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## Characterizing the medium in heavy ion collisions through jets

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- I. Introduction: radiative energy loss.
- 2. Successes and problems.
- 3. Going beyond (what we are currently using).
- 4. Summary and outlook.

Characterizing the medium in HIC through jets.

# I. Introduction: radiative energy loss

I.I.Theoretical setup.

I.2. Models.

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## I.I.Theoretical setup:



Two parameters define the medium: qhat or gluon density plus mean free path, and length (geometry, dynamical expansion).

Characterizing the medium in HIC through jets: 1. Introduction.

## 1.2. Models (I) (Majumder, nucl-th/0702066):

I/2. BDMPS/GLV: static medium.



$$\omega \frac{dI}{d\omega d\mathbf{k}_{\perp}} = \frac{\alpha_s C_F}{(2\pi)^2 \omega^2} 2 \operatorname{Re} \int_0^\infty dy_I \int_{y_I}^\infty d\bar{y}_I e^{i\bar{q}(y_I - \bar{y}_I)} \\ \times \int d\mathbf{u} e^{-i\mathbf{k}_{\perp} \cdot \mathbf{u}} \exp\left(\frac{1}{2} \int_{\bar{y}_I}^\infty d\xi n(\xi) \sigma(\mathbf{u})\right) \\ \times \frac{\partial}{\partial \mathbf{y}} \cdot \frac{\partial}{\partial \mathbf{u}} \int_{\mathbf{y}=0=\mathbf{r}(y_I)}^{\mathbf{u}=\mathbf{r}(\bar{y}_I)} \mathcal{D}\mathbf{r} \exp\left[i \int_{y_I}^{\bar{y}_I} d\xi \frac{\omega}{2} \left(\dot{\mathbf{r}}^2 - \frac{n(\xi) \sigma(\mathbf{r})}{i\omega}\right)\right]$$

Exact solution unknown, two approximations: I. Harmonic oscillator (Brownian motion): multiple soft scatterings.  $n(\xi)\sigma(\mathbf{r}) \simeq \frac{1}{2}\hat{q}(\xi)\mathbf{r}^{2}$ 

2. Opacity expansion: N=I, single hard scattering, corrects Brownian motion.  $[n(\xi)\sigma(\mathbf{r})]^{\mathbf{N}}$ Comparison for massless and massive: SW '03,ASW '04. Characterizing the medium in HIC through jets: I. Introduction.

## I.2. Models (II):

3. AMY: rates order  $\alpha_s$ , dynamical medium, no interference of emissions in/out medium, expansion.

4. GW(M): FF in DIS on nuclei, first corrections in  $L/k_T^2$ , modification of DGLAP splitting functions, virtuality (see also Majumder et al. '07).



$$\tilde{D}(z_{1,}\mu^{2}) = D(z_{1,}\mu^{2}) + \frac{\alpha_{s}}{2\pi} \int_{0}^{\mu^{2}} \frac{dl_{\perp}^{2}}{l_{\perp}^{2}} \int \frac{dy}{y} \left( \frac{1+y^{2}}{1-y} f(x,y,Q^{2},l_{\perp}) + V.C. \right) D(z_{1}/y,\mu^{2})$$

$$f = \frac{C_{A} 2\pi \alpha}{l_{T}^{2} + k_{T}^{2}} \frac{\int dy dy_{1} dy_{2} \left\langle A \left| \overline{\psi}(y) F(y_{1}) F(y_{2}) \psi(0) \right| A \right\rangle e^{i \, factors}}{N_{c} \, f^{A}(x)}$$



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## 2. Successes and problems:

2.1. Light hadrons: R<sub>AA</sub> and back-to-back suppression. :-)

2.2. Non-photonic electrons and more differential observables. :-(

2.3. qhat: dependence on medium modeling. :-(

2.4. Limitations of the formalism. :-(

2.5. Jets (see the talk by G. Salam at HP2008). :-))))))

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#### 2.1. RAA and back-to-back for light:

 $= \int d\Delta E P(\Delta E) \left( \frac{d\sigma^{\rm vac}(p_{\perp} + \Delta E)/dp_{\perp}^2}{d\sigma^{\rm vac}(p_{\perp})/dp_{\perp}^2} \right)$ 



 $Q(p_{\perp}) = \frac{d\sigma^{\text{med}}(p_{\perp})/dp_{\perp}^2}{d\sigma^{\text{vac}}(p_{\perp})/dp_{\perp}^2}$ 

Characterizing the medium in HIC through jets: 2. Successes and problems.

16

12

4.0

 $\approx$ 

 $D_{h/q}^{(\text{med})}(x,Q^2) = \int_0^1 d\epsilon P(\epsilon) \frac{1}{1-\epsilon} D_{h/q}\left(\frac{x}{1-\epsilon},Q^2\right)$ 

AuAu→π<sup>0</sup> 200GeV 0-10%

R<sub>AA</sub> p<sub>r</sub>=8±.5GeV

R<sub>AA</sub> p<sub>T</sub>=20±.5GeV

PbPb $\rightarrow \pi^{\circ}$  5.5TeV 0-10%

2.5

1.5

2.0

ε<sub>o</sub> (GeV/fm)

--- I<sub>AA</sub> p<sub>T</sub><sup>trig</sup>=8±.5GeV p<sub>T</sub><sup>asso</sup>=6±.5GeV

•••••• I p\_t<sup>ig</sup>=20±.5GeV p\_-<sup>asso</sup>=10±.5GeV

3.0

3.5

#### 2.1. RAA and back-to-back for light:

 $\frac{d\sigma^{\rm vac}(p_{\perp} + \Delta E)/dp_{\perp}^2}{d\sigma^{\rm vac}(p_{\perp})/dp_{\perp}^2}$ 



 $= \int d\Delta E P(\Delta E)$ 

 $d\sigma^{\rm med}(p_\perp)/dp_\perp^2$ 

 $Q(p_{\perp}) = \frac{d\sigma \cos(p_{\perp})}{d\sigma^{\operatorname{vac}}(p_{\perp})/dp_{\perp}^{2}}$ 



#### BDMS '01; Wang et al '96 Medium modeling $\rightarrow < T_0 qhat > = I - I 5 GeV^2$







## 2.2. e's, differential observ.:

BAA

10<sup>-1</sup>

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- Heavy quarks radiate less: non-photonic electrons not conclusive: benchmark (Armesto et al '05), hadronization (Adil et al '06), collisional (Djordjevic et al '06, Ayala et al. '07), resonances (van Hees et al '06), dynamical medium (Djordjevic et al. '08),...
- PseudoFF not well understood: no broadening at high  $p_T$  in the near side, trigger bias?







#### 2.3. qhat: medium modeling

$$\hat{q}(\xi) = K \cdot 2 \cdot \epsilon^{3/4}(\xi)$$

$$\langle \hat{q} \rangle = \frac{2}{L^2 - \tau_0^2} \int_{\tau_0}^L d\tau \tau \hat{q}_0 \frac{\tau_0}{\tau} \simeq \frac{2\tau_0 \hat{q}_0}{L} \approx \frac{\hat{q}_0}{2 \div}$$

Gyulassy et al. '01, Salgado et al. '02

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Phenomenological implementation	qhat (GeV²/fm)
fixed length	<~l (average)
Woods-Saxon (PQM)	4-14 (average)
dilution	increases, factor 2-5
dynamical medium (Djordjevic et al.)	decreases
flow (Armesto et al., Baier et al.)	no effect
hydro (Eskola et al., Bass	K~3-4, late times
et al.)	important

### 2.4. Limitations of the formalism:

- Calculations done in the high-energy approximation: only soft emissions.
- Energy-momentum conservation imposed a posteriori in the single inclusive spectrum (GLV; Salgado et al. '03).
- Multiple gluon emission: Quenching Weights (BDMS '01), independent (Poissonian) gluon emission: assumption!

$$P(\Delta E) = \sum_{n=0}^{\infty} \frac{1}{n!} \left[ \prod_{i=1}^{n} \int d\omega_i \frac{dI(\omega_i)}{d\omega} \right] \delta \left( \Delta E - \sum_{i=1}^{n} \omega_i \right) \exp \left[ -\int_0^{\infty} d\omega \frac{dI}{d\omega} \right]$$

No role of virtuality in medium emissions (but GWM!).

• Medium and vacuum treated differently.

## 2.5. Jets (I):

• Jets are theoretical well-defined objects which are as close to QCD partons as we can achieve. They require a definition (clustering algorithm).

• Difficulty in heavy ion collisions (ALICE, ATLAS, CMS): background with large EBE fluctuations, substraction! (Cacciari-Salam).



## 2.5. Jets (II):



At RHIC, the measurement is still biased (a high-p<sub>T</sub> particle required); only at the LHC (RHIC-II?) truly unbiased measurements will be available (e.g. ff. for z<0.7 or  $\xi$ <5). Characterizing the medium in HIC through jets: 2. Successes and problems.



## 3. Going beyond:

Recent attempts to go beyond (JHEP0802 (2008) 048, with L. Cunqueiro, C.A. Salgado, *Santiago* and W.-C. Xiang, *Bielefeld*); also with G. Corcella (*Pisa*).

Motivation: to check radiative eloss, more differential and unbiased observables (particle correlations and jets) have to be studied (others: Borghini et al. '05-..., Wang et al. '01-..., Vitev '05, Polosa et al. '06) → Monte Carlo for in-medium parton branching.

3.1. Medium-modified DGLAP evolution of frag. funct. (FF).

3.2. PYTHIA with in-medium branching: Q-PYTHIA.

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#### 3.1. Medium-modified SF and Sudakovs:

In the vacuum, the formalism gives collinear  $(z \rightarrow I)$  SFs:

$$\frac{dI^{\text{vac}}}{dz \, d\mathbf{k}_{\perp}^{2}} = \frac{\alpha_{s}}{2\pi} \frac{1}{\mathbf{k}_{\perp}^{2}} P^{\text{vac}}(z), \quad P^{\text{vac}}(z) \simeq \frac{2C_{R}}{1-z} \qquad \omega = (1-z)E \text{ and } \mathbf{k}_{\perp}^{2} = z(1-z)E$$
  
In the medium, we make the analogy (ansatz!!!) (Polosa et al. '06):  
$$P^{\text{tot}}(z) = P^{\text{vac}}(z) + \Delta P(z,t) \qquad \Delta P(z,t) \simeq \frac{2\pi t}{\alpha_{s}} \frac{dI^{\text{med}}}{dzdt}$$
  
Medium-modified Sudakovs:  
$$\Delta_{i}(t) = \exp\left[-\int_{t_{0}}^{t} \frac{dt'}{t'} \int_{z_{\min}(t')}^{1-z_{\min}(t')} dz \frac{\alpha_{s}(t',z)}{2\pi} \sum_{j} P_{i \rightarrow j}(z,t')}{z_{min}(t)} \sum_{j} P_{i \rightarrow j}(z,t')} \int_{0}^{1} \frac{\lne^{s/2}}{\lne^{s/2}} \frac{\lne^{s/2}}{\lne^{s/2}} \frac{\lne^{s/2}}{\lne^{s/2}} \int_{0}^{1} \frac{\lne^{s/2}}{\lne^{s/2}} \frac{$$

## 3.1. Medium modified DGLAP evolution of FF (I):

• Medium-modified DGLAP evolution of FF (from KKP i. c.):



## 3.1. Medium modified DGLAP evolution of FF (I):

- Improvements: virtuality in medium emissions, medium and vacuum treated on the same footing, energy momentum conservation.
- Drawbacks: formation time of the gluons does not affect the medium length seen by the radiating partons; energy degradation of the parent not considered; no elastic scattering, no

conversions, no color reconnections with the medium included.





## 3.1. Medium modified DGLAP evolution of FF (II):

Comparison with experimental data gives qhat~I GeV<sup>2</sup>/fm (as with QW for fixed L) or qhat~I0 GeV<sup>2</sup>/ fm (for cylinder or sphere).

QW reproduced (numerically and analytically) for high energies and virtualities.



#### 3.2. Q-PYTHIA: implementation

• We modify PYTHIA FSR routine PYSHOW, for t-ordered branching, introducing the medium-modified splittings into the dicing of both t and z.  $P_{tot}(z) = P_{vac}(z) \rightarrow P_{tot}(z) = P_{vac}(z) + \Delta P(z, t, \hat{q}, L, E)$ 



Energy degradation and formation time considered.
Rejections undone: agreement between vacuum and default as a check.



p<sub>T</sub> (GeV)

ξ=In(E\_i/p)

Characterizing the medium in HIC through jets: 3. Beyond.

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## 3.2. Q-PYTHIA: program

- Fortran program, uses PYTHIA-6.4.18 defaults (arXiv:0809.4433 [hep-ph]).
- We modify \*only\* PYSHOW, providing additional auxiliary routines: black box for the user.
- User-defined:
  - \* Position of hard scattering of jet origin (x<sub>0</sub>,y<sub>0</sub>,z<sub>0</sub>,t<sub>0</sub>): QPYGIN.
  - \*Values of [qhat L] and  $[\omega_c = qhat L^2/2]$  at point of branching  $(x,y,z,t,\beta_x,\beta_y,\beta_z)$ : QPYGEO; which contains medium modeling.



## 3.3. Q-PYTHIA: results (I)



## 3.3. Q-PYTHIA: results (II)

At parton level and for small energy, some multiplicity enhancement and modest broadening.



## 3.3. Q-PYTHIA: results (III)

At parton level and for high energy, high multiplicity enhancement and modest broadening. Energy constraints and multiple splitting very important.



## 3.3. Q-PYTHIA: results (IV)

Two back-toback gluons in singlet: an extreme example of hadronization. It kills most of the multiplicity enhancement, medium effects (soft stuff) less evident.



## 3.3. Q-PYTHIA: results (V)



### 3.3. Q-PYTHIA: others

• PQM (Dainese-Loizides-Paic): PYTHIA with radiative eloss via QW.

• PYQUEN (Lokhtin-Snigirev): elastic scattering on top of radiative energy loss à la BDMPS with additional included gluons (path length considerations).

• JEWEL (Zapp-Ingelman-Rathsman-Stachel-Wiedemann): elastic scattering in DGLAP evolution plus radiative energy loss through a multiplicative constant in the collinear part of the splitting functions (Borghini-Wiedemann).

• YaJEM (Renk): increase of virtuality (i.e. of the length of the evolution) through qhatL computed at every point.

## 4. Summary and outlook:

- To check radiative energy loss as the explanation for jet quenching, differential probes needed: relation energy degradation / radiation enhancement / p<sub>T</sub> broadening.
- We have supplemented vacuum splitting functions with medium terms: virtuality, energy conservation, vacuum and medium treated on the same footing. A modified DGLAP evolution for fragmentation functions has been performed, its compatibility with QW for high virtualities showed.
- A publicly available (<u>http://igfae.usc.es/qatmc/</u>) medium-modified Monte Carlo parton shower: Q-PYTHIA.I.0 has been shown (Q-HERWIG to follow), required for correlations and jet shape studies. Future: jet reconstruction, color reconnections, energy flow from/to the medium, theoretical basis.

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