



Monochromatization for Direct Higgs Production

in Future Circular e+e- Colliders

M. Alan Valdivia valdivia@fisica.ugto.mx Universidad de Guanajuato

RAFAE 17, Hotel San Francisco, Tlaxcala City, 280917

Outline

- Motivation
- Objectives
- Introduction: Accelerator Physics
- Standard Monochromatization
- Radiation Effects In FCC
- Optimized Monochromatization
- Conclusion
- Discussion

Motivation

- Questions still unanswered about Higgs boson properties
- Post-LHC high energy resolution requirements
- Upcoming radiation Effects
- FCCe+e- proposal opportunity

Objectives

- Characterize the radiation effects at the FCCe+e-
- Quantify the IP parameters modification due to the radiation effects.
- Define the IP parameters to produce monochromatization
- Develop an optimized monochromatization scheme for the FCCe+e-

Colliding Particles A Qualitative Description

- Particles are produced at the **source**
- Particles are injected into the vacuum chamber
- Stored particles loss energy by radiation, compensated by the RFS
- Particles are focused by guiding fields toward a **designed orbit**
- Accelerating field collects particles into circulating **bunches**
- Amplitud **damping** occurs due to energy loss by radiation
- Amplitude excitation occurs as a quantum effect
- Between damping and excitation a **balance** is reached
- Collision are produced at defined **interaction points**
- Monochromatization requires parameters optimization



Colliders Luminosity



A proposal Future Circular e+e- Collider



- Luminosity is expected of order 10e36
- Energy range per beam should be from 40 to 175 GeV
- Circumference would be of 100 km
- Energy Resolution requires an improvement for direct Higgs production

Beam Dynamics A Quantitative Description



Horizontal Plane

$$F_x = \dot{p_x} = \frac{d}{dt}(\gamma m v_x) = \gamma m a_x$$

Vertical Plane $F_y = \dot{p_y} = \frac{d}{dt}(\gamma m v_y) = \gamma m a_y$

$$B_x = \left(\frac{\partial B_y}{\partial x}\right) y$$

$$B_y = B_0 + \left(\frac{\partial B_y}{\partial x}\right) x$$

Linear Approximation + Paraxial Approximation $R = \rho + x$ $x \ll \rho$

$$\frac{\partial^2 y}{\partial \theta^2} + ny = 0 \qquad \qquad \frac{\partial^2 x}{\partial \theta^2} + (1 - n)x = 0 \qquad \qquad n \equiv -\frac{\rho}{B_0} \left(\frac{\partial B_y}{\partial x}\right)$$

Hill Differential Equation

$$x'' + K(s)x = 0 \qquad y'' - K(s)y = 0, \qquad K(s) \equiv \frac{1}{(B\rho)} \left(\frac{\partial B_y}{\partial x}\right)(s)$$
$$K(s+C) = K(s)$$

Ansatz
$$x(s) = A\omega(s)\cos(\phi(x) + \phi_0)$$

Courant-Snyder Parameters

$$\beta(s) \equiv rac{\omega^2(s)}{k} \qquad \alpha(s) \equiv -rac{1}{2} eta\prime(s) \qquad \gamma(s) \equiv rac{1+lpha^2(s)}{eta(s)}$$

$$\begin{array}{c} & & & \\ & & \\ \hline & & \\ &$$

Emittance

$$\epsilon = \pi \mathcal{W} = \pi [\gamma(s)x^{2}(s) + 2\alpha(s)x(s)x'(s) + \beta(s)x'^{2}(s)]$$

Size of the Distribution

$$\sigma_{RMS} = \sqrt{\epsilon \beta(s)}$$

Lattice Design Dispersion



General Contribution

 $x'' + K(s)x = \frac{\delta}{\rho(s)}$

Dispersion Contribution

Dispersion Function $D(s) = S(s) \int_0^s \frac{C(\tau)}{\rho(\tau)} d\tau - C(s) \int_0^s \frac{S(\tau)}{\rho(\tau)} d\tau$ $x(s) = C(s)x_0 + S(s)x'_0 + D(s)\delta_0$



Equation of Motion for Off-Energy Particles

$$x'' = -K_x x + G\left(\frac{\epsilon}{E_0}\right)$$

Dispersion Equation $D''(s) = -K_x D(s) + G(s)$ $G(s) = \rho^{-1}(s)$

Beam Parameters Distributions

Transverse Displacement $x = x_{\beta} + D_x \epsilon_0$, $y = y_{\beta} + D_y \epsilon_0$ $\epsilon_0 = \frac{\epsilon}{E_0}$ Average at the IP $\langle \mathcal{A} \rangle^{\pm} = \int f^{\pm}(X^{\pm})\mathcal{A}(X^+, X^-)dX^{\pm}$ $\langle \mathcal{A} \rangle^* = \left\langle \langle \mathcal{A} \rangle^+ \right\rangle^- = \left\langle \langle \mathcal{A} \rangle^- \right\rangle^+$

$$f^{\pm}(x,y,\epsilon) = \frac{f^{\pm}(p_x,p_y,z)}{\sqrt{8\pi^3\beta_x^*\epsilon_{xc}\beta_y^*\epsilon_{yc}\sigma_\epsilon^2}} \exp\left\{-\frac{(x_\beta + D_x\epsilon_0)^2}{2\beta_x^*\epsilon_{xc}} - \frac{(y_\beta + D_y\epsilon_0)^2}{2\beta_y^*\epsilon_{yc}} - \frac{\epsilon^2}{2\sigma_\epsilon^2}\right\}$$

Average over distribution

$$\sigma_x^* = \sqrt{\langle x^2 \rangle^{\pm}} = \sqrt{\beta_x^* \epsilon_{xc} + D_x^{*2} \sigma_\epsilon^2}$$
$$\sigma_y^* = \sqrt{\langle y^2 \rangle^{\pm}} = \sqrt{\beta_y^* \epsilon_{yc} + D_y^{*2} \sigma_\epsilon^2}$$
$$\sigma_\epsilon^* = \sqrt{\langle \epsilon_0^2 \rangle^{\pm}}$$

Position-Energy Correlation

$$egin{aligned} &\langle x\epsilon_0
angle^{\pm} = D_x^{*\pm} \sigma_\epsilon^2 \ &\langle y\epsilon_0
angle^{\pm} = D_y^{*\pm} \sigma_\epsilon^2 \end{aligned}$$

Notes

- **Relative Momentum Deviation** of particles differs from ideality
- **Dispersion** quantifies the effects on displacement due to off-momentum particles
- **Emittance** is a measure of how far our distribution is from ideal orbit
- **Beta Function** is related to the physical transverse size of bunches

Collision Energy Spread

 $\sigma_w = \sqrt{2}E_0\sigma_\epsilon$

Beam Energy Spread

$${\sigma_\epsilon}^2 \propto \frac{55 \hbar c {E_0}^2}{32 \sqrt{3(mc^2)^3}} \frac{I_3}{I_2} \frac{1}{J_\epsilon}$$

Function of radius and J_{ε}

$$\sigma_w \propto (\rho J_\epsilon)^{-1/2}.$$

 $J_\epsilon \in [0.5, 2.5]$

Typical Options

$$\left. \begin{array}{c} \rho >> \rho_0 \\ J_{\epsilon} > J_{0\epsilon} \end{array} \right\} \Rightarrow \sigma_w < \sigma_{0w}$$

Baseline Scheme

 $egin{aligned} D^*_{x^+} &= -D^*_{x^-} &= 0 \ D^*_{y^+} &= D^*_{y^-} &= 0 \ \mathcal{L} &= \mathcal{L}_0 \end{aligned}$

Average over distribution

$$\sigma_x^* = \sqrt{\langle x^2 \rangle^{\pm}} = \sqrt{\beta_x^* \epsilon_{xc} + D_x^{*2} \sigma_\epsilon^2} \qquad \langle x \epsilon_0 \rangle^{\pm} = D_x^{*\pm} \sigma_\epsilon^2$$

$$\sigma_y^* = \sqrt{\langle y^2 \rangle^{\pm}} = \sqrt{\beta_y^* \epsilon_{yc} + D_y^{*2} \sigma_\epsilon^2} \qquad \langle y \epsilon_0 \rangle^{\pm} = D_y^{*\pm} \sigma_\epsilon^2$$

Monochromatization principle A "simple explanation"

Standard collision Dispersion has the same sign in the IP

$$e^{-} \xrightarrow{E + \Delta E} \overleftarrow{E} + \Delta E} \overrightarrow{E} e^{-}$$

 $\xrightarrow{E - \Delta E} \overleftarrow{E} - \Delta E} e^{-}$
 $w = 2(E_0 + \epsilon)$

Monochromatization Dispersion has opposite sign in the IP

$$e^{-} \xrightarrow{E + \Delta E} \stackrel{E}{\longleftrightarrow} \stackrel{E - \Delta E}{\xleftarrow{E}} e^{+}$$

 $w = 2E_0 + 0(\epsilon)^2$

Monochromatization Factor

$$\lambda \equiv rac{\mathcal{L}_0}{\mathcal{L}}$$

Both Opposite Sign

$$D_{x^+}^* = -D_{x^-}^* = D_x^* \qquad D_{y^+}^* = -D_{y^-}^* = D_y^*$$

$$\mathcal{L} = \frac{\mathcal{L}_0}{\sqrt{1 + \sigma_\epsilon^2 \left(\frac{D_x^{*2}}{\sigma_{x\beta}^{*2}} + \frac{D_y^{*2}}{\sigma_{y\beta}^{*2}}\right)}} \qquad \Sigma_w = \frac{\sqrt{2E_0\sigma_\epsilon}}{\sqrt{1 + \sigma_\epsilon^2 \left(\frac{D_x^{*2}}{\sigma_{x\beta}^{*2}} + \frac{D_y^{*2}}{\sigma_{y\beta}^{*2}}\right)}} \qquad \lambda = \sqrt{1 + \sigma_\epsilon^2 \left(\frac{D_x^{*2}}{\sigma_{x\beta}^{*2}} + \frac{D_y^{*2}}{\sigma_{y\beta}^{*2}}\right)}$$

$$D_{x^+}^* = -D_{x^-}^* = D_x^*$$
 $D_{y^+}^* = D_{y^-}^* = D_y^*$

$$\Sigma_{w} = \frac{\sqrt{2}E_{0}\sigma_{\epsilon}}{\sqrt{1 + \sigma_{\epsilon}^{2}\left(\frac{D_{x}^{*2}}{\sigma_{x\beta}^{*2}}\right)}} \qquad L = \frac{L_{0}}{\sqrt{1 + \sigma_{\epsilon}^{2}\left(\frac{D_{x}^{*2}}{\sigma_{x\beta}^{*2}}\right)}\sqrt{1 + \sigma_{\epsilon}^{2}\left(\frac{D_{y}^{*2}}{\sigma_{y\beta}^{*2}}\right)}} \qquad \lambda = \sqrt{1 + \sigma_{\epsilon}^{2}\left(\frac{D_{x}^{*2}}{\sigma_{x\beta}^{*2}}\right)}$$

Monochromatization Factor λ

Standard Monochromatization

Monochromatization Factor

$$\lambda = \sqrt{1 + \sigma_{\epsilon}^2 \left(\frac{D_x^{*2}}{\sigma_{x\beta}^{*2}}\right)}$$

 $\mathcal{L} \propto rac{1}{\lambda}$

 $D_{y^+}^* = D_{y^-}^* = 0$





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Notes

- **Baseline Parameters** fail to produce the required center-of-mass energy resolution for direct Higgs production at 125 GeV
- Higgs Width of ~4.2 MeV requires an improvement in resolution by at least a factor of 10 since baseline produced a ~40 MeV width distribution
- **Double Ring System** of the FCCee may produced dispersion at the interaction point as required for monochromatization.
- **Radiation Effects** could compromise the performance of a monochromatization scheme

Radiation Effects: Synchrotron radiation Radiation Damping



Length Increment

$$\delta l_\epsilon = \oint G(s) x_\epsilon ds = rac{\epsilon}{E_0} \oint G(s) D_x(s) ds$$

Dilation Factor

$$rac{\delta l_\epsilon}{L} = lpha rac{\epsilon}{E_0} \qquad \qquad lpha = rac{1}{L} \oint G(s) D_x(s) ds$$

Relative Revolution Time Increment

$$\frac{\delta t}{T_0} = \frac{\delta l}{L} = \alpha \frac{\epsilon}{E_0}$$

Time Displacement Evolution

$$\delta z = -\alpha \frac{\epsilon}{E_0} L$$

$$\frac{d\tau}{dt} = -\alpha \frac{\epsilon}{E_0}$$

Synchrotron radiation Radiation Damping

Energy Change $\delta U = eV(\tau_1) - U_{rad}(\epsilon)$ RF Energy Supply $eV(\tau) = U_{rf}(\bar{t}_s - \tau)$ $U_{rad} = U_0 + D\epsilon, \qquad D = \left(\frac{dU_{rad}}{d\epsilon}\right)_{E_0}$ **Electric Field** \otimes ⊗ Particle Trajectory Vacuum $eV(\tau)$ ⊗ \otimes $\int U_0$ Chamber \odot Magnetic **RF** Cavity \odot Field **RF** Power Feed In Complex Notation **Energy Evolution** Time Displacement Evolution $\frac{d\tau}{dt} = -\alpha \frac{\epsilon}{E_0}$ $rac{d\epsilon}{dt} = rac{1}{T_0}(e\dot{V}_0 au - D\epsilon)$ $\widetilde{\epsilon} = -i \frac{\Omega E_0}{\widetilde{\tau}} \widetilde{\tau}$ Time Displacement Evolution

$$\frac{d^2\tau}{dt^2} + 2\alpha_{\epsilon}\frac{d\tau}{dt} + \Omega^2\tau = 0 \qquad \qquad \alpha_{\epsilon} = \frac{D}{2T_0} \qquad \Omega^2 = \frac{\alpha e\dot{V}_0}{T_0E_0}$$

Synchrotron radiation Quantum Radiation Effects



$$\frac{dN_{\gamma}}{dt} = \int_0^\infty \frac{dn_{\gamma}}{dt}(u)du \qquad < u > = \left(\frac{dN_{\gamma}}{dt}\right)^{-1} \int_0^\infty u \frac{dn_{\gamma}}{dt}du \qquad < u^2 > = \left(\frac{dN_{\gamma}}{dt}\right)^{-1} \int_0^\infty u^2 \frac{dn_{\gamma}}{dt}du$$

 $\frac{dN_{\gamma}}{dt} = \frac{15\sqrt{3}}{8} \frac{P_{\gamma}}{u_c} \qquad < u > = \frac{8}{15\sqrt{3}} u_c \qquad < u^2 > = \frac{11}{27} u_c^2$

Synchrotron radiation Quantum Radiation Effects

Energy Deviation Oscillation $i\Omega(t-t_{2})$

 $\epsilon = A_0 \exp^{i\Omega(t-t_0)}$

Photon Emission

$$\epsilon = A_0 \exp^{i\Omega(t-t_0)} - u \exp^{i\Omega(t-t_i)}$$
$$\epsilon = A_1 \exp^{i\Omega(t-t_1)}$$

New Amplitude

$$A_1^2 = A_0^2 + u^2 - 2A_0 \cos \Omega (t_i - t_0)$$

Probable Amplitude Change

$$<\delta A^2>==u^2$$



Probable Amplitude Squared

$$\frac{d\left\langle A^2\right\rangle}{dt} = -2\frac{\left\langle A^2\right\rangle}{\tau_{\epsilon}} \qquad \qquad \left\langle \frac{dA^2}{dt}\right\rangle = \frac{d\left\langle A^2\right\rangle}{dt} = \frac{dN_{\gamma}}{dt}u^2$$

For Sinusoidal Energy Oscillation

$$\left\langle A^2 \right\rangle = rac{1}{2} au_\epsilon rac{dN_\gamma}{dt} u^2 \qquad \sigma_\epsilon^2 = \left\langle \epsilon^2 \right\rangle = rac{\left\langle A^2 \right\rangle}{2} = rac{1}{4} au_\epsilon rac{dN_\gamma}{dt} u^2$$

Excitation Term

 $< u^2 > \left(\frac{dN_{\gamma}}{dt}\right) = \frac{55}{24\sqrt{3}} r_e \hbar m c^4 \frac{\gamma^7}{\rho^3}$

Beamstrahlung Sokolov-Ternov Theory



Beamstrahlung Sokolov-Ternov Theory



Self-Consistent Equations

$$\epsilon_{x,\text{tot}} = \epsilon_{x,\text{SR}} + \frac{\tau_x n_{\text{IP}}}{4T_{\text{rev}}} \{ n_\gamma < u^2 > \} \mathcal{H}_x^* \qquad \qquad \sigma_{\delta,\text{tot}}^2 = \sigma_{\delta,\text{SR}}^2 + \frac{n_{\text{IP}} \tau_{E,\text{SR}}}{4T_{\text{rev}}} \{ n_\gamma < u^2 > \}$$

Bunch Length
$$\sigma_{z,\text{tot}} = \frac{\alpha_{\text{C}}C}{2\pi Q_s} \sigma_{\delta,\text{tot}}$$

$$\mathcal{H}_{x}^{*} \equiv rac{\left(eta_{x}^{*} {D_{x}^{\prime}}^{*} + lpha_{x}^{*} {D_{x}^{*}}
ight)^{2} + {D_{x}^{*2}}}{eta_{x}^{*}}$$

Baseline monochromatization

E_e [GeV]	62.5
scheme	m.c.
	basel.
I_b [mA]	408.3
N_b [10 ¹⁰]	3.3
n_b [1]	25760
n_{IP} [1]	2
β_x^* [m]	1.0
$\beta_y^* [\text{mm}]$	2
$\check{D_x^*}$ [m]	0.22
$\epsilon_{x,\mathrm{SR}} \mathrm{[nm]}$	0.17
$\epsilon_{x,\mathrm{tot}} \mathrm{[nm]}$	0.21
$\epsilon_{y,\mathrm{SR}} \mathrm{[pm]}$	1
$\sigma_{x,\mathrm{SR}}~[\mathrm{\mu m}]$	132
$\sigma_{x,\mathrm{tot}}~[\mathrm{\mu m}]$	144
$\sigma_y [{ m nm}]$	45
$\sigma_{z,\mathrm{SR}} [\mathrm{mm}]$	1.8
$\sigma_{z,{ m tot}} [{ m mm}]$	1.8
$\sigma_{\delta,\mathrm{SR}}$ [%]	0.06
$\sigma_{\delta,\mathrm{tot}}$ [%]	0.06
$\theta_c \; [\mathrm{mrad}]$	0
circ. C [km]	100
$lpha_{ m C}~[10^{-6}]$	7
$f_{ m rf}~[m MHz]$	400
$V_{ m rf}~[m GV]$	0.4
$U_{0,\mathrm{SR}} \; [\mathrm{GeV}]$	0.12
$U_{0,\mathrm{BS}} \; \mathrm{[MeV]}$	0.01
$ au_E/T_{ m rev}$	509
Q_s	0.030
$\Upsilon_{ m max}$ $[10^{-4}]$	0.3
$\Upsilon_{\rm ave} [10^{-4}]$	0.1
$\theta_c \; [\mathrm{mrad}]$	0
$\xi_x [10^{-2}]$	1
$\xi_y [10^{-2}]$	4
λ [1]	9.2
$L [10^{35}]$	1.0
$cm^{-2}s^{-1}$]	
$\sigma_w [{ m MeV}]$	5.8

Width of standard model Higgs 4-5 MeV requires $\lambda \ge 10$ N_b = 3.3x10¹⁰ n_b = 25760 β_y * = 2 mm

Monochromatization Factor 0.5 1 1.5 2 Ĕ ∗__× 2.5 9 3 3.5 4 4.5 5 0.05 0.1 0.15 0.2 0.25 D _ * [m] 2 4 6 8 10 12 14 16 18 20

 \star : λ = 9.2, L = 1 x 10³⁵ cm⁻²s⁻¹

Baseline monochromatization

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Width of standard model Higgs 4-5 MeV requires $\lambda \ge 10$ $N_b = 3.3 \times 10^{10} n_b = 25760 \beta_y^* = 2 mm$ Luminosity contours -00e+35 4.5 3.5 1.36e+35 2.08e+35 1.90e+35 1.54e+35 1.18e+35 1.72e+35 2.26e+35 44e+35 2e+35 <u>ع</u> * 13035 × 2.5 β 1.358138.59e+35 1.5 0.5 0.05 0.1 0.15 0.2 0.25 D _ * [m]

 \star : λ = 9.2, L = 1 x 10³⁵ cm⁻²s⁻¹

Optimized Monochromatization Luminosity_max and λ

S =
$$[0.1, 3]$$
, T = $[0.1, 3]$; $\beta_x = \beta_{0x} * S^2$, $D_x * = D_{0x} * S$; $N_b = N_{0b} / T$, $n_b = n_{0b} * T$

 $\beta_{0x} = 1.0 \text{ m}, \beta_{0y} = 1.0 \text{ mm}, D_{0x} = 0.22 \text{ m}, N_{0b} = 3.3 \times 10^{10}, n_{0b} = 25760 \quad \bigstar : \lambda_0 = 10.17321$



Luminosity contours





Optimized Monochromatization Luminosity max and λ

S = [0.1, 3], T = [0.1, 3]; $\beta_x = \beta_{0x} * S^2$, $D_x * = D_{0x} * S$; $N_b = N_{0b} / T$, $n_b = n_{0b} * T$

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Optimized Monochromatization Luminosity_max and λ

S = [0.1, 3], T = [0.1, 3]; $\beta_x = \beta_{0x} * S^2$, $D_x * = D_{0x} * S$; $N_b = N_{0b} / T$, $n_b = n_{0b} * T$

 $\beta_{0x} = 1.0 \text{ m}, \beta_{0y} = 1.0 \text{ mm}, D_{0x} = 0.22 \text{ m}, N_{0b} = 3.3 \times 10^{10}, n_{0b} = 25760$



Monochromatization Factor

 \star : λ = 5.07, β = 1.96 m, D_x = 0.308 m, L = 3.736 x 10³⁵ cm⁻²s⁻¹

Optimized Monochromatization Conclusions

E_e [GeV]	45.6	62.5	62.5	62.5	80
scheme	CW	ho.	m.c.	m.c.	CW
			basel.	opt'd	
I_b [mA]	1450.3	408.3	408.3	408.3	151.5
N_b [10 ¹⁰]	3.3	1.05	3.3	11.1	6.0
n_b [1]	91500	80960	25760	7728	5260
n_{IP} [1]	2	2	2	2	2
β_x^* [m]	1	1.0	1.0	1.96	1
β_y^* [mm]	2	2	2	1	2
D_x^* [m]	0	0	0.22	0.308	0
$\epsilon_{x,\mathrm{SR}} \; [\mathrm{nm}]$	0.09	0.17	0.17	0.17	0.26
$\epsilon_{x,\mathrm{tot}} \; [\mathrm{nm}]$	0.09	0.17	0.21	0.70	0.26
$\epsilon_{y,\mathrm{SR}} \; [\mathrm{pm}]$	1	1	1	1	1
$\sigma_{x,\mathrm{SR}} \ [\mu\mathrm{m}]$	9.5	9.2	132	185.7	16
$\sigma_{x,{ m tot}} \; [\mu{ m m}]$	9.5	9.2	144	188.5	16
$\sigma_y \text{ [nm]}$	45	45	45	32	45
$\sigma_{z,\mathrm{SR}} [\mathrm{mm}]$	1.6	1.8	1.8	1.8	2.0
$\sigma_{z,\mathrm{tot}}$ [mm]	3.8	1.8	1.8	1.8	3.1
$\sigma_{\delta,\mathrm{SR}}$ [%]	0.04	0.06	0.06	0.06	0.07
$\sigma_{\delta,\mathrm{tot}}$ [%]	0.09	0.06	0.06	0.06	0.10
$\theta_c \text{ [mrad]}$	3 0	0	0	0	30
circ. C [km]	100	100	100	100	100
$lpha_{ m C}~[10^{-6}]$	7	7	7	7	7
$f_{ m rf}~[m MHz]$	400	400	400	400	400
$V_{ m rf}~[m GV]$	0.2	0.4	0.4	0.4	0.8
$U_{0,\mathrm{SR}} \; [\mathrm{GeV}]$	0.03	0.12	0.12	0.12	0.33
$U_{0,\mathrm{BS}} \; \mathrm{[MeV]}$	0.5	0.05	0.01	0.01	0.21
$ au_E/T_{ m rev}$	1320	509	509	509	243
Q_s	0.025	0.030	0.030	0.030	0.037
$\Upsilon_{ m max}$ $[10^{-4}]$	1.7	0.8	0.3	0.85	4.0
$\Upsilon_{\rm ave} \ [10^{-4}]$	0.7	0.3	0.1	0.35	1.7
$\theta_c [\text{mrad}]$	30	0	0	0	30
$\xi_x [10^{-2}]$	5	12	1	2.22	7
$\xi_y [10^{-2}]$	13	15	4	6.76	16
λ [1]	1	1	9.2	5.08	1
$L \ [10^{35}]$	9.0	2.2	1.0	3.74	1.9
$cm^{-2}s^{-1}]$					
$\sigma_w ~[{ m MeV}]$	58	53	5.8	10.44	113
+					

- Monochromatization scheme can be implemented
- Beamstrahlung effects may be predicted
- Simulation (!) supports predictions
- Lattice designed is still in progress and the required modification should be possible
- **Theory** confirmation could be achieved at the FCCe+e-
- Alternative optimization techniques under study

Participations

2017 IMC17: Talk, Cancún City

2017 Seminario IFUNAM:' Talk, IF-UNAM, Mexico City

2017 División de Partículas y Campos Meeting, Talk, (Co-Author) Poster (2), CINVESTAV, Mexico City

2016 IPAC (Author) Poster (1) Contribution Paper (1) to Proceedings Conference, Copenhagen, Denmark

2016 Red FAE Meeting & Talk Pachuca, Hidalgo

2016 CERN-BINP Workshop Talk / Contribution Paper to Proceedings Workshop Geneva, Switzerland

2016 Latin-American Conference High Energy Physics: Particle and Strings Talk, Havana, Cuba

2016 IPAC (Presenter) Poster (2) Contribution Paper (2) to Proceedings Conference, Busan, Korea

2016 USPAS; Fundamentals of Accelerator Physics and Technology, School Austin, Texas, USA

2015 Mexican Particle Accelerator School & Talk; Guanajuato, Guanajuato, México

2015 FCC-ee Optics Design 33rd meeting *Talk*, CERN, France/Switzerland.

2015 Summer project at CERN as associated member of the personnel, CERN, Talk France/Switzerland.

2014 Seminario DCI (IFUG); Talk, León, Guanajuato.

2014 Science Undergraduate Summer Laboratory Internships; lectures on HEP School León, Guanajuato.





Monochromatization for Direct Higgs Production

in Future Circular e+e- Colliders

M. Alan Valdivia valdivia@fisica.ugto.mx Universidad de Guanajuato

RAFAE 17, Hotel San Francisco, Tlaxcala City, 280917