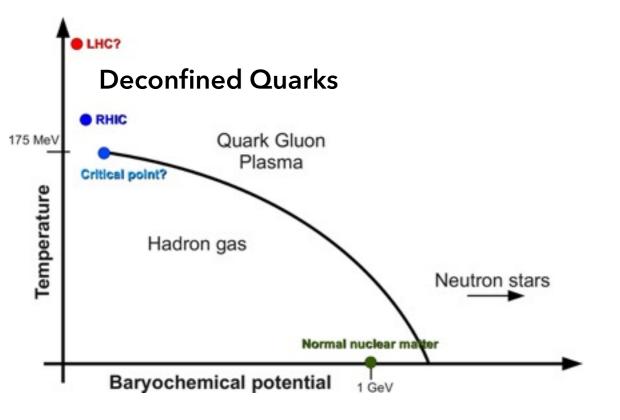




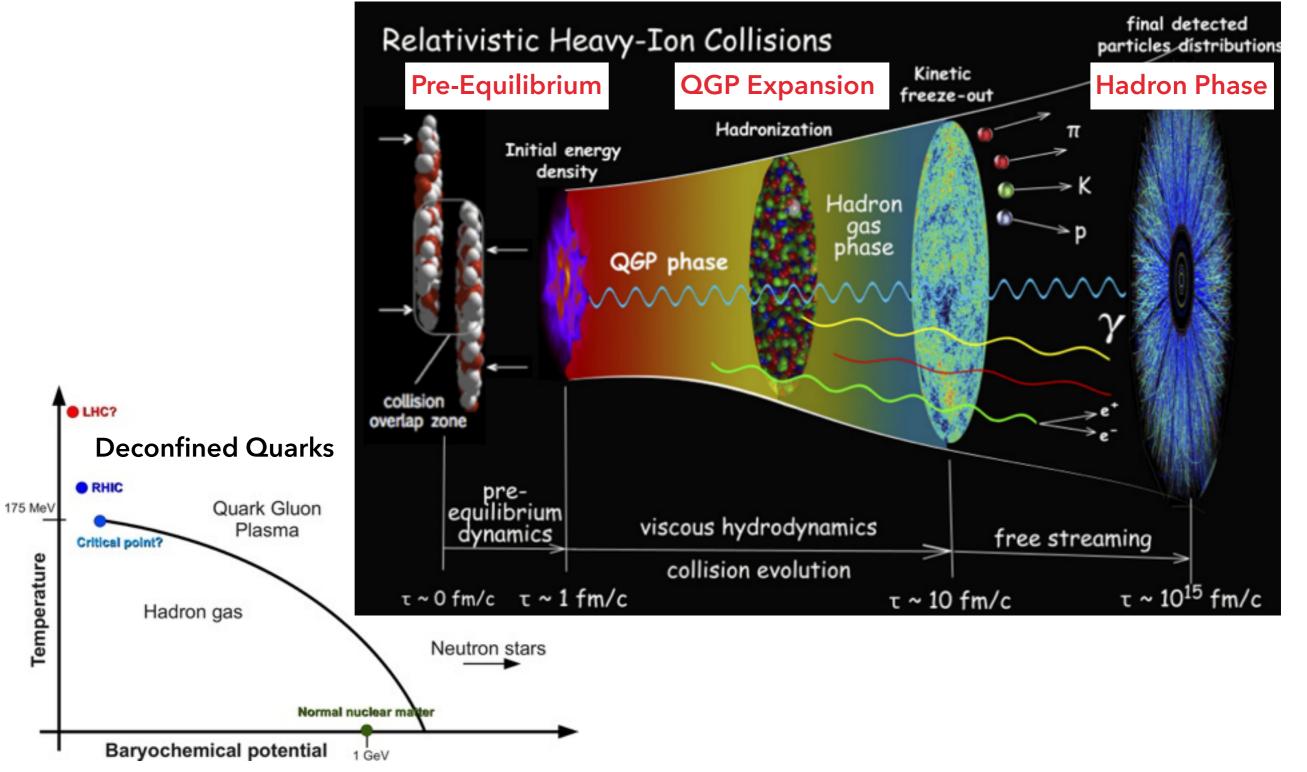


A+A COLLISIONS ———— SMALL SYSTEM

QUARK GLUON PLASMA (QGP)

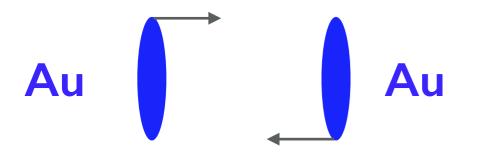


QUARK GLUON PLASMA EVOLUTION

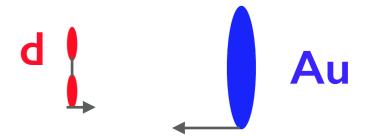


WHY ARE SMALL SYSTEMS INTERESTING?

Au+Au collisions were thought to be the necessary for the generation of QGP

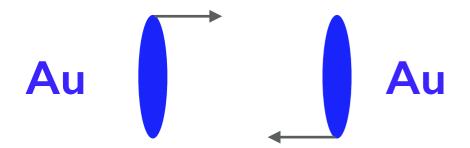


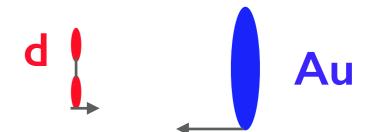
Small systems like d+Au were thought of as the control test to measure cold nuclear matter effects.



WHY ARE SMALL SYSTEMS INTERESTING?

Au+Au collisions were thought to be the necessary for the generation of QGP Small systems like d+Au were thought of as the control test to measure cold nuclear matter effects.

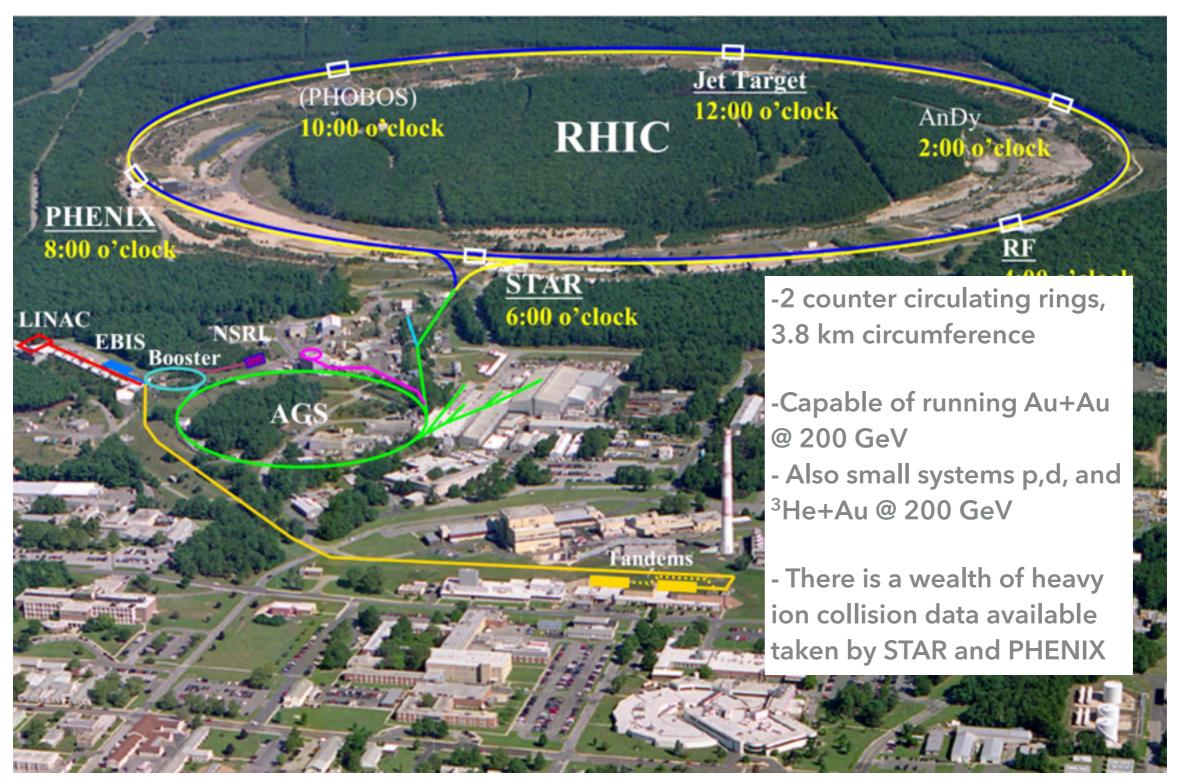




However, recent measurements of flow and jet quenching in small systems have yielded surprising results.

Small systems allow for control over the initial collision geometry.

RELATIVISTIC HEAVY ION COLLIDER AT BNL



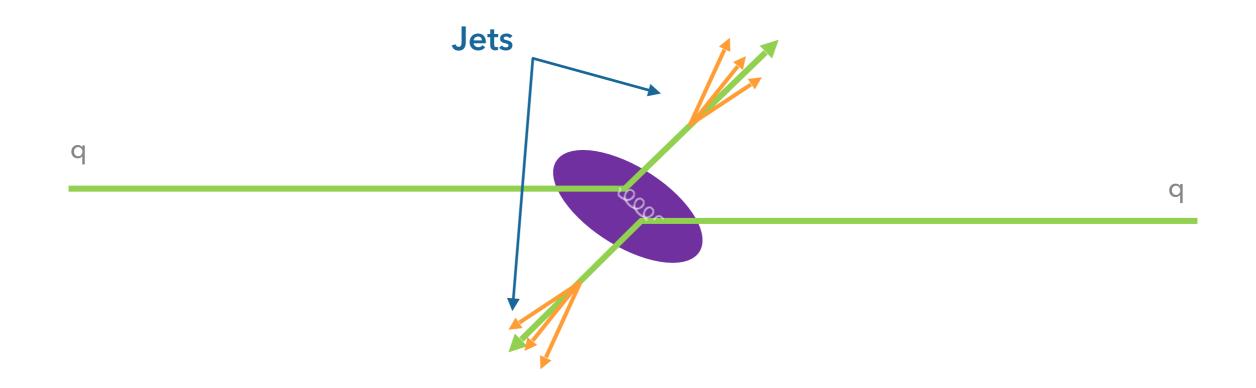




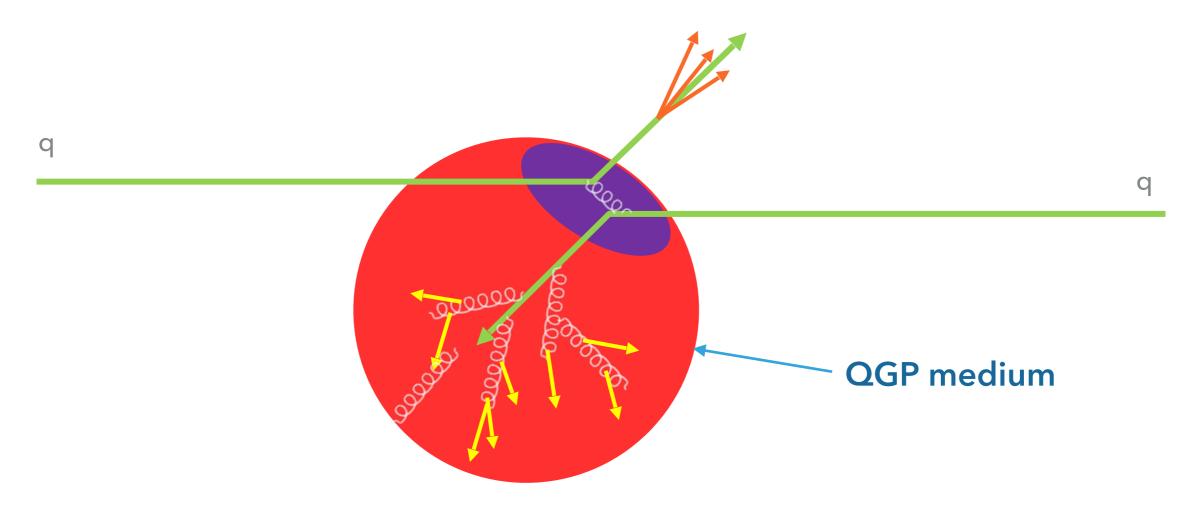
A+A COLLISIONS

SMALL SYSTEM

HARD SCATTERING



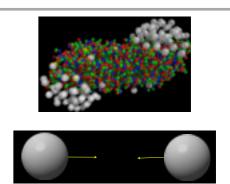
THE QGP IS OPAQUE



Energy loss via Gluon Bremsstrahlung

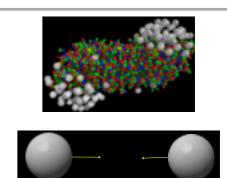
Suppression in jets.

 $\frac{\text{Physics in the Medium}}{\text{Physics in the Vacuum}}$



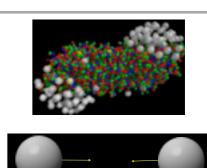
$$R_{AA} = \frac{1}{\langle N_{coll} \rangle} \frac{dN_{AA}/dp_T}{dN_{pp}/dp_T}$$

 $R_{AA} = \frac{1}{\langle N_{coll} \rangle} \frac{dN_{AA}/dp_T}{dN_{pp}/dp_T} \sim \frac{\text{Physics in the Medium}}{\text{Physics in the Vacuum}}$



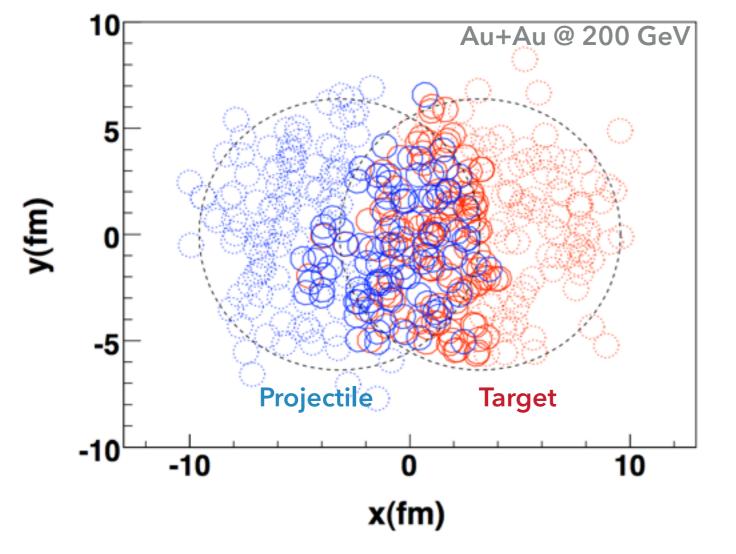
Number of binary collisions

$$R_{AA} = \frac{1}{\langle N_{coll} \rangle} \frac{dN_{AA}/dp_T}{dN_{pp}/dp_T} \sim \frac{\text{Physics in the Medium}}{\text{Physics in the Vacuum}}$$



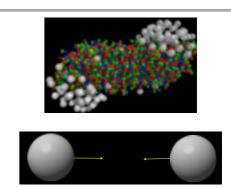
Number of binary collisions

Glauber Monte Carlo



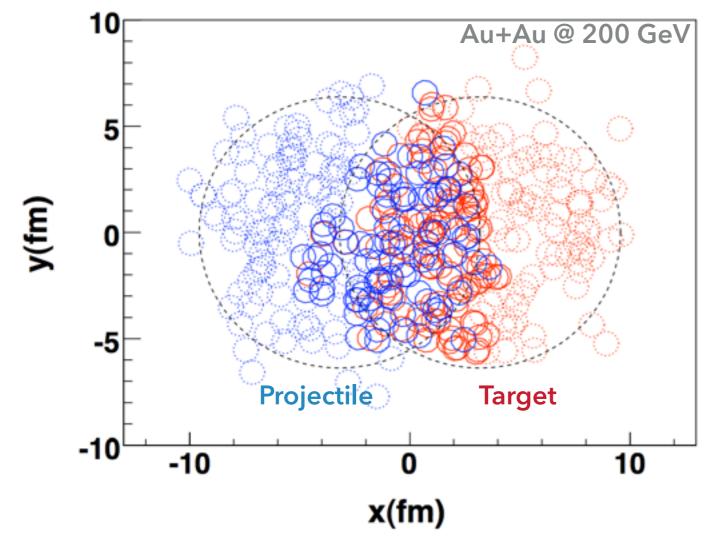
$$R_{AA} = \frac{1}{\langle N_{coll} \rangle} \frac{dN_{AA}/dp_T}{dN_{pp}/dp_T}$$

 $R_{AA} = \frac{1}{\langle N_{coll} \rangle} \frac{dN_{AA}/dp_T}{dN_{pp}/dp_T} \sim \frac{\text{Physics in the Medium}}{\text{Physics in the Vacuum}}$

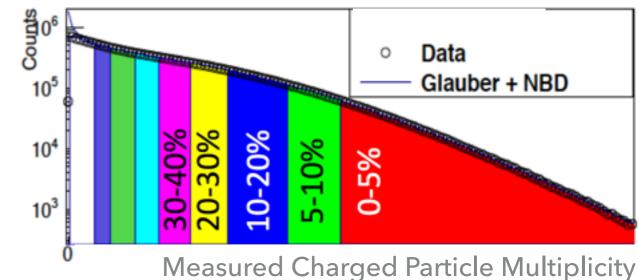


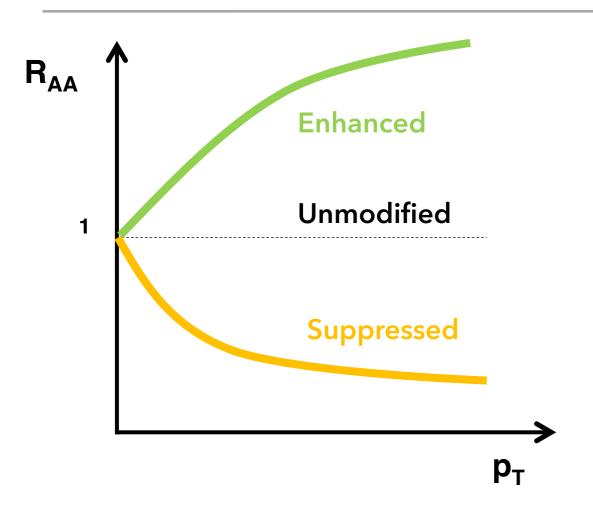
Number of binary collisions

Glauber Monte Carlo

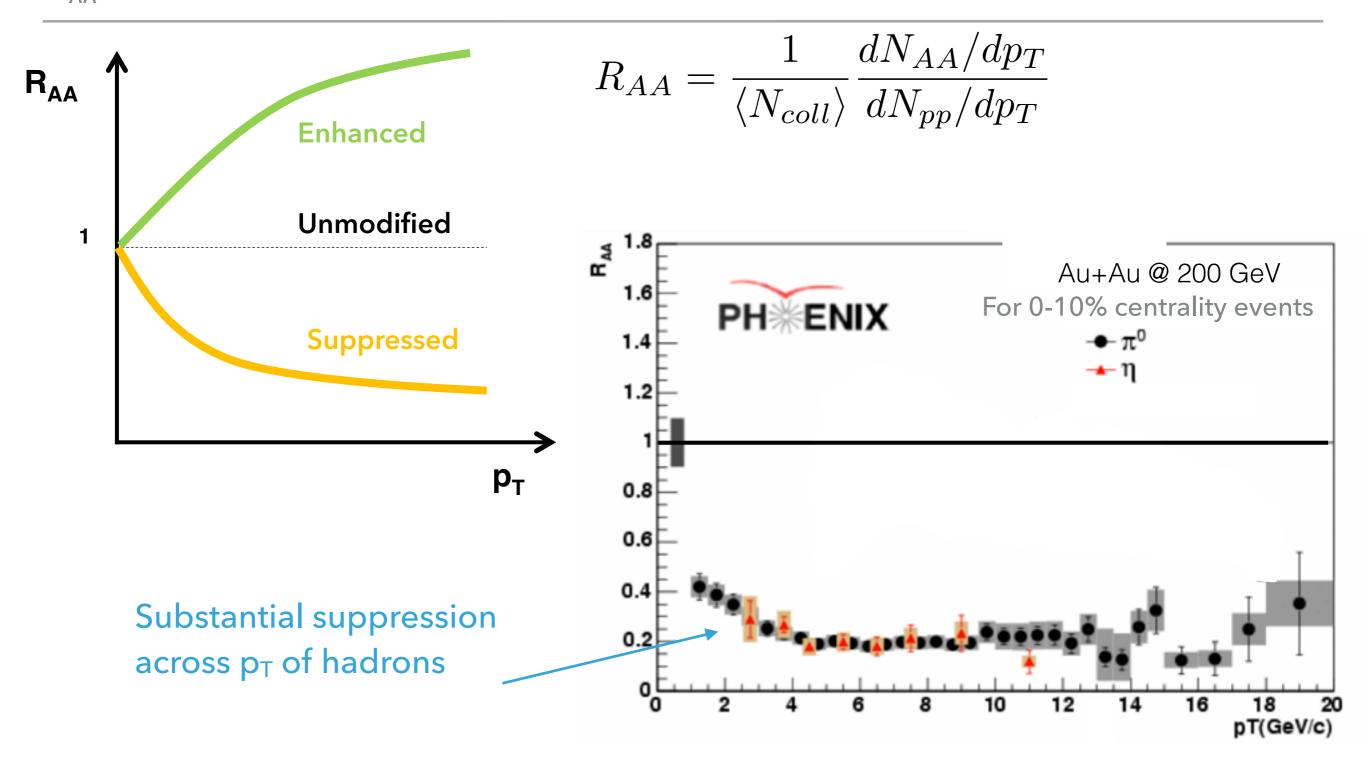


And we can relate Ncoll to charged particle multiplicity to determine centrality classes: 0 -100% (0% is most central)





$$R_{AA} = \frac{1}{\langle N_{coll} \rangle} \frac{dN_{AA}/dp_T}{dN_{pp}/dp_T}$$



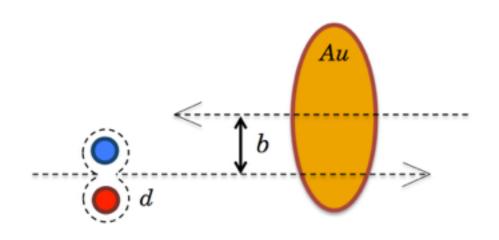
This measurement is consistent with jet quenching due to the medium.

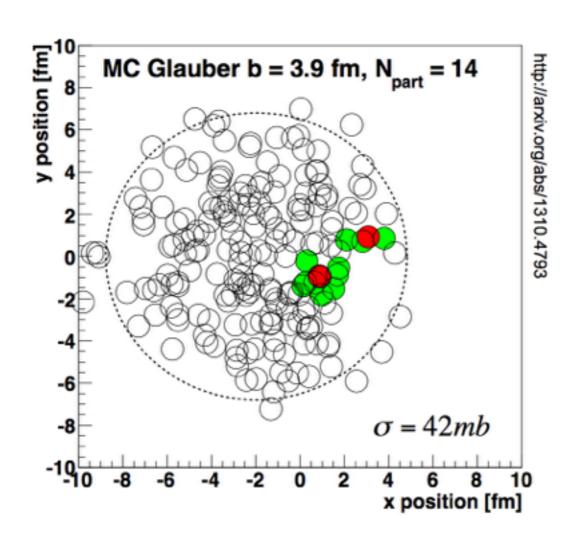




A+A COLLISIONS

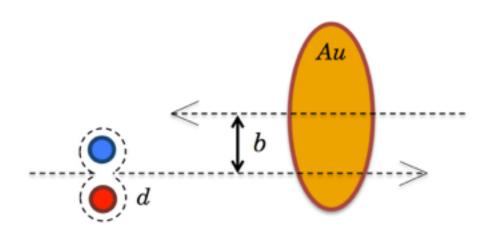
SMALL SYSTEM

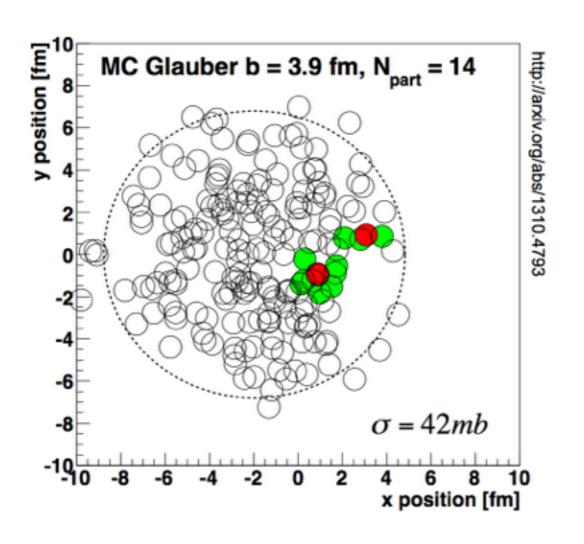




$$R_{\mathrm{d}A} = \frac{1}{\langle N_{coll} \rangle} \frac{dN_{\mathrm{d}A}/dp_T}{dN_{pp}/dp_T}$$

The number of binary collisions in d+Au is much lower than in Au+Au.



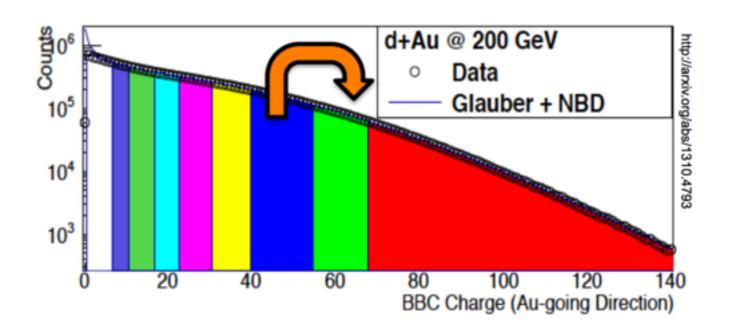


$$R_{\mathrm{d}A} = \frac{1}{\langle N_{coll} \rangle} \frac{dN_{\mathrm{d}A}/dp_T}{dN_{pp}/dp_T}$$

The number of binary collisions in d+Au is much lower than in Au+Au.

NOT ALL BINARY COLLISIONS ARE CREATED EQUAL

Hard scattering collisions will bias the centrality of events upwards.

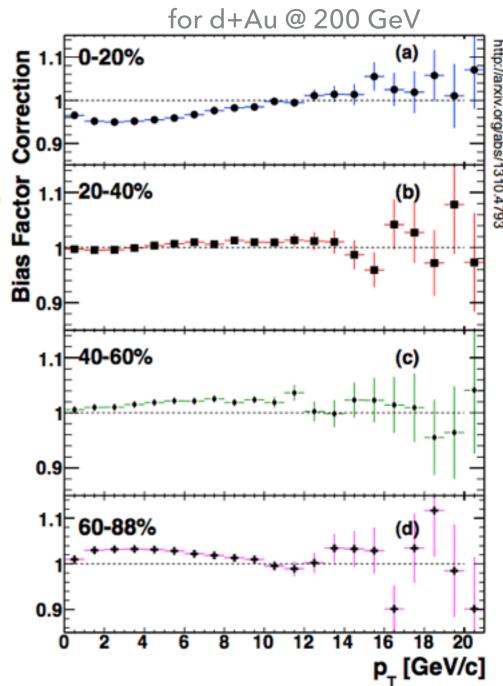


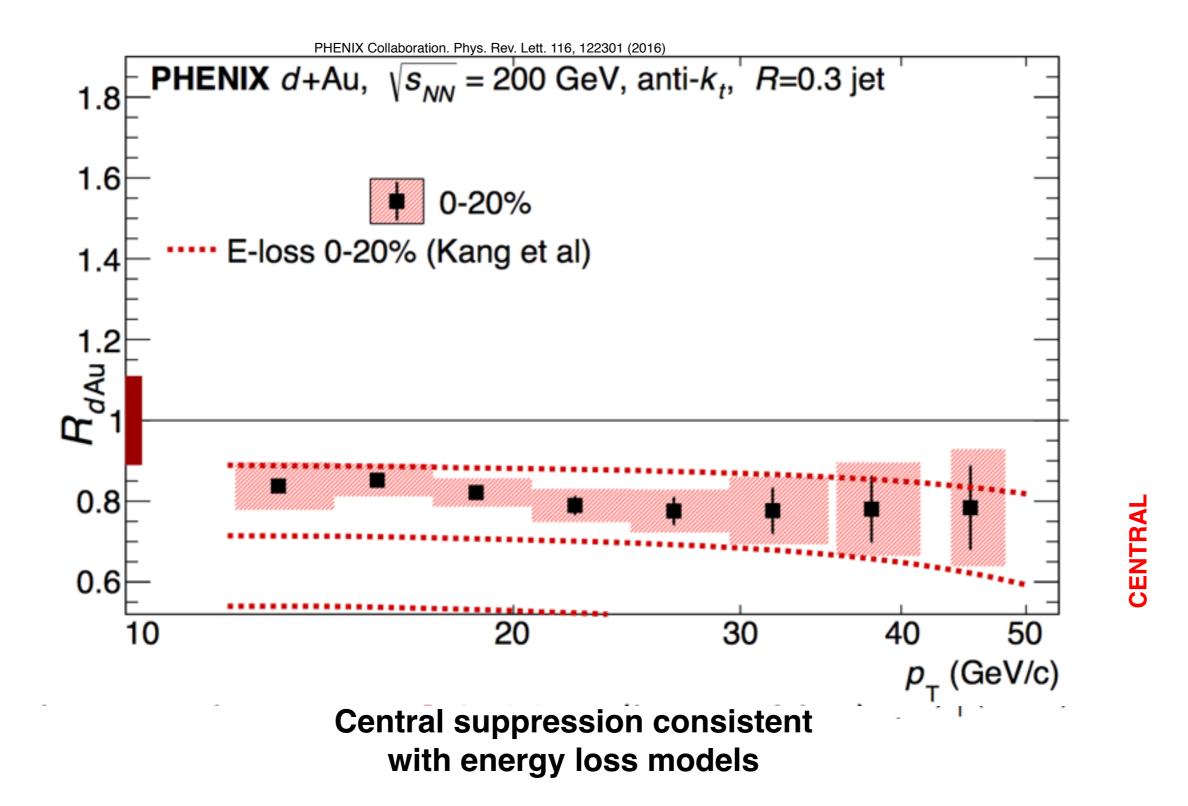
 Using the MC Glauber, we can calculate the centrality dependent bias factors

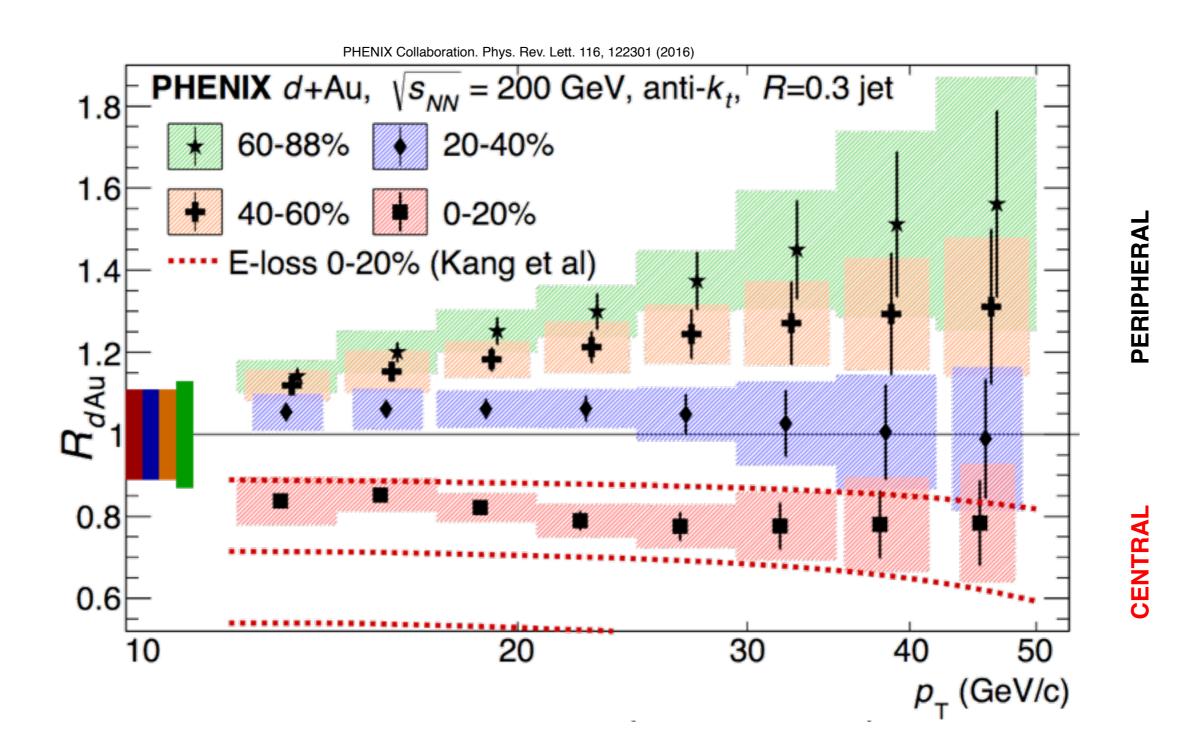
	Glauber Model	
Centrality	Bias Correction Factor	
0-20%	0.94 ± 0.01	
20-40%	1.00 ± 0.01	
40-60%	1.03 ± 0.02	
60-88%	1.03 ± 0.06	

- Using the MC Glauber, we can calculate the centrality dependent bias factors
- We can go farther to calculate the pT dependent bias factors using HIJING (Heavy Ion Jet INteraction Generator)

	Glauber Model	HIJING	
Centrality	Bias Correction Factor	Mean Bias Factor 1< p _⊤ <5 GeV/c	
0-20%	0.94 ± 0.01	0.951 ± 0.001	
20-40%	1.00 ± 0.01	0.996 ± 0.001	
40-60%	1.03 ± 0.02	1.010 ± 0.001	
60-88%	1.03 ± 0.06	1.030 ± 0.001	

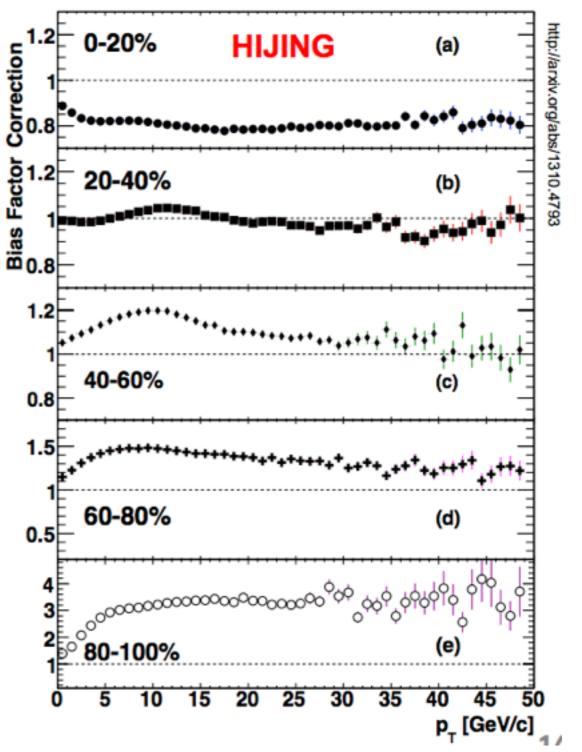




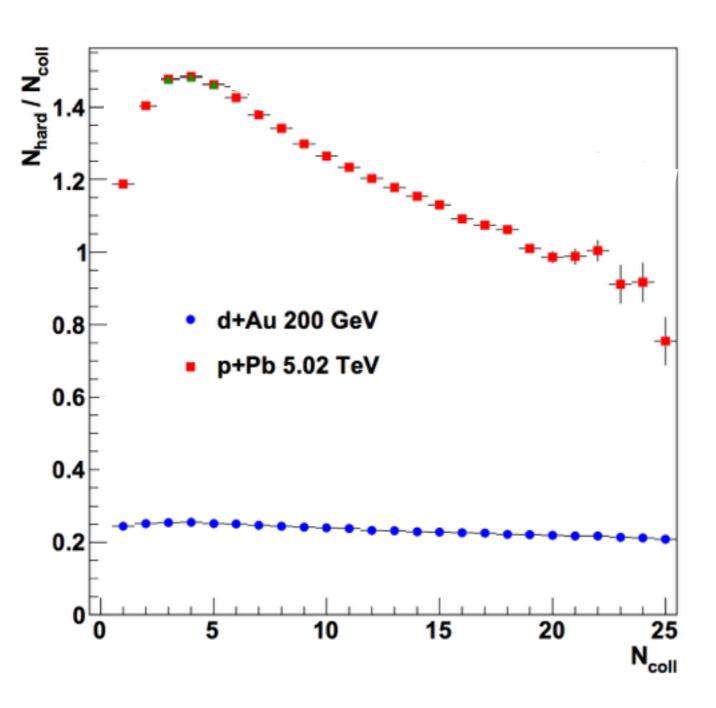


 We can calculate the bias factors for LHC p+Pb 5.02 TeV events.

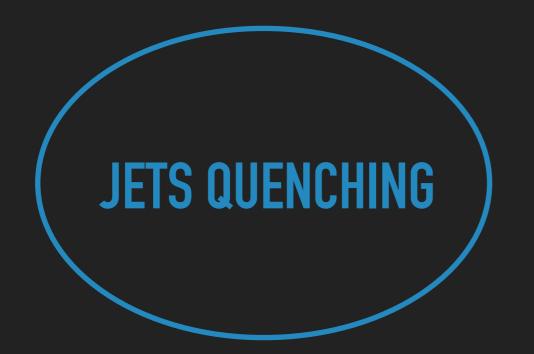
The centrality bias effect is much larger at LHC energies, MPI could play a role.



Theodore Koblesky MPI 2016



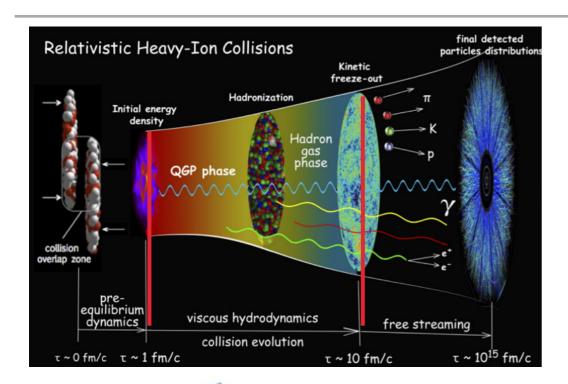
- Using HIJING, we can compute the number of hard scatterings per nucleon nucleon collision
- There is nearly an order of magnitude difference between the LHC and RHIC.





A+A COLLISIONS

SMALL SYSTEM



- The medium becomes locally equilibrated
- Initial state geometric anisotropy gets translated into final state momentum anisotropy.

 We can measure flow by looking at the long range angular correlations in the spray of particles

$$rac{dN}{d\phi} \propto 1 + \sum_{n=1} 2 \, v_n \, \cos(n [\phi - \Psi_n]),$$

VN are Flow Coefficients

$$v_{\rm N} = \frac{\langle \sum \cos \mathbb{N} (\phi - \Psi_{\rm N}) \rangle}{\text{Res}(\Psi_{\rm N})}$$

 $\Psi_{\scriptscriptstyle N}$ is the generalized participant Event Plane

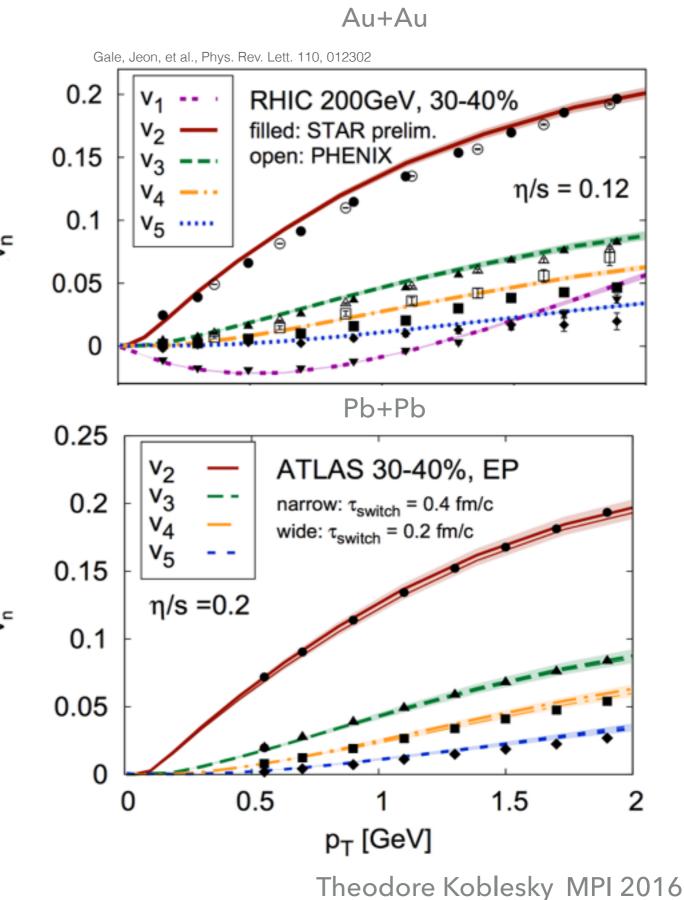
$$rac{dN}{d\phi} \propto 1 + \sum_{n=1}^{\infty} 2 \, v_n \, \cos(n [\phi - \Psi_n]),$$

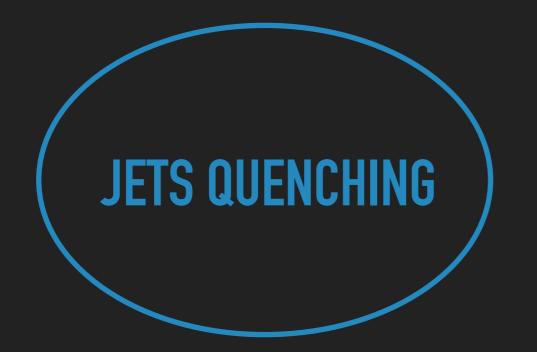
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 $\Psi_{\scriptscriptstyle N}$ is the generalized participant Event Plane

Hydrodynamics describes the data at both energies up to the 5th harmonic order.





COLLECTIVE FLOW

A+A COLLISIONS

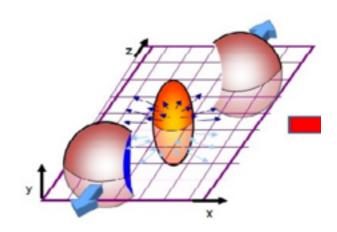
SMALL SYSTEM

p+Au d+Au ³He+Au 200 GeV

0-5% central p+Au d+Au ³He+Au 200 GeV

	$p+\mathrm{Au}$	$d+\mathrm{Au}$	³ He+Au
$\langle N_{ m coll} angle$	9.7 ± 0.6	18.1 ± 1.2	26.1 ± 2.0
$\langle N_{ m part} angle$	10.7 ± 0.6	17.8 ± 1.2	25.1 ± 1.6
Glauber $\langle \varepsilon_2 \rangle$	0.23 ± 0.01	0.54 ± 0.04	0.50 ± 0.02

 $\boldsymbol{\epsilon}_{\mathbf{2}}$ is the second order initial collision eccentricity



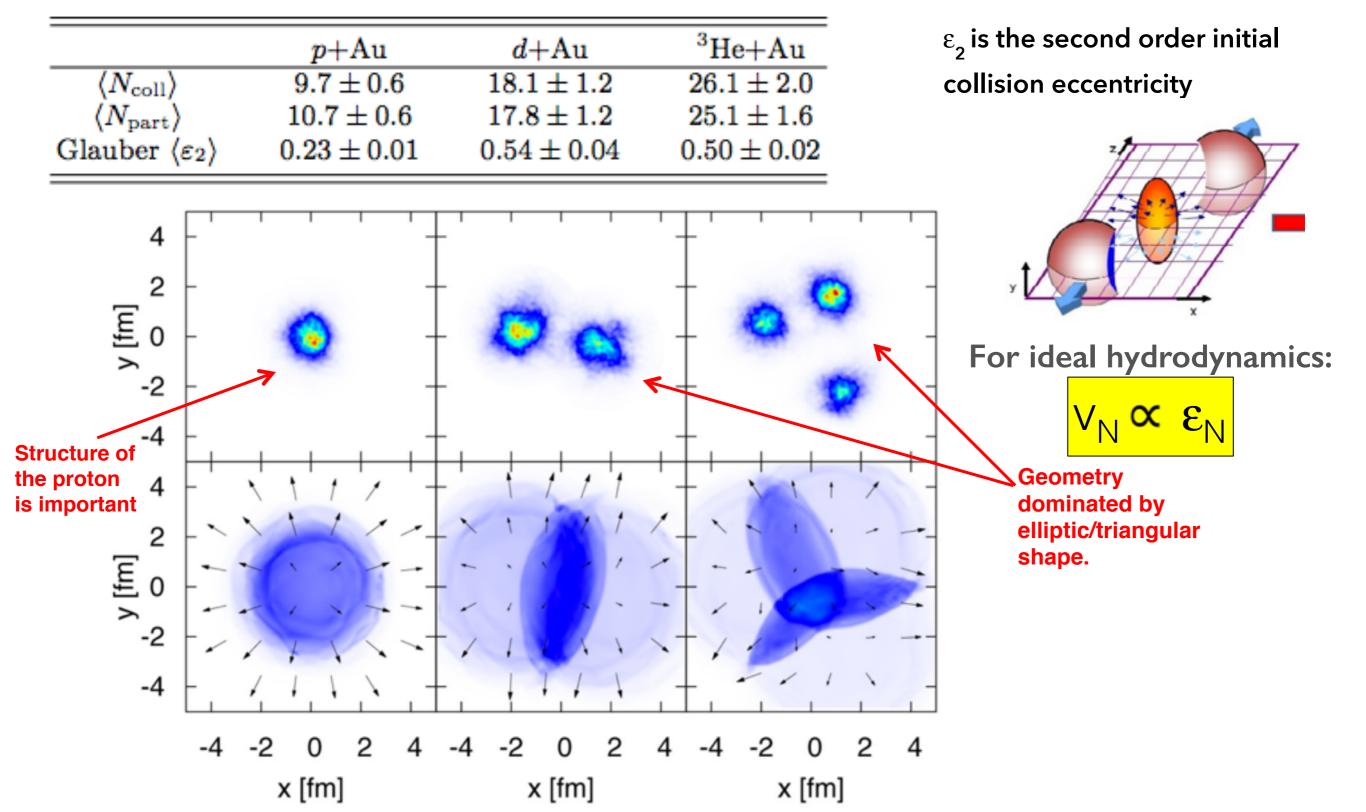
For ideal hydrodynamics:



0-5% central p+Au

d+Au

³He+Au 200 GeV



p+Au

 9.7 ± 0.6

0-5% central p+Au

 $\langle N_{
m coll}
angle$

d+Au

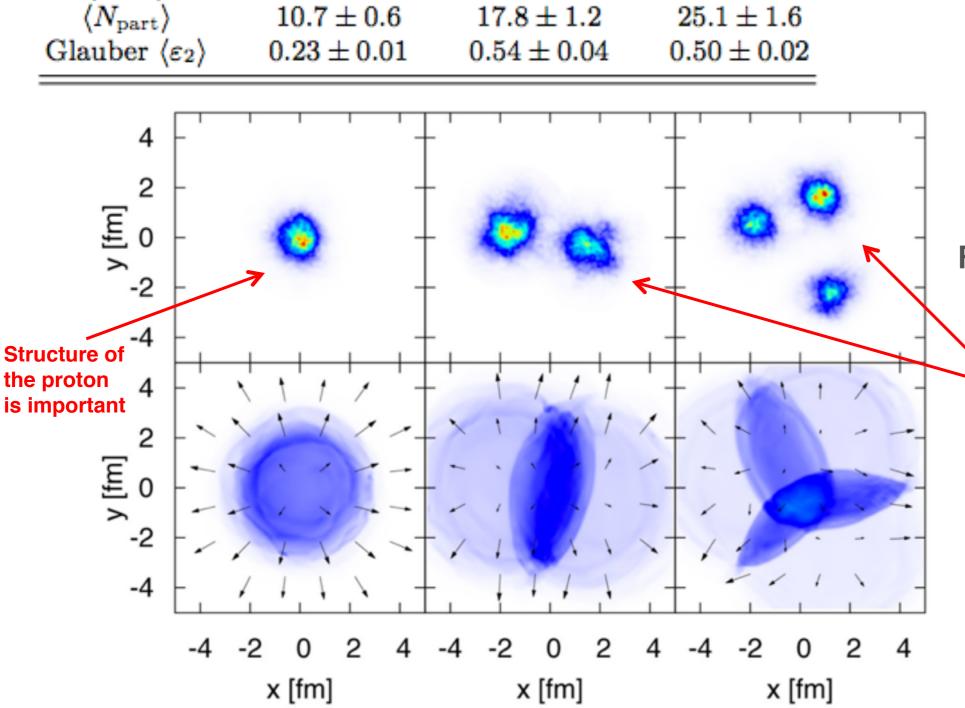
d+Au

 18.1 ± 1.2

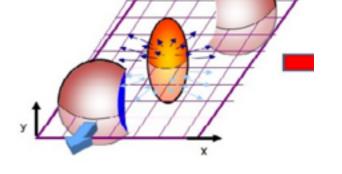
³He+Au 200 GeV

³He+Au

 26.1 ± 2.0



 ε_2 is the second order initial collision eccentricity



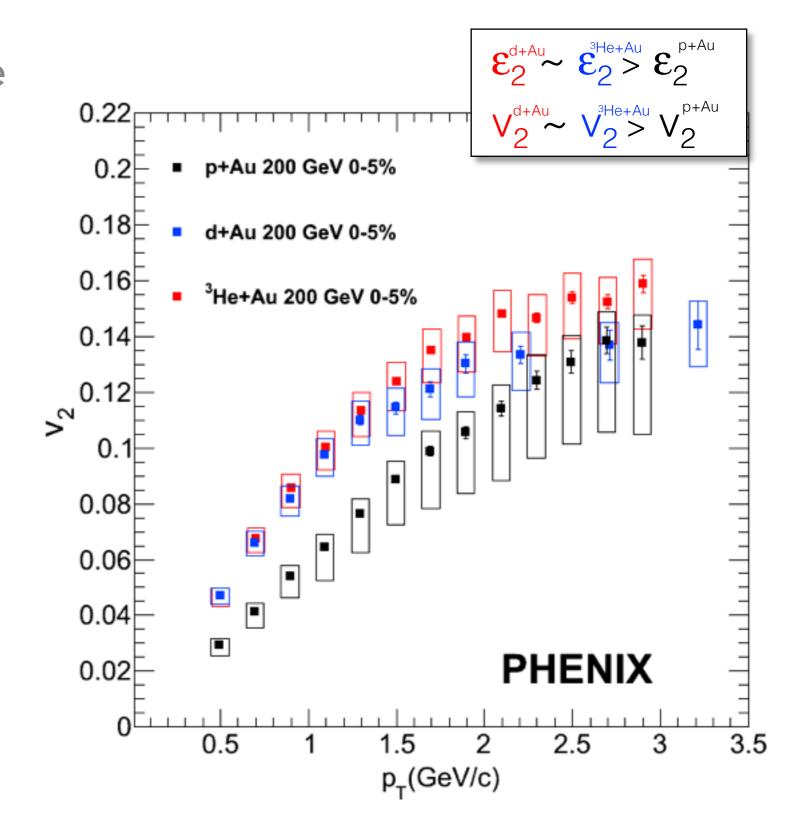
For ideal hydrodynamics:

 $V_N \propto \varepsilon_N$

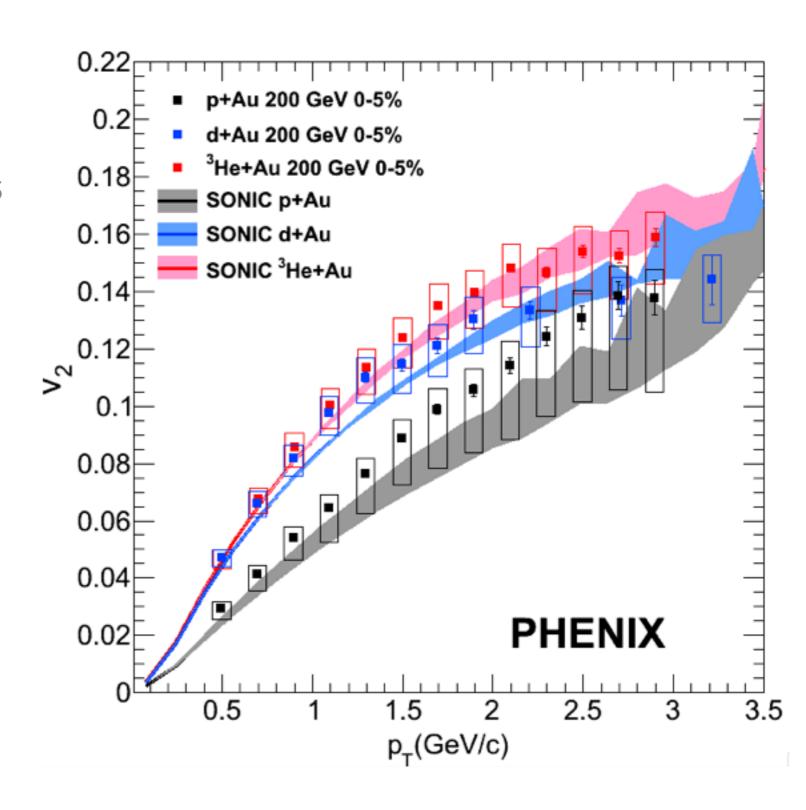
Geometry dominated by elliptic/triangular shape.

New 2016 d+Au Beam Energy Scan (200, 62, 39, 20 GeV) dataset (not in this talk)

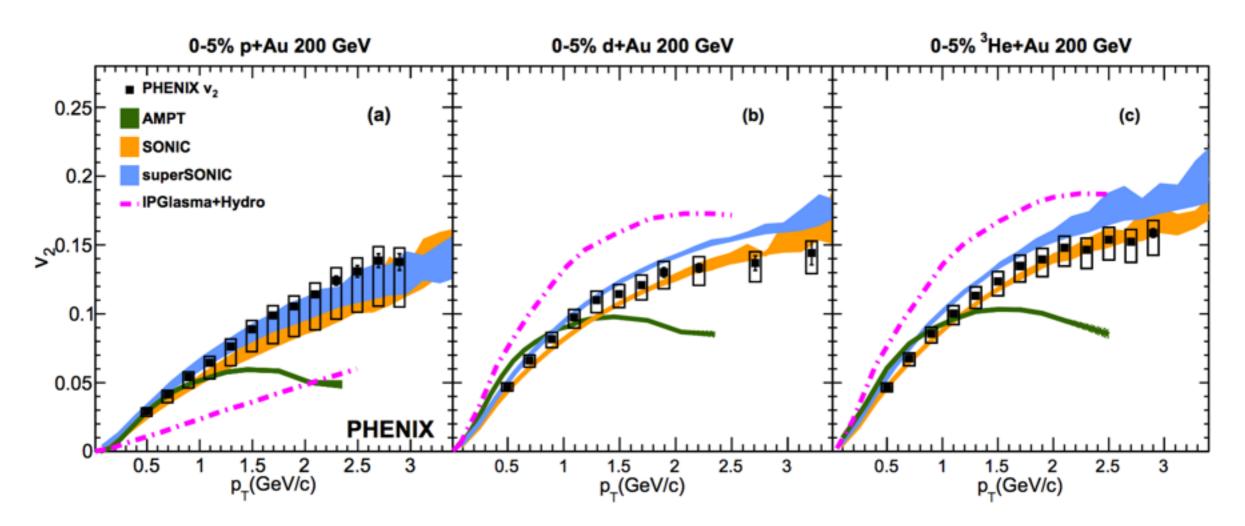
- v2 measured across three distinct small systems roughly follows Glauber ε2 scaling.
- The large non-flow systematic error on the p +Au points gives it the largest errors of all 3 systems.



- SONIC is a model which includes:
 - -MC Glauber
 - -Viscous Hydrodynamics
 - eta/s = 0.08
 - Hadronic cascade at T
 - = 170 MeV
 - Centrality matching
- The data is consistent with a viscous hydrodynamic model
- The epsilon scaling is not perfect.



SONIC HAS THE LARGEST AGREEMENT

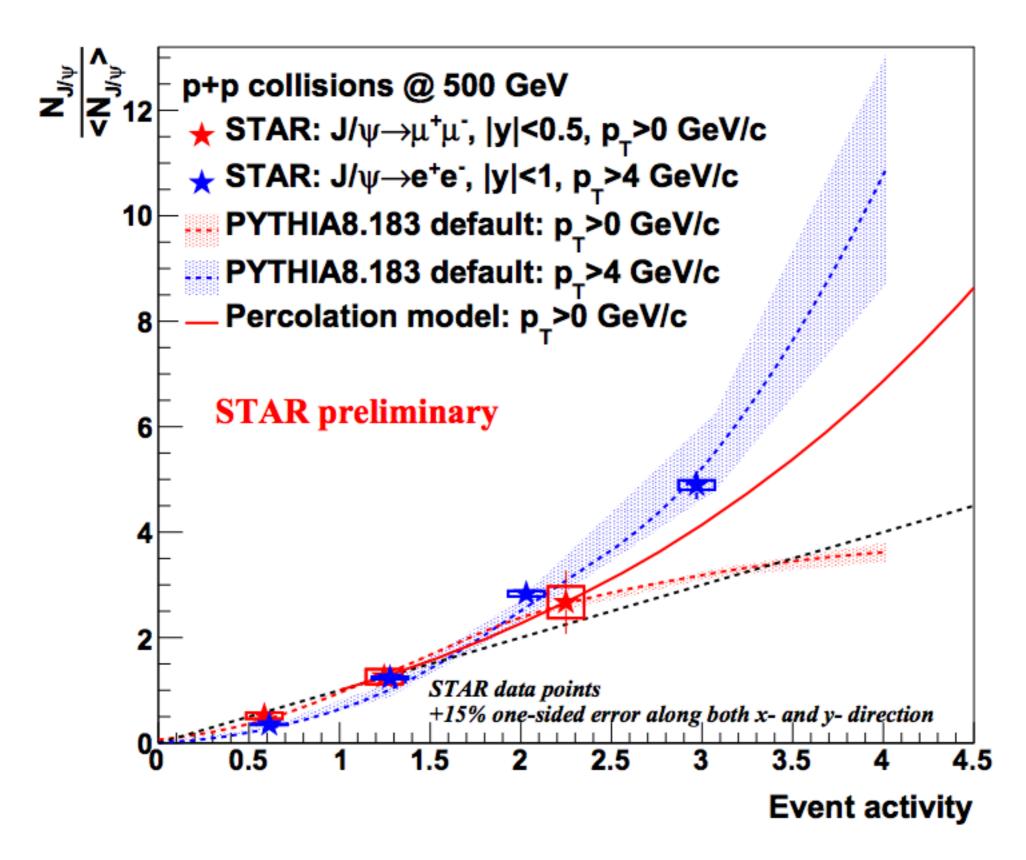


- ▶ IP Glasma (initial conditions) + Hydro can not simultaneously agree with all three systems.
- AMPT (a multi transport model) uses string-melty and a tunable parton scattering cross section.
 - Is in agreement with all three systems up to ~ 1.5 GeV
 - Does not use viscous hydrodynamics

- Centrality bias factors due to hard scattering must be calculated in small systems.
 - MPI probably plays a role in the centrality bias factors difference at the LHC in comparison to RHIC

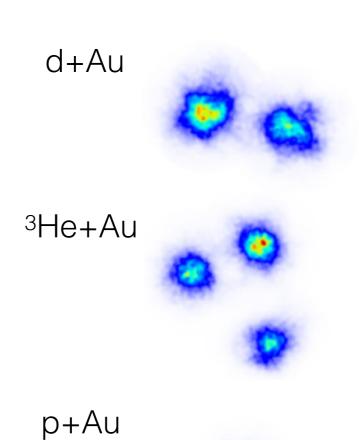
- Substantial flow coefficients are observed in p+Au, d+Au, and He3+Au at RHIC.
 - These observations are consistent with hydrodynamic models.
 - Could be evidence of QGP

THANK YOU



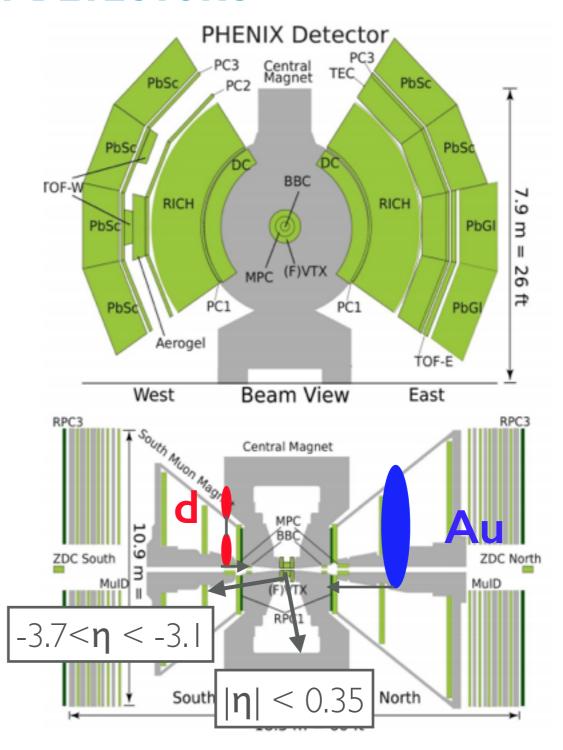
DATASET INFORMATION

- 2008 d+Au 200 GeV
 - delivered luminosity: 437 nb⁻¹
- 2014 He3+Au 200 GeV
 - delivered luminosity: 134 nb⁻¹
- 2015 p+Au 200 GeV
 - delivered luminosity: 1270 nb⁻¹
- New 2016 d+Au Beam Energy Scan (200, 62,
- 39, 20 GeV) dataset (not in this talk)
 - delivered luminosity: (289, 44, 7.2, 19.5) nb⁻¹



RELEVANT DETECTORS

- Beam Beam Counter (BBC): 3.1< |η|< 3.7
 - -Does minimum bias event triggering and centrality characterization
- Forward Silicon Vertex Detector (FVTX): $1 < |\eta| < 3$
 - Does precision charged particle measurements near the collision vertex
- Drift Chamber (DC): |η| < 0.35
 Does precision charged particle tracking and momentum measurement in mid-rapiditiy

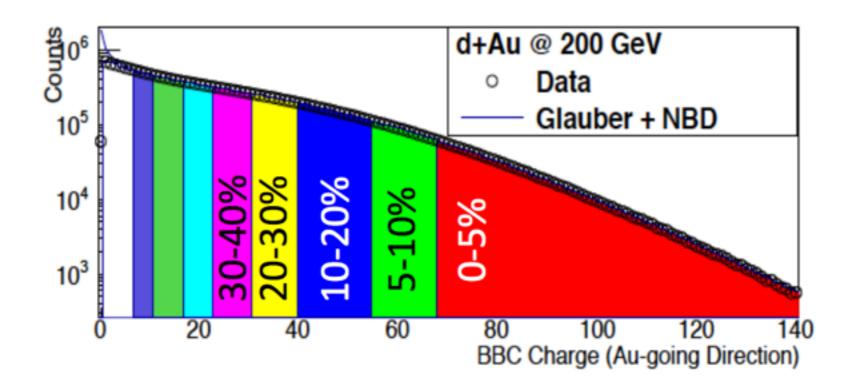


CENTRALITY RELATED TO NUMBER OF BINARY COLLISIONS

Using Negative Binomial Distribution (to model charge fluctuations):

$$NBD(x; \mu, \kappa) = \left(1 + \frac{\mu}{\kappa}\right) \frac{(\kappa + x - 1)!}{x!(\kappa - 1)!} \left(\frac{\mu}{\mu + \kappa}\right)^x$$

And fold in the MC Glauber: $P(x) = \sum_{n=1}^{N_{binary}(max)} Gl(n) \times NBD(x; n\mu, n\kappa)$



CENTRALITY BIAS FROM HARD SCATTERING

- In p+p 200 GeV, the BBC Minbias trigger fires 23% more often if there is a charged particle in mid-rapidity.
- Consider the inelastic cross section:

$$\sigma_{inel} = \sigma_{non-diff} + \sigma_{s-diff} + \sigma_{d-diff}$$
 get rid of

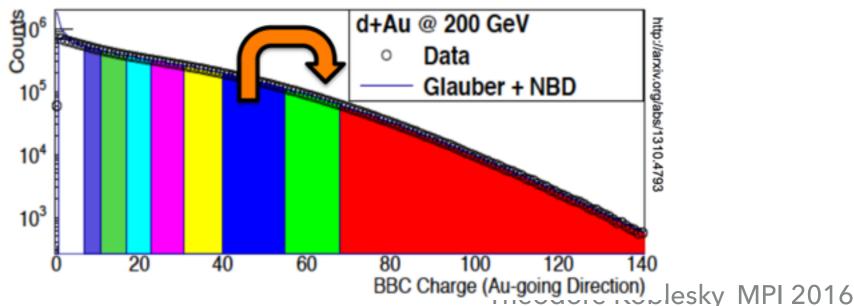
non-diffractive

single-diffractive

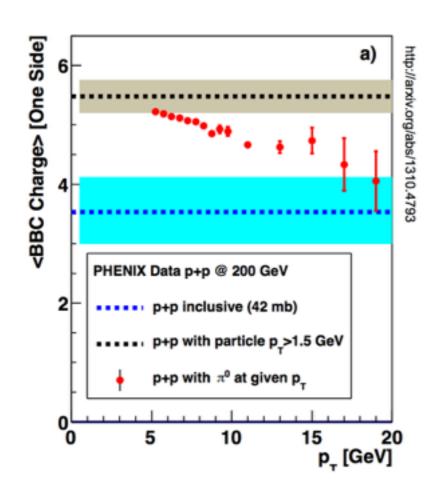
double-diffractive

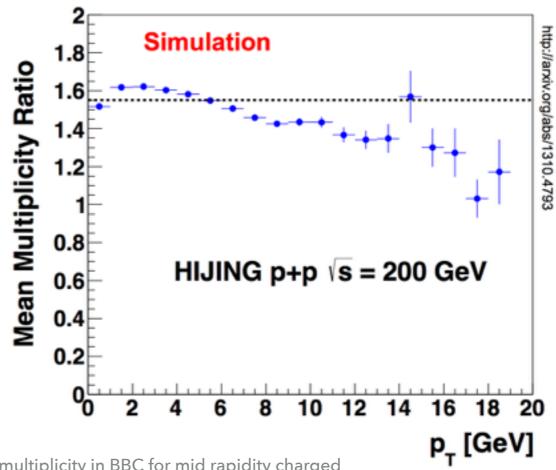
Hard scatterings tend to bias the charge into higher

centrality classes.



- Use HIJING (Heavy Ion Jet INteraction Generator),
 - A is a successful MC for heavy ion and p+p collisions.
- HIJING is able to reproduce a p+p bias and general pT dependence.



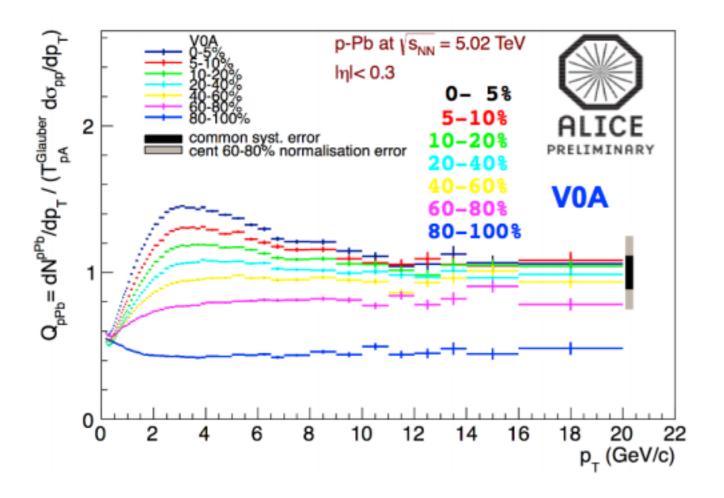


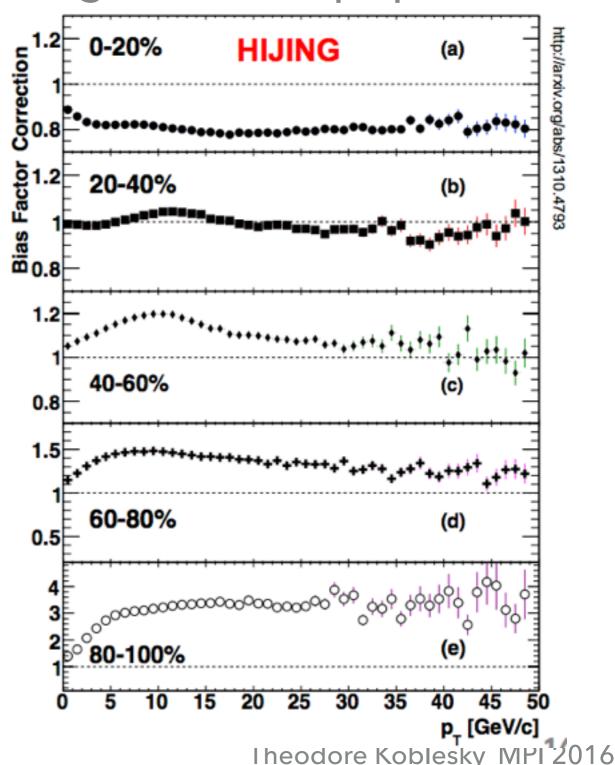
(ratio of mean multiplicity in BBC for mid rapidity charged particle triggered events to all inelastic collision events)
Theodore Koblesky MPI 2016

Because the HIJING simulation is successful, we can apply it to LHC p+Pb 5.02 TeV events using simulated p+p 5.02

TeV events

The centrality bias effect is larger at LHC energies, MPI could play a role.





Type

Systematic Uncertainty

2.0%

5.0%

3%

NON-FLOW IS THE DOMINANT SYSTEMATIC SOURCE

Source

Track Background

Event Pile-up

Non-Flow

Beam Angle

Event-Plane Detectors

- Systematic types:
 - Type A: point-to-point uncorrelated between pT bins
 - Type B: point-to-point correlated
 - Type C: overall normalization uncertainty
- Non-flow is estimated using p+p 200 GeV c2 measurements scaled down by the mean BBC south (Au going direction) charge ratio.
- There are other methods of estimating non-flow contributions, but it's not clear which is preferable.

