

Vector boson tagged jets and jet substructure

Ivan Vitev^{1,*}

¹Los Alamos National Laboratory, Theoretical Division, Mail Stop B283
Los Alamos, NM 87545, USA

Abstract. In these proceedings, we report on recent results related to vector boson-tagged jet production in heavy ion collisions and the related modification of jet substructure, such as jet shapes and jet momentum sharing distributions. Z^0 -tagging and γ -tagging of jets provides new opportunities to study parton shower formation and propagation in the quark-gluon plasma and has been argued to provide tight constraints on the energy loss of reconstructed jets. We present theoretical predictions for isolated photon-tagged and electroweak boson-tagged jet production in Pb+Pb collisions at $\sqrt{s_{NN}} = 5.02$ TeV at the LHC, addressing the modification of their transverse momentum and transverse momentum imbalance distributions. Comparison to recent ATLAS and CMS experimental measurements is performed that can shed light on the medium-induced radiative corrections of predominantly quark-initiated jets. The modification of parton splitting functions in the QGP further implies that the substructure of jets in heavy ion collisions may differ significantly from the corresponding substructure in proton-proton collisions. Two such observables and the implication of tagging on their evaluation is also discussed.

1 Introduction

In ultrarelativistic collisions of heavy nuclei a new deconfined state of matter, the quark-gluon plasma (QGP), is expected to be formed. The attenuation of the production rate of high transverse momentum particles in nucleus-nucleus (A+A) relative to the binary collision-scaled proton-proton (p+p) reactions has been proposed a long time ago as a signature of its formation [1]. The quenching of hadrons opposite a direct photon was studied [2] shortly after – a precursor of the modern tagged reconstructed jet observables [3, 4]. The production of a vector boson in association with jets is a powerful channel to probe the fundamental properties of Quantum Chromodynamics (QCD). Vector boson-tagged jets are also very well suited to studying the effects of the QGP. The tagging bosons escape the region of the hot dense medium unaffected and can be used to constrain the energy of the away-side parton shower. Previous studies of vector boson tagged jet production in heavy ion collisions have been performed in the framework of perturbative QCD [5, 6]. A Boltzmann transport model [7], an event generator JEWEL [8] and a hybrid strong/weak coupling model [9] have also been used. On the experimental side, isolated γ -tagged and Z^0 -tagged jets results in Pb+Pb collisions at $\sqrt{s_{NN}} = 5.02$ TeV have become available at the LHC from the ATLAS and CMS collaborations [10–12]. These new measurements call for improved theoretical calculations [13] to interpret the experimental data.

*e-mail: ivitev@lanl.gov

Vector-boson tagged jet cross section modification cannot be understood in isolation. The emergence of in-medium parton branching, qualitatively different from the one which determines the jet properties in $e^+ + e^-$, $e^- + p$, and $p + p$ collisions, also implies non-trivial modification of jet substructure observables. These include jet shapes, jet fragmentation functions and jet momentum sharing distributions. In these proceedings we also report on calculation of two of the above [14, 15] in a theoretical framework that uses the same fundamental input – the in-medium splitting functions.

2 Vector boson tagged jets

Let us now turn to the new theoretical results on V+jet modification at the LHC.

2.1 Proton collisions

Before we proceed to heavy ion collisions, it is important to validate the perturbative calculations in the simpler proton-proton reactions. We use Pythia 8 [16], which is a popular high energy phenomenology event generator. It utilizes leading-order perturbative QCD matrix elements+parton shower, combined with the Lund string model for hadronization and can describe well the main properties of a p+p event structure. This is shown in the left panel of figure 1 in comparison to Z^0 +jet measurements made by the CMS collaboration at $\sqrt{s} = 7$ TeV [17]. We also find good agreement with γ -tagged jet measurements, see [18].

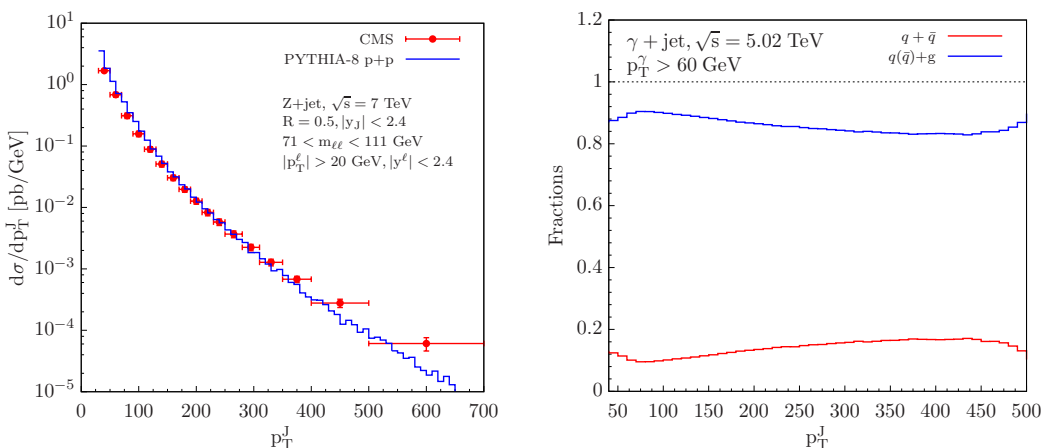


Figure 1. Left panel: Pythia 8 simulations are compared to CMS measurements of Z^0 +jet production in p+p collisions at the LHC at $\sqrt{s} = 7$ TeV. Results are shown differentially versus the jet p_T^J . Right panel: the fractional contributions of different subprocesses to isolated- γ +jet production cross sections in p+p collisions at $\sqrt{s} = 5.02$.

At leading order two dominant channels for V+jet are implemented in Pythia, namely $q + \bar{q} \rightarrow V + g$ and $q(\bar{q}) + g \rightarrow V + q(\bar{q})$. We have checked that the $g + g \rightarrow V + g$ channel contribution to the cross section is negligible. The utility of vector boson tagging beyond constraining the recoil jet energy is that it can be used for flavor selection. We have shown this in the right panel of figure 1 on the example of photon-tagged jets at $\sqrt{s} = 5.1$ TeV. Results indicate that γ tagging or Z^0 tagging preferentially selects quark jets and at high p_T the sample has $\sim 80\%$ purity. The accumulated statistics allows us to use the fully differential $d\sigma/dp_T^V dp_T^J$ to high transverse momenta of the jet and the vector boson.

2.2 Heavy ion collisions collisions

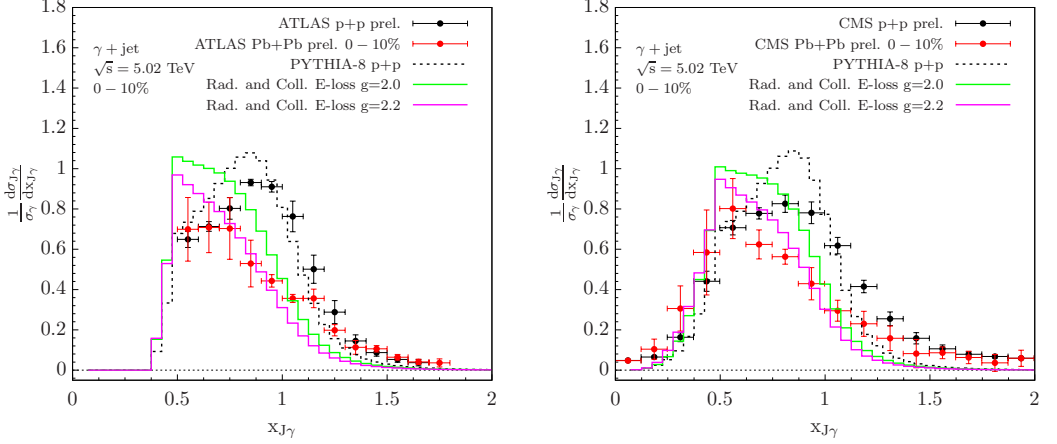


Figure 2. Left panel: The isolated photon-tagged jet asymmetry distributions are shown and compared to ATLAS data in central Pb+Pb collisions at the LHC [12]. The transverse momenta for the isolated photon and the jet are $p_T^\gamma > 60$ GeV and $p_T^J > 30$ GeV, respectively. The jet radius parameter is $R = 0.4$. Right panel: same comparison but for CMS data [11] with jet radius $R = 0.3$.

We start from the baseline p+p cross sections obtained from Pythia. However, in heavy ion collisions we take into account the reaction geometry, the hydrodynamic expansion of the QGP and radiative and collisional energy losses. At an impact parameter $|\mathbf{b}_\perp|$ we evaluate the cross sections as

$$\begin{aligned} \frac{d\sigma^{AA}(|\mathbf{b}_\perp|)}{dp_T^V dp_T^J} &= \int d^2 \mathbf{s}_\perp T_A \left(\mathbf{s}_\perp - \frac{\mathbf{b}_\perp}{2} \right) T_A \left(\mathbf{s}_\perp + \frac{\mathbf{b}_\perp}{2} \right) \sum_{q,g} \int_0^1 d\epsilon \frac{P_{q,g}(\epsilon; s_\perp, |\mathbf{b}_\perp|)}{1 - f_{q,g}^{\text{loss}}(R; s_\perp, |\mathbf{b}_\perp|) \epsilon} \\ &\times \frac{d\sigma_{q,g}^{NN}(p_T^V, p_T^J / (1 - f_{q,g}^{\text{loss}}(R; s_\perp, |\mathbf{b}_\perp|) \epsilon))}{dp_T^J dp_T^V}. \end{aligned} \quad (1)$$

Here hard production follows the binary collision density $\propto T_A(\mathbf{s}_\perp - \mathbf{b}_\perp/2) T_A(\mathbf{s}_\perp + \mathbf{b}_\perp/2)$ and we use an optical Glauber model with inelastic nucleon-nucleon scattering cross sections $\sigma_{\text{in}} = 70$ mb at $\sqrt{s_{NN}} = 5.02$ TeV. In equation 1 $P_{q,g}(\epsilon)$ is the probability that a fraction ϵ of the parent parton is redistributed through medium-induced multiple gluon emission. We employ the soft gluon emission limit of the of the SCET_G [19, 20] splitting functions [21, 22] used to predict inclusive light hadron attenuation in heavy ion collisions [23, 24]. We note that the spectrum coincides with the original energy loss derivation of [25]. Of course, only the fraction of the energy that is redistributed outside of the jet cone leads to jet cross section suppression. We evaluate this fraction as follows

$$f_{q,g}^{\text{loss}}(R; \text{rad} + \text{coll}) = 1 - \left(\int_0^R dr \int_{\omega_{\min}}^E d\omega \frac{dN_{q,g}^g(\omega, r)}{d\omega dr} \right) / \left(\int_0^{R_{\max}} dr \int_0^E d\omega \frac{dN_{q,g}^g(\omega, r)}{d\omega dr} \right). \quad (2)$$

Equation 2 treats radiative and collisional energy losses on the same footing since for convenience we have express this total collisional energy loss as an integral over the spectrum of the medium induced

gluons

$$\Delta E_{q,g}^{\text{coll}}(\text{tot.}) = \sum_{i=1}^{N_{q,g}^{\text{tot. partons}}} \int_{z_i}^{\infty} \frac{d\Delta E_i^{\text{coll}}}{d\Delta z} d\Delta z, \quad \Delta E_{q,g}^{\text{coll}}(\text{tot.}) = \int_0^{\omega_{\min}} d\omega \int_0^{R_{\max}} dr \omega \frac{dN_{q,g}^g(\omega, r)}{d\omega dr}. \quad (3)$$

To lowest leading order in perturbative QCD, the transverse momentum of the vector boson is balanced by the transverse momentum of the jet, $p_T^V = p_T^J$. Higher order processes and the development of parton showers introduce deviations from this equality. The exact differential distribution of $d\sigma/dp_T^V dp_T^J$ is also affected by the jet reconstruction algorithm, jet radius choice, experimental cuts, and detector resolution. Nevertheless, the quenching and shift of this distribution to smaller values of p_T^J are currently the best proxies for jet energy loss. To study these effects one introduces the transverse momentum imbalance $x_{JV} = p_T^J/p_T^V$. distribution can be obtained from the double differential distribution of V+jet cross section

$$\frac{d\sigma}{dx_{JV}} = \int_{p_T^{J,\min}}^{p_T^{J,\max}} dp_T^J \frac{p_T^J}{x_{JV}^2} \frac{d\sigma(p_T^V = p_T^J/x_{JV}, p_T^J)}{dp_T^V dp_T^J}, \quad (4)$$

where $p_T^{J,\min}$ and $p_T^{J,\max}$ are matched to the desired cuts of the experimental measurements.

In figure 2 we plot the normalized momentum imbalance distributions for the γ +jet final state in p+p and Pb+Pb collisions at the LHC. Comparison of the calculations to the ATLAS measurements [12] are shown in the left panel. Comparison of the calculations to the CMS measurements [11] are shown in the right panel. The black dashed histogram shows Pythia 8 simulations for the the x_{JZ} distribution in p+p collisions. The black solid points represent the ATLAS and CMS results in the elementary nucleon-nucleon collisions. It can be seen that the x_{JZ} distribution from Pythia 8 simulation is narrower than the one measured by the experiments. Part of the reason is that experimental measurements are not unfolded for detector resolution effects. This can be seen on the example of Z^0 +jet momentum imbalance by CMS in Ref. for the p+p reference [10]. We applied the same smearing functions that experiment applies to Monte Carlo simulations, to our differential p_T distributions for p+p and Pb+Pb collisions. We got broader x_{JZ} distributions and better agreement between the curves to the data points in both p+p and Pb+Pb collisions. In figure 2 we also give results of our theoretical calculations in Pb+Pb collisions. The green and magenta histograms correspond to jet-medium coupling strengths $g = 2.0$ and $g = 2.2$, respectively, values that have worked well in describing the light inclusive hadron [23, 24] and jet suppression data [18] at the LHC. Our results include both medium-induced radiative energy loss and the parton showers energy dissipation of through collisional processes in the QGP. We see the downshift of x_{JV} , which quantitatively agrees with measurements in terms of the difference between p+p and Pb+Pb collisions.

Another often studied observable sensitive to nuclear modification effects in V+jet systems is I_{AA} . It is defined as ratio of the differential cross section for tagged jets in A+A collisions to the binary collision scaled p+p result. We present I_{AA} in isolated γ +jet production at the LHC in 0 – 30% Pb+Pb collisions in figure 3. By comparing to CMS experimental data, we find that our theoretical results agree with data for a wide kinematic range. In each p_T^γ window, the jet energy loss effects are shown in four cases and coded with different colors - two different jet-medium coupling strengths and simulations where collisional energy loss effects are either included or excluded. Figure 3 illustrates that there is kinematic sensitivity of I_{AA} and the largest suppression is observed along the diagonal region of the transverse momenta of the trigger γ and the recoil jet. The reason for this behavior in $p_T^\gamma \approx p_T^J$ arises from the steeper falling cross section in that region. The cross section is suppressed in the region $p_T^J > p_T^\gamma$ and is enhanced for $p_T^J < p_T^\gamma$. This is characteristic of in-medium tagged-jets as discussed in previous works [5, 6].

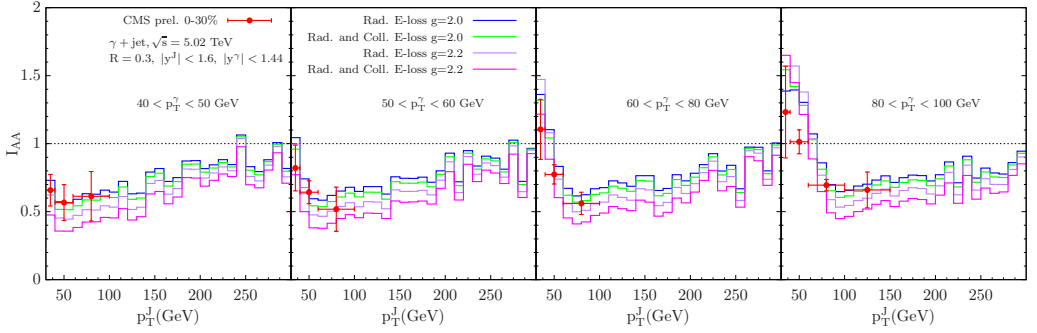


Figure 3. Theoretical results for I_{AA} compared to CMS data with several different transverse momentum cuts. Shown are four different scenarios for energy loss effects with and without collisional energy loss and with couplings between the jet and the medium $g = 2.0$ and $g = 2.2$

Table 1. Theoretical predictions for the difference of the average x_{JZ} between p+p and Pb+Pb 0 – 30% central collisions. The center of mass energy is chosen to be $\sqrt{s} = 5.02$ TeV, the transverse momentum cut for the recoil jet is $p_T^J > 30$ GeV, and data is from CMS [10].

	$\Delta\langle x_{JZ} \rangle$			
p_T^Z (GeV)	40 – 50	50 – 60	60 – 80	80 – 120
CMS [10]	0.061 ± 0.059	0.123 ± 0.051	0.124 ± 0.052	0.068 ± 0.042
Rad. + Coll. $g = 2.0$	0.022	0.050	0.075	0.086
Rad. + Coll. $g = 2.2$	0.024	0.058	0.093	0.119

To better quantify the shift toward smaller values of the x_{JV} distribution, we introduce the mean value of x_{JV} and the difference between proton and heavy ion collisions

$$\langle x_{JV} \rangle = \left(\int dx_{JV} x_{JV} \frac{d\sigma}{dx_{JV}} \right) / \left(\int dx_{JV} \frac{d\sigma}{dx_{JV}} \right), \quad \Delta\langle x_{JV} \rangle = \langle x_{JV} \rangle_{pp} - \langle x_{JV} \rangle_{PbPb}. \quad (5)$$

The calculated downshifts of the x_{JV} distribution is consistent with the experimental data [10] within the measurement uncertainties for different p_T^Z cuts. More importantly, when the downshift is recast into energy loss of the jet recoiling against the Z^0 boson, we find that this experimentally determined jet energy loss corresponds within 15% accuracy to the theoretically computed energy redistributed out of the jet cone due to in-medium interactions of the parton shower

3 Jet substructure

We now turn to the question of jet substructure in heavy ion collisions.

3.1 Momentum sharing distributions

A new jet substructure observable, called the groomed momentum sharing, was recently proposed and studied using the soft drop jet grooming procedure [26, 27]. It is sensitive to the hard branching in the jet formation and is controlled by the leading-order Altarelli-Parisi splitting functions. Having reconstructed a jet of radius R using the anti- k_T algorithm, one can recluster the jet using the

Cambridge/Aachen algorithm and go through the branching history, dropping the soft branch until $z_{cut} < \min(p_{T_1}, p_{T_2})/(p_{T_1} + p_{T_2}) \equiv z_g$. One can also demand that the angular separation between the two branches is greater than an angular scale Δ , $\Delta < \Delta R_{12} \equiv r_g$. By selecting the groomed jet radius r_g , one gains access to the momentum sharing distribution $p(z_g)$ at different splitting angles, and can study the angular distribution ρz_g . We evaluate

$$p_i(z_g) = \frac{\int_{k_\Delta}^{k_R} dk_\perp \bar{\mathcal{P}}_i(z_g, k_\perp)}{\int_{z_{cut}}^{1/2} dx \int_{k_\Delta}^{k_R} dk_\perp \bar{\mathcal{P}}_i(x, k_\perp)} \quad \text{where} \quad k_\Delta = \omega x(1-x) \tan \frac{\Delta}{2}, \quad k_R = \omega x(1-x) \tan \frac{R}{2}. \quad (6)$$

The symmetrized splitting functions are $\bar{\mathcal{P}}_i(x, k_\perp) = \sum_{j,l} [\mathcal{P}_{i \rightarrow jl}(x, k_\perp) + \mathcal{P}_{i \rightarrow jl}(1-x, k_\perp)]$ and in the presence of a QGP, $\mathcal{P}_{i \rightarrow jl}(x, k_\perp) = \mathcal{P}_{i \rightarrow jl}^{vac}(x, k_\perp) + \mathcal{P}_{i \rightarrow jl}^{med}(x, k_\perp)$. The medium-induced component were calculated using soft-collinear effective theory with Glauber gluon interactions (SCET_G) [19–22].

Analytic and numerical considerations show that for $x < 1/2$ in Eq. (6) medium-induced component goes as $1/x^2$ [21]. We thus expect that the momentum sharing distribution will be enhanced at small z_g and suppressed near $z_g = 1/2$. Numerical results [15] are shown in the left panel of figure 4 and compared to preliminary CMS data. The parameters for the groomed soft-dropped jets are $\beta = 0$, $z_{cut} = 0.1$ and $\Delta R_{12} > 0.1$. We estimate the theoretical uncertainty by varying the jet-medium coupling $g = 2.0 \pm 0.2$. Analogous calculations in the soft gluon emission limit have also been presented [28, 29]. The finalized CMS data [30] changed significantly and is now in excellent agreement with the theoretical predictions of [15]. This signifies the need for reliable pQCD theory to interpret heavy ion data and motivates further studies, such as the groomed jet radius distribution. The anticipated modification is shown in the right panel of figure 4 for several different p_T bins. The peak of this distribution reflects the medium enhancement of large-angle splitting for hard branching processes.

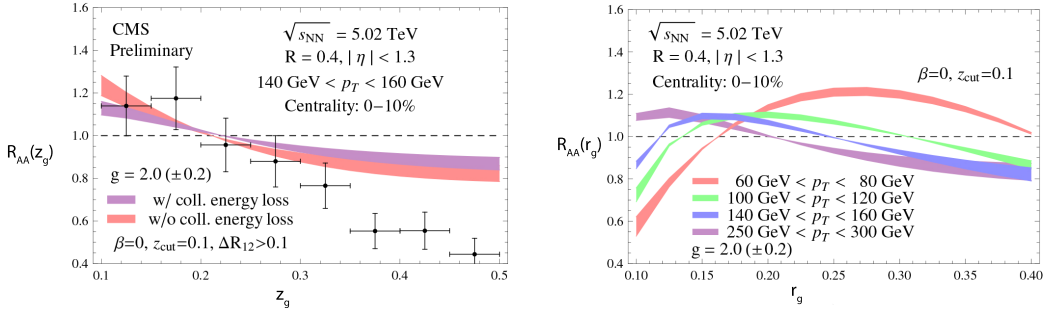


Figure 4. Left panel: theoretical calculations compared to preliminary CMS data for the ratio of momentum sharing distributions of inclusive anti- k_T $R = 0.4$ jets in central Pb+Pb and p+p collisions at $\sqrt{s_{NN}} = 5.02$ TeV. Purple and red bands correspond to calculations with and without the implementation of collisional energy loss. Right panel: the groomed jet radius modification of inclusive jets in proton-proton and central lead-lead collisions at $\sqrt{s_{NN}} = 5.02$ TeV. We show results for four p_T bins with $60 \text{ GeV} < p_T < 80 \text{ GeV}$, $100 \text{ GeV} < p_T < 120 \text{ GeV}$, $140 \text{ GeV} < p_T < 160 \text{ GeV}$ and $250 \text{ GeV} < p_T < 300 \text{ GeV}$.

3.2 Jet shapes

A complementary observable, in fact the first one that has been quantitatively studied in relation to inclusive jet modification in heavy ion collisions [31], is the jet shape. The jet shape [32] describes

the transverse energy profile inside a jet of size R , reconstructed using a jet algorithm. The integral jet shape, defined as the fraction of the transverse energy E_T of the jet within a subcone of size r , reads

$$\Psi_J(r) = \frac{\sum_{i, d_{i\tilde{n}} < r} E_T^i}{\sum_{i, d_{i\tilde{n}} < R} E_T^i}, \quad \text{where } d_{i\tilde{n}} = \sqrt{(\eta_i - \eta_{jet})^2 + (\phi_i - \phi_{jet})^2} \quad (7)$$

In heavy ion collisions, the differential jet shape $\rho(r) = \frac{d}{dr}\Psi(r)$ is more sensitive to in-medium modification, which can be presented as $M_\rho(r) = \rho^{AA}(r)/\rho^{pp}(r)$.

A robust theory of jets in high energy nuclear collisions should describe, within its regions of applicability, the modification of the jet shapes in addition to the attenuation of the inclusive and dijet cross sections [33]. We have taken the opportunity to address this key issue in [14], building upon an improved description of $\rho^{pp}(r)$ [34] in the framework of SCET [35]. Our results are shown in figure 5, where we impose cuts on the jet transverse momentum $p_T > 100$ GeV and pseudo-rapidity $0.3 < |\eta| < 2.0$ of jets. We take the coupling between the jet and the QGP to be $g = 2$ and the theoretical uncertainty is estimated by varying the jet energy scales $\frac{1}{2}\mu_{j_R} < \mu < 2\mu_{j_R}$ in the calculations. Several contributions to the jet shape modification are shown. The blue band corresponds to the initial state cold nuclear matter effects. The effect of quenching is added in the red band, which enhances the fraction of quark jets that are narrower than the gluon jets. The green band correspond to the full calculation the green band is the full calculation and adds the event-by-event broadening. It is evident that the jet shape modification is highly nontrivial and we find that it compares well to the data from CMS [36] with jet radius $R = 0.3$.

We also present predictions for the modification of jet shapes in Pb+Pb collisions at $\sqrt{s_{NN}} \approx 5.1$ TeV, focusing on the differences between inclusive jets and tagged jets. The results are given in the right panel of figure 5. While $M_\rho(r)$ for jet shapes at $\sqrt{s_{NN}} \approx 5.10$ TeV is similar to the one at $\sqrt{s_{NN}} = 2.76$ TeV, the modification of γ -tagged jet shapes is more pronounced. Photon-tagged jet shapes also exhibit a different modification pattern. Since photon tagging selects primarily quark jets, see previous section, the quenching does not play a significant role and the enhancement toward the periphery of jets is more pronounced. Other examples of jet shape calculations include [37, 38].

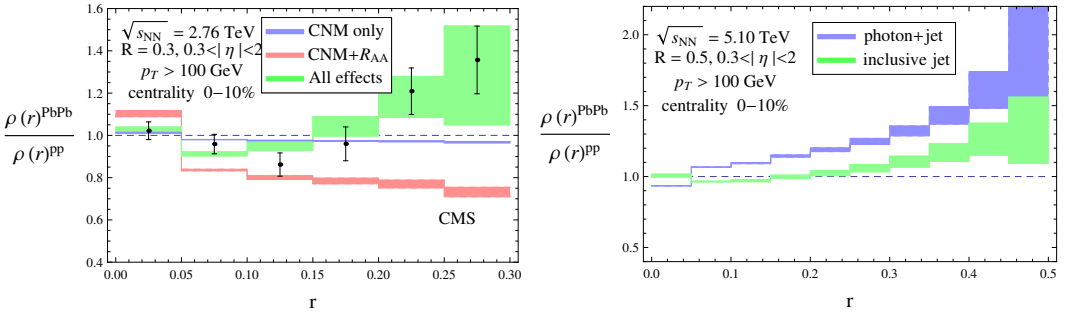


Figure 5. Left panel: the modification of differential jet shapes of inclusive jets in Pb+Pb central collisions at $\sqrt{s_{NN}} = 2.76$ TeV at the LHC. Different contributions are shown and compared to CMS data. Right panel: theoretical predictions for the differential jet shape modification of inclusive and photon-tagged jets with $R = 0.3$ in Pb+Pb collisions at $\sqrt{s_{NN}} \approx 5.1$ TeV at the LHC.

4 Conclusions

To summarize, in these proceedings we reported results on vector boson-tagged jet production [13] and jet substructure [14, 15] in heavy ion collisions. We presented a new study of γ -tagged or Z^0 -tagged jets in Pb+Pb reactions at the LHC a center-of-mass energy per nucleon pair of 5.02 TeV. Good description of the mean momentum imbalance shift $\Delta\langle x_{JV} \rangle$ and the tagged jet nuclear modification factor I_{AA} can be achieved in a framework that accounts for the QGP-induced radiative and collisional dissipation of the parton shower energy outside of the jet cone. At the heart of such quenching effects for reconstructed jets is the in-medium parton branching – qualitatively and quantitatively different from the standard Altarelli-Parisi splittings in the vacuum. New techniques, such as the study of the subjet momentum sharing distributions, have recently emerged and can probe precisely this parton dynamics. We also presented the first calculation of this subjet momentum sharing distribution in heavy ion collisions. The validation of the theoretical predictions by recent experimental measurements complements our earlier results on the modification of the jet shape and corroborates the perturbative picture of parton shower modification in the medium. Finally, we found that the modification of γ -tagged and Z^0 -tagged jet substructure can differ substantially from the corresponding substructure modification of inclusive jets. Future experimental results for such observables can significantly enhance our understanding of in-medium QCD dynamics.

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