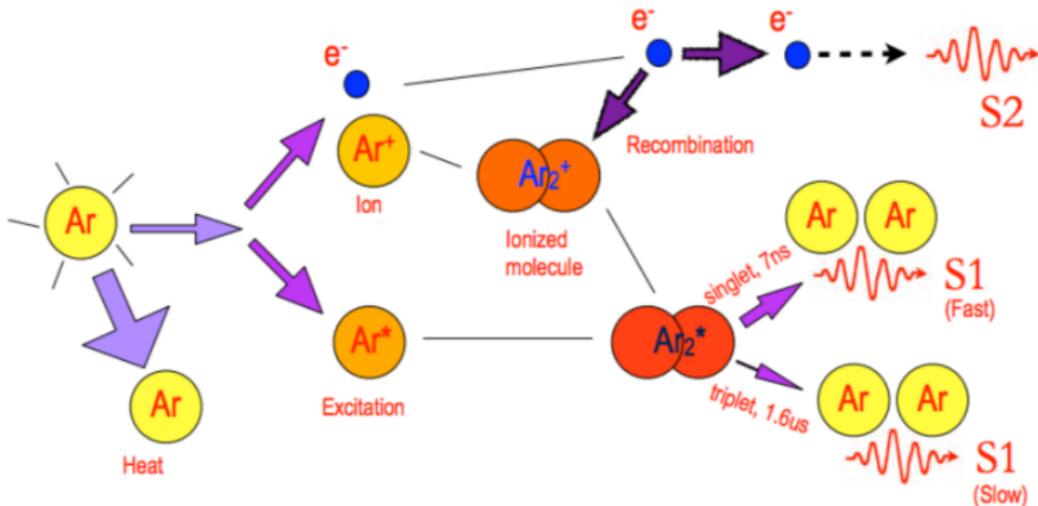
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Argon, nuclear recoil



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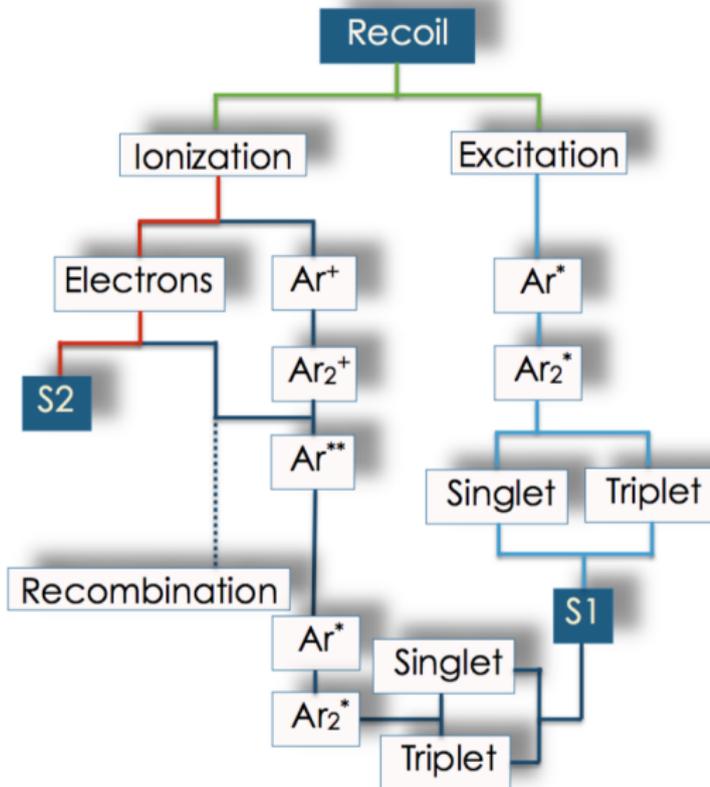
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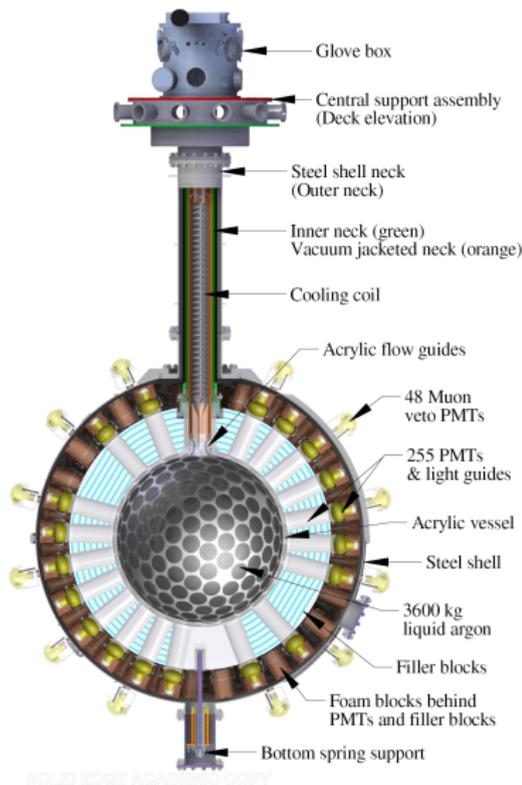
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- Single phase liquid argon: simple, scalable, inexpensive
- 3600 kg argon (1000 kg fiducial) in ultra-clean AV
- Vessel is "resurfaced" in-situ to remove Rn daughters
- TPB wavelength shifter deposition: in-situ vacuum evaporation
- 255 Hamamatsu R5912 HQE 8@ PMTs (32% QE, 75% coverage)
- 50 cm light guides and PE shielding for neutron moderation
- Detector immersed in 8 m water shield tank in Cube Hall



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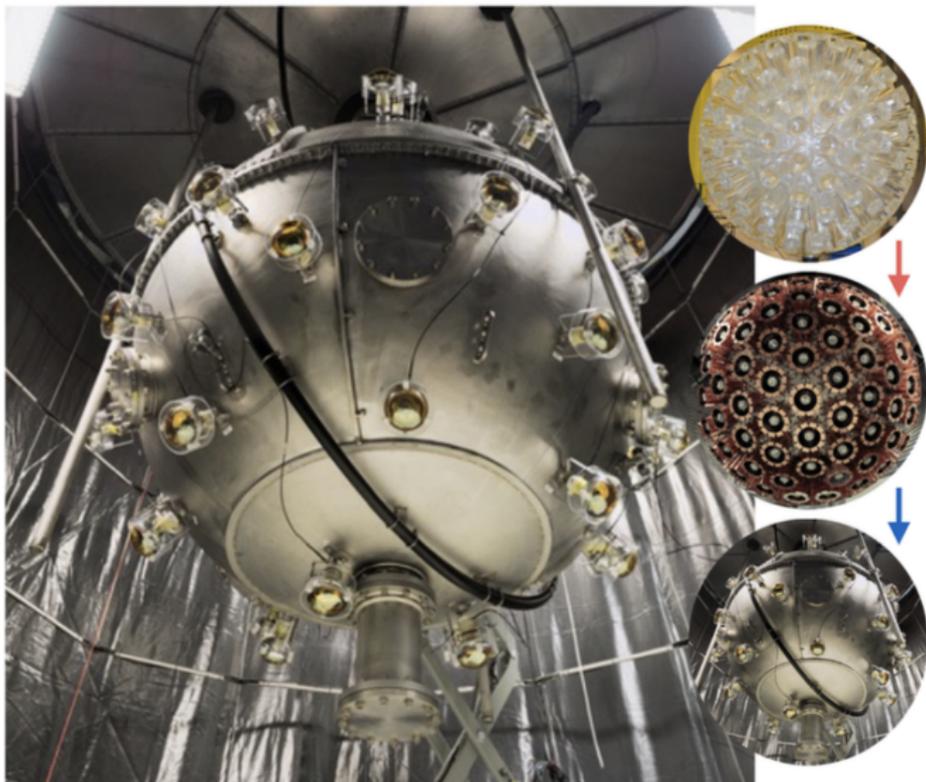
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The three types of interaction that neutrons may undergo in LAr are the following:

- **Elastic scattering** producing nuclear recoils
- **Inelastic collision** leading to γ emission and nuclear recoils
- **Neutron capture** with subsequent emission of a γ and Auger electrons.

Since its event topology is similar as WIMPs, neutron elastic scattering might be an issue because both particles produce a recoiling argon nucleus.

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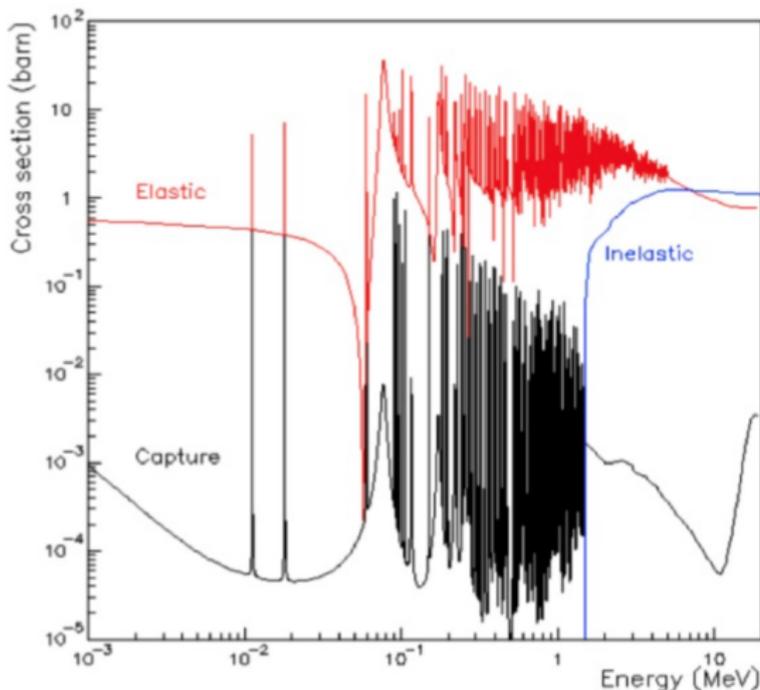


Figure 2.1: Neutron elastic scattering (red), inelastic scattering (blue) and capture (black) cross section on LAr. PhysRevC.85.065811

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- It is well known that for nuclear recoils in liquid noble gas, only a fraction of the energy deposit leads to ionization and scintillation.
- The rest of the energy is transferred to atomic motion and lost in heat without electrically exciting or ionizing argon target.
- This effect is called **nuclear quenching** and is described by the Lindhard theory¹.

¹*Mat. Fys. Medd. Dn. Vid. Selks.* 33(1963), no. 10, 1-42.

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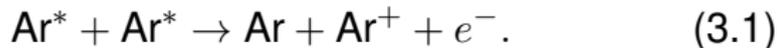
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The luminescence quenching depends upon other processes:

- A.Hitachi and T.Doke²:



- The penning process:



- Superelastic collisions quenching the singlet states to the triplet states³.
- Purity effects.

²PhysRevB.27.5279

³doi:10:1016/0009-2614(76)80566-0

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In general

$$\left(\frac{dE}{dx}\right)_{total} = \left(\frac{dE}{dx}\right)_{elec} + \left(\frac{dE}{dx}\right)_{nucl}. \quad (3.3)$$

Lindhard suggested:

$$E_R = \eta(E_R) + \nu(E_R). \quad (3.4)$$

Then:

$$f_n(E_R) \equiv \frac{\eta(E_R)}{E_R} = \frac{\eta(E_R)}{\eta(E_R) + \nu(E_R)} = \frac{\int_0^{E_R} (dE/dx)_{elec} dE}{\int_0^{E_R} (dE/dx)_{tot} dE}. \quad (3.5)$$

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The last expression needs to be evaluated for any possible recoil energies and can be approximated by

$$f_n = \frac{k \cdot g(\varepsilon)}{1 + k \cdot g(\varepsilon)} \quad (3.6)$$

where $k \approx 0.133Z^{2/3}A^{-1/2}$ and

$$g(\varepsilon) \approx 3\varepsilon^{0.15} + 0.7\varepsilon^{0.6} + \varepsilon. \quad (3.7)$$

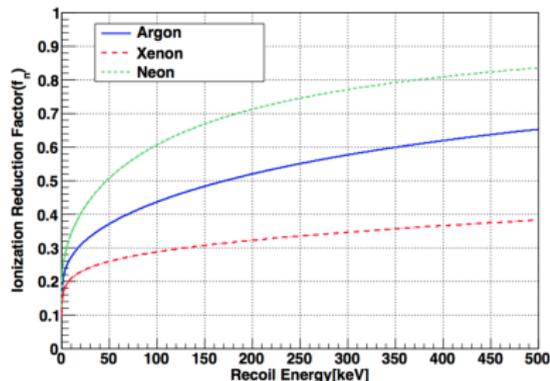


Figure 3.1: Ionization energy reduction factor (f_n) as a function of the recoil energy for argon.

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In the presence of the luminescence quenching, the scintillation light response is described by the Birk's law saturation:

$$\frac{dS}{dx} = \frac{A\left(\frac{dE}{dx}\right)_{elec}}{1 + kB\left(\frac{dE}{dx}\right)_{elec}} \quad (3.8)$$

\implies

$$f_l := \frac{1}{1 + kB\left(\frac{dE}{dx}\right)_{elec}} \quad (3.9)$$

For LAr⁴ $kB = 7.4 \times 10^{-4} \text{MeV}^{-1} \text{gcm}^{-2}$.

⁴PhysRevB.54.15724

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In order to fully describe the luminescence quenching for noble liquid, Mei *et al.* arXiv:0712.2470, combined Lindhard theory f_n and Birk's saturation law f_l .

$$q_n = f_n \times f_l \quad (4.1)$$

where q_n is called the quenching factor or the total scintillation efficiency.

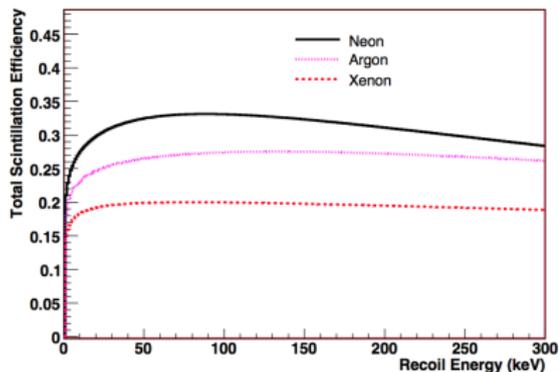


Figure 4.1: Ionization energy reduction factor (f_n) as a function of the recoil energy for argon. arXiv:0712.2470



Relative scintillation efficiency \mathcal{L}_{eff} from Lindhard and Hitachi models

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The relative scintillation efficiency is usually denoted by \mathcal{L}_{eff} and is defined as the ratio of the scintillation yield of nuclear recoil to the scintillation yield of electron recoil from a photoabsorbed γ -source.

$$\mathcal{L}_{eff} = \frac{L_{y,nr}(E_{nr})}{L_{y,er}(E_{er})} \quad (4.2)$$

where the subscripts "nr" and "er" are used for nuclear and electron recoil.



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Two techniques have been employed to measure \mathcal{L}_{eff} :

- **Indirect measurement:** Comparison between a continuum-energy source of neutrons with a simulated spectrum. \mathcal{L}_{eff} is obtained by applying iteratively a fit procedure until the best fit with the measured spectrum is achieved.⁵
- **Direct measurement:** Study of the measured response by recording at fixed angle the single elastic scattered from a monochromatic neutron source.⁶

⁵arXiv:1007.3746

⁶PhysRevC.84.045805 & PhysRevC.85.065811

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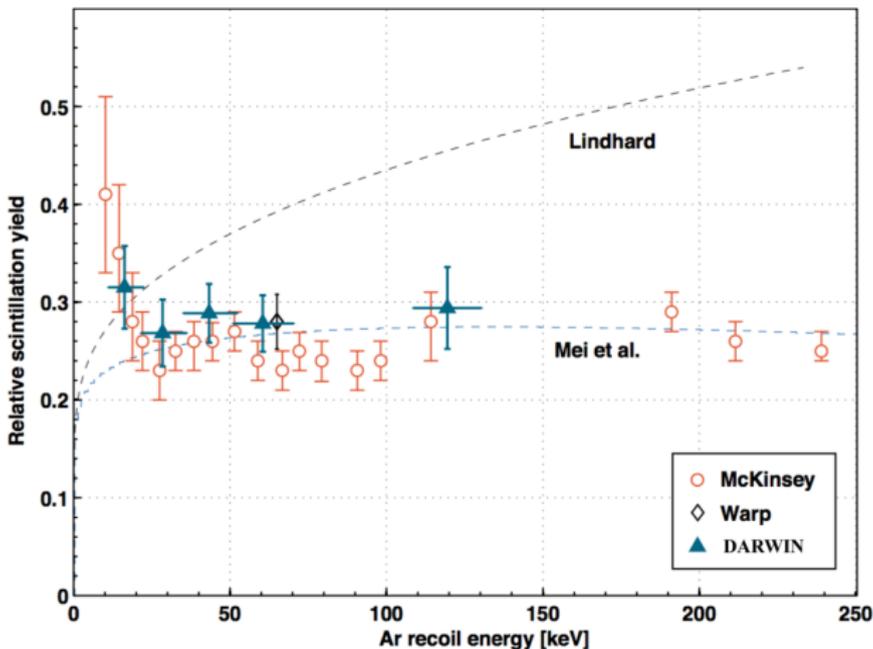


Figure 5.1: Measurements of the relative scintillation efficiency in LAr \mathcal{L}_{eff} as a function of recoil energy. arXiv:1203.0849v1 [astro-ph.IM]

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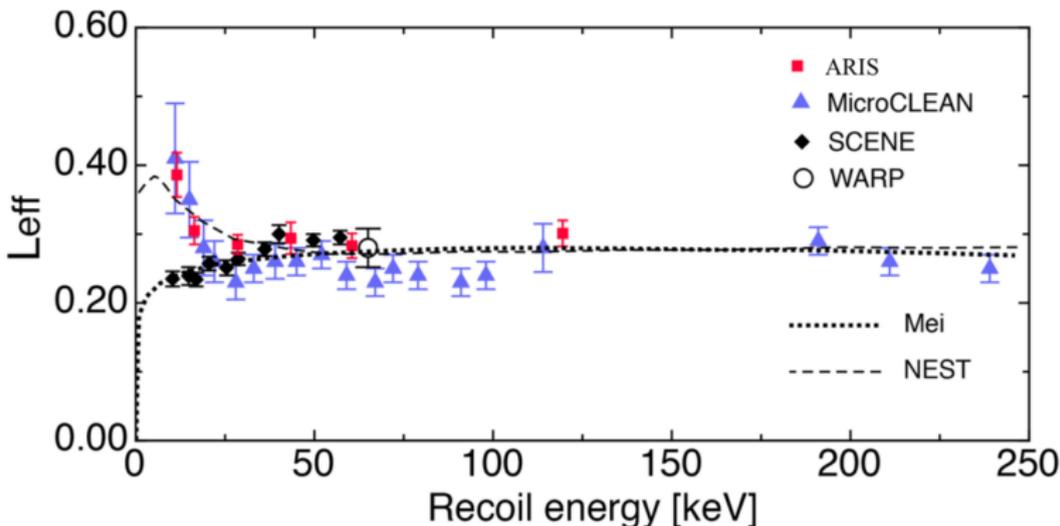


Figure 5.2: Measurements of the relative scintillation efficiency in LAr \mathcal{L}_{eff} as a function of recoil energy. ARIS, P.Agnes.

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- It is necessary to measure the scintillation efficiency down to the energy threshold so as to quantify the WIMP detection sensitivity.
- Plan to measure the LAr response at different NR energies from an AmBe source in the DEAP-3600 experiment (largest operating dark matter detector).

Thank You.

A. Míchell Martínez Mendoza

