FINAL STATE MULTIPLICITY AND PARTICLE CORRELATION IN SMALL SYSTEMS

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OUTLOOK

Final state variables and particle correlation results will be shown and discussed under a Multiple Parton Interaction (MPI) interpretation.

- Final state multiplicity
 - Pseudorapidity and Transverse-momentum distributions of charged particles
 - Hadronic Event Shape
 - Forward Energy Measurement
- Particle correlation
 - Long-Range Near-Side Two particle angular correlation results at 13 TeV
 - Collectivity of strange hadrons
 - MPI as a way to understand LRNS

Measurements of particle yields and kinematic distributions are essential in exploiting the energy regimes of particle collisions at the LHC.

Eur. Phys. J. C 74 (2014) 3053

Charged particle pseudorapidity distribution:

 $\frac{1}{N_{events}}\frac{dN_{ch}}{d\eta} = \frac{C_{T2}\Sigma_M\Sigma_{pT}N_{tracks}(M,pT,\eta)\omega_{tracks}(M,pT,\eta)\omega_{event}(M,n_{T2})}{\Delta\eta\Sigma_MN_{evt}(M)\omega_{event}(M,n_{T2})}$

where ω_{tracks} and ω_{events} are correction factors and C_{T2} accounts for the track reconstruction efficiency

Charged particle pT distribution:

 $\frac{1}{N_{events}} \frac{dN_{ch}}{dpT_{leading}} = \frac{\Sigma_{\eta}N_{tracks}(\eta, pT_{leading}) \cdot C(pT_{leading}) \cdot C_{T2}(pT_{leading})}{N_{events} \cdot \Delta pT_{leading}}$

where C is the correction to stable particle level

Eur. Phys. J. C 74 (2014) 3053

8 TeV





- Studies on pseudorapidity and transverse momentum distributions led to the formulation of MPI theories in order to explain the disagreement data-MC
- From the 8 TeV analysis: interesting study on a wide pseudorapidity spectrum triggered by TOTEM
- Tunes based on Underlying Event variables do the best job in describing data (Gunnellini's talk)
- Comparison data-MC shows that models tuned on MPI observables better describe data.

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Energy dependence of pseudorapidity and pT



HADRONIC EVENT SHAPE



Tranverse thrust: $\tau_{\perp} = 1 - max_{\hat{\eta}_T} \frac{\sum_i |\vec{p}_{T,i} \cdot \hat{\eta}_T|}{\sum_i \vec{p}_{T,i}}$. $\tau_{\perp} = 0$ for perfectly balanced two-jet events and $\tau_{\perp} = (1-2/\pi)$ in isotropic multijet events. **Sphericity**: $S = \frac{3}{2}(\lambda_2 + \lambda_3)$ and **Transverse Sphericity**: $S_{\perp} = \frac{2\lambda_2}{\lambda_1 + \lambda_2}$ where λ_1, λ_2 and λ_3 are the normalized eigenvalues ($\lambda_1 < \lambda_2 < \lambda_3$) of the momentum tensor.

Events with a large number of MPI are expected to appear with a spherical shape, especially for high multiplicity.



- Transverse trust describe an higher isotropic contribution than expected in jet events
- Sphericity is higher in high-pT (and high multiplicity) events than expected
- Data/MC disagreement at large ΣpT

FORWARD ENERGY SPECTRUM

Total Energy [GeV]



28/11/2016

MULTIPLICITY FOR MPI STUDIES

- Final state multiplicity
 - Pseudorapidity and Transverse-momentum distributions of charged particles
 - Hadronic Event Shape
 - Forward Energy Measurement

So far we saw how Multiple Parton Interaction can help in the description of the final state multiplicity variables and hence the understanding of their dynamics

- Particle correlation
 - Long-Range Near-Side Two particle angular correlations
 - Strangeness particles production study to access LRNS

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Multiplicity plays a key role also in particle correlation, interplay with MPI can help in the results interpretation

PARTICLE CORRELATIONS

- Two-particle angular correlations for charged particles are studied in:
 - Short range: $|\Delta \eta| < 2$
 - Long range: $2 < |\Delta \eta| < 4.8$

Given:

- Signal function: $S_N(\Delta \eta, \Delta \phi) = \frac{1}{N(N-1)} \frac{d^2 N^{sign}}{d\Delta \eta \Delta \phi}$ charged two-particle pair density in the same events
- $B_N(\Delta\eta,\Delta\phi) = \frac{1}{N^2} \frac{d^2 N^{mixed}}{d\Delta\eta\Delta\phi}$ Background function: distribution of uncorrelated particle pairs from two randomly selected events
- Correlation function is defined as:

$$R (\Delta \eta, \Delta \phi) = \left((\langle N \rangle - 1) \left(\frac{S_N(\Delta \eta, \Delta \phi)}{B_N(\Delta \eta, \Delta \phi)} - 1 \right) \right)_{bins}$$









Phys. Rev. Lett. 116 (2016) 172302

p-p collisions results at 13 TeV:



For the **low-multiplicity sample** ($N_{trk}^{offline} < 35$), the dominant features is the peak near ($\Delta\eta, \Delta\phi$) = (0, 0) for pairs of particles originating from the same jet. The elongated structure at $\Delta\phi \approx \pi$ corresponds to pairs of particles from back-to-back jets.



Phys. Rev. Lett. 116 (2016) 172302

p-p collisions results at 13 TeV:



In **high-multiplicity pp events** ($N_{trk}^{offline} \ge 105$), in addition to these jet-like correlation structures, a "ridge"-like structure is clearly visible at $\Delta \phi \approx 0$, extending over a range of at least 4 units in $|\Delta \eta|$.

Confirmed what was observed at 7 TeV

At lower energy observed in p-A and A-A collisions

No such long-range correlations are predicted by PYTHIA.



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LRNS evolution with system size:



The long-range near-side yields have been measured for p-p, p-Pb and Pb-Pb collisions in CMS.

The ridge-like correlations become significant at a multiplicity value of about 40 in all three systems and exhibit a nearly linear increase for higher value.

For a given multiplicity value the associated yield in pp collision is roughly 10 % and 25 % of those observed in PbPb and pPb collissions respectively.

There a strong collision system size dependence of the longrange near-side correlations

CMS HIN-16-010 Strange hadron production and correlations in small colliding systems provide additional insights into the physical origin of the LRNS correlation



- The observed long-range ($|\Delta \eta| > 2$) correlations are quantified in terms of azimuthal anisotropy Fourier harmonics (v_n)
- The elliptic v₂ and triangular v₃ flow Fourier harmonics are extracted from long-range two-particle correlations at different values of center of mass energy and for different system size



- V₂:
 - No energy dependence
 - Qualitatively similar shape for pp, p-Pb, and Pb-Pb

- V₃:
- No energy dependence
- Values for pp are slightly different from p-Pb and Pb-Pb at **higher multiplicity** (N > 60)

The v_2 term is studied as a function of pT and particle species: **at high multiplicity** a deviation of v_2 term among various particle species is observed.

At low pT:

- K_s^0 is higher than $\Lambda/\overline{\Lambda}$
- the lighter particle species exhibit a stronger azimuthal anisotropy signal
- similar trend observed in A-A and p-Pb collisions

At high pT:

- $\Lambda/\overline{\Lambda}$ higher than K_s^0
- Reverse ordering is similar to previous observation in p-Pb and Pb-Pb collisions

Qualitatively consistent with the hydrodynamic models.



WHICH ROLE PLAYED BY MPI IN LONG-RANGE NEAR-SIDE CORRELATIONS?

- 1. For large impact parameter b the MPI tend to lie in the collision plane of the hardest interaction and the final state particles will have similar azimuthal angle ϕ (near-side)
- 2. MPI would require enough interactions to explain the high multiplicity events
- 3. Incoming partons have very different x_{bj} hence will have interactions in a broad pseudorapidity range η (long range)

Adding a modification in PYTHIA6, introducing a correlation between the azimuth of the event plane of individual MPI and the event plane of the hardest interaction



With this modification **PYTHIA** shows the ridge structure for the high-multiplicity moderate pT events.

BUT high multiplicity events are generally central collisions with an impact parameters $b \approx 0$.

CONCLUSION

- We can study MPI in two different dynamic regimes, multiplicity studies focus on the soft dynamics and constitute a complementary input on the Underlying Event analysis
- MPI are unavoidable:
 - Experimental evidences that MPI mechanisms are needed for a complete description of LHC final states
 - > To explain the high multiplicity events in the correlation effects
- High multiplicity in the final state plays a key role:
 - Still not completely understood (large deviation MC/Data in high multiplicity)
 - > MPI dynamics characterization and the system size dependence
 - Final state correlation, i.e. the «ridge effect» (Hydro? MPI alone? CGC?AMPT?)

THANK YOU FOR THE ATTENTION!



BACKUP

Charged particle pseudorapidity distribution:

 $\frac{1}{N_{events}} \frac{dN_{ch}}{d\eta} = \frac{C_{T2} \Sigma_M \Sigma_{pT} N_{tracks}(M, pT, \eta) \omega_{tracks}(M, pT, \eta) \omega_{event}(M, n_{T2})}{\Delta \eta \Sigma_M N_{evt}(M) \omega_{event}(M, n_{T2})}$ where ω_{tracks} and ω_{events} are correction factors and C_{T2} accounts for the track
reconstruction efficiency. M is the track multiplicity $\omega_{event}(M, n_{T2}) = \frac{1}{\epsilon_{trig}(n_{T2}) \epsilon_{PV}(M)} \qquad \omega_{track}(M, p_T, \eta) = \frac{1 - f_{np}(M, p_T, \eta)}{\epsilon_{track}(M, p_T, \eta) (1 + f_m(M, p_T, \eta))}$ $f_{np} = \frac{N_{reco}^{not matched tracks}(M, p_T, \eta)}{N_{rec}^{all track candidates}(M, p_T, \eta)}$

Charged particle pT distribution:

$$\frac{1}{N_{events}} \frac{dN_{ch}}{dpT_{leading}} = \frac{\sum_{\eta} N_{tracks}(\eta, pT_{leading}) \cdot C(pT_{leading}) \cdot C_{T2}(pT_{leading})}{N_{events} \cdot \Delta pT_{leading}}$$

where C is the correction to stable particle level
$$C(p_{T, \text{leading}}) = \frac{\left(\frac{1}{N} \frac{dN_{ch}}{dp_{T, \text{leading}}}\right)^{\text{gen}}}{\left(\frac{1}{N} \frac{dN_{ch}}{dp_{T, \text{leading}}}\right)^{\text{reco}}}$$

CENTRAL-FORWARD MULTIPLICITY ANALYSIS AT 8 TEV

Pseudorapidity and transverse momentum distribution were studied by CMS collaboration at 8 TeV (*Eur. Phys. J. C 74 (2014) 3053*) with a different trigger:

- Minimum Bias events are triggered by TOTEM T2 telescopes that cover the pseudorapidity region 5.3 < |η|
 < 6.6 for tracks with pT> 40 MeV.
- The measurements was performed for tracks with pT > 0.1 GeV and pT > 1 GeV in two consitions:
 - Inclusive sample with tracks reconstructed in the TOTEM T2 in either hemisphere
 - Sample enhanced in non-single diffractive dissociation events by requiring tracks in T2 both forward and backward hemispheres
- Selection criteria:
 - Rejection of the backgrounds requiring at least one reconstructed primary vertex with at least two tracks and with |z| < 15cm around the position of the nominal interaction
 - High purity tracks are selected with pT > 0.1 GeV or pT > 1 GeV and relative transverse momentum uncertainty less than 10 % within the pseudorapidity range $|\eta| < 2.4$
 - Track-vertex association applied requiring $d_{xy}/\sigma_{xy} < 3$ and $d_z/\sigma_z < 3$
 - For the measurement of the leading-track pT distribution the threshold for the tracks is 0.4 GeV V. MARIANI 28/11/2016 23

CENTRAL MULTIPLICITY ANALYSIS AT 13 TEV

Phys. Lett. B 751 (2015) 143

13 TeV results by CMS Collaborations:

- Measurements of $dN_{ch}/d\eta$ in the range $|\eta| < 2$ for inelastic proton-proton collision with 2015 data taken at 0 Tesla during a special low intensity beam configuration
- N_{ch} is defined to include decay products of particle with decay length cτ < 1 cm, products of secondary interactions are excluded
- Data are compared to PYTHIA8 v208 and EPOS LHC (Energy-conserving quantum mechanical multiple scattering approach, based on Parton, Off-shell remnants, and Splitting of parton ladders)

Event selection:

Selection of inelastic collision events:

- Online: a coincidence of signals form both the BPTX devices is required (both proton bunches crossing the IP)
- Offline: at least one reconstructed interaction vertex is required

HADRONIC EVENT SHAPE



Figure 1: Mean transverse sphericity as a function of charged particle multiplicity for pp collisions at $\sqrt{s} = 7$ TeV. The statistical errors are displayed as error bars and the systematic uncertainties as the shaded area. The results are shown for the different event classes: (a) "bulk," (b) "soft" and (c) "hard."

MB events are analyzed

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- The sphericity in data is steadily rising with multiplicity suggesting a more isotropic distribution of tracks in azimuth than the models.
- The general agreement between models is better for "soft" events while for the "hard" ones the disagreement is up to ~ 20% at low and high multiplicity

FORWARD ENERGY SPECTRUM





Event at $\sqrt{s} = 0.9$ Pythia6 D6T without multiple parton interaction completely fails the data description



Comparison between 13 TeV (red) and 7 TeV data,

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Comparison between CMS data at 7 TeV (CMS-QCD-10-002) and PYTHIA8 in 4 range of pT bins.

Two discrepancies:

- The strength of the away-side correlation is over –or underpredicted for almost all the bins
- PYTHIA8 fails to reproduce the local maximum near $\Delta \phi \approx 0$ in any of the pT or multiplicity bins.

The long range, near side correlation increases in strength with increasing multiplicity and is stronger in the bin 1 < pT < 2 GeV

Deeper study on v_2 term is done evaluating this variables from simultaneously correlating several (no less than four) particles.

- Suppress the short-range two particle correlations such as jets and resonance decays and as a
- Powerful tool to directly probe the collective nature of the observed azimuthal correlations.



- v₂{2}≈v₂{4}≈v₂{6} in pp collisions (left)
- Qualitatively similar results seen in high multiplicity pp and pPb, as well as peripheral PbPb for v_2 {4} and v_2 {6}
- The ratio of v2{4} to v2{2} is related to the total number of fluctuating sources in the initial state of a collision. The comparable magnitudes of v2{2} and v2{4} signals observed in pp collisions may indicate a smaller number of initial fluctuating sources that drive the long-range correlations seen in the final state.

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Strong evidence for the collective nature of the long-range correlations observed in pp collisions.