

Reunión Anual de la División de Partículas y Campos Sociedad Mexicana de Física

*Efectos de Nueva física en el autoacoplamiento hhh
inducido por fermiones espejos del MHP+T*

Dr. Agustín Moyotl Acuahuitl, Dr. Miguel Ángel Pérez Angón
Departamento de Física, CINVESTAV

BUAP, 23 de mayo de 2016



TABLE OF CONTENTS

- 1.- *Introduction.*
- 2.- *Higgs physics*
- 3.- *hhh at one loop level in SM*
- 4.- *Littlest Higgs Model with T-parity*
- 5.- *hhh at one loop in LHM with T-parity*
- 7.- *Conclusions*

INTRODUCTION

Observation of a new boson at a mass of 125 GeV with the CMS experiment at the LHC ☆

CMS Collaboration ☆

CERN. This content
Observation of a new particle in the search for the Standard Model Higgs boson with the ATLAS detector at the LHC ☆

ATLAS Collaboration

This paper is dedicated to the contributions to the

$\sqrt{s} = 7$ TeV in 2011 and 5.8 fb^{-1} at $\sqrt{s} = 8$ TeV in 2012

Article
Received
Received
Accepted
Available
Editor:
Keywords
CMS
Physics
Higgs

ARTICLE INFO

Article history:

Received 31 July 2012

Received in revised form 8 August 2012

Accepted 11 August 2012

Available online 14 August 2012

Editor: W.-D. Schlatter

ABSTRACT

A search for the Standard Model Higgs boson in proton–proton collisions with the ATLAS detector at the LHC is presented. The datasets used correspond to integrated luminosities of approximately 4.8 fb^{-1} collected at $\sqrt{s} = 7$ TeV in 2011 and 5.8 fb^{-1} at $\sqrt{s} = 8$ TeV in 2012. Individual searches in the channels $H \rightarrow ZZ^{(*)} \rightarrow 4\ell$, $H \rightarrow \gamma\gamma$ and $H \rightarrow WW^{(*)} \rightarrow e\nu\mu\nu$ in the 8 TeV data are combined with previously published results of searches for $H \rightarrow ZZ^{(*)}$, $WW^{(*)}$, $b\bar{b}$ and $\tau^+\tau^-$ in the 7 TeV data and results from improved analyses of the $H \rightarrow ZZ^{(*)} \rightarrow 4\ell$ and $H \rightarrow \gamma\gamma$ channels in the 7 TeV data. Clear evidence for the production of a neutral boson with a measured mass of 126.0 ± 0.4 (stat) ± 0.4 (sys) GeV is presented. This observation, which has a significance of 5.9 standard deviations, corresponding to a background fluctuation probability of 1.7×10^{-9} , is compatible with the production and decay of the Standard Model Higgs boson.

© 2012 CERN. Published by Elsevier B.V. All rights reserved.

mass of 126.0 ± 0.4 (stat) ± 0.4 (sys) GeV

$H \rightarrow ZZ^{(*)} \rightarrow 4\ell$, $H \rightarrow \gamma\gamma$ and $H \rightarrow WW^{(*)} \rightarrow e\nu\mu\nu$

Higgs Decay modes at LHC

Direct measurement

$$h^0 \rightarrow ZZ^* \rightarrow 4\ell$$

ATLAS, *Phys.Lett. B* 716, 1 (2012) $m_h = 126.0 \pm 0.4 \pm 0.4 \text{ GeV}$

$$h^0 \rightarrow \gamma\gamma^* \rightarrow 4\ell$$

ATLAS, *Phys.Rev.D* 90, 052004 (2014) $m_h = 125.36 \pm 0.41 \text{ GeV}$

CMS, *Phys.Rev.Lett.* 114 191803 (2015) $m_h = 125.09 \pm 0.21 \pm 0.11 \text{ GeV}$

$$h^0 \rightarrow WW^* \rightarrow \nu\ell\nu\ell$$

CMS, *Phys.Lett.B* 716, 30 (2012) $m_h = 125.3 \pm 0.4 \pm 0.5 \text{ GeV}$

$$\ell = e, \mu$$

$$h^0 \rightarrow \tau\tau, \tau \rightarrow \ell\nu\bar{\nu}$$

CMS, *Nature Phys.* 10 557 (2014) $m_h = 125 \text{ GeV}$

ATLAS, *JHEP* 1504, 117 (2015) $m_h = 125.36 \text{ GeV}$

$$h^0 \rightarrow b\bar{b}$$

ATLAS, *JHEP* 1501, 069 (2015) $m_h = 125.36 \text{ GeV}$

CMS, *Phys.Rev.D* 92 032008 (2015) $m_h = 125 \text{ GeV}$

Upper bound

$$h^0 \rightarrow \mu\mu$$

ATLAS, *Phys.Lett.B* 738, 68 (2014) $m_h = 125 \text{ GeV}$

$$h^0 \rightarrow \mu\mu(ee)$$

CMS, *Phys.Rev.D* 92 032008 (2015) $m_h = 120 - 150 \text{ GeV}$

$$h^0 \rightarrow Z\gamma \quad Z \rightarrow \ell\ell$$

ATLAS, *Phys.Lett.B* 732, 8 (2014) $m_h = 120 - 150 \text{ GeV}$

CMS, *Phys.Lett.B* 753 341 (2016) $m_h = 125 \text{ GeV}$

HIGGS PHYSICS

Higgs Decay modes at LHC

Direct measurement

$$h^0 \rightarrow ZZ^* \rightarrow 4\ell$$

$$h^0 \rightarrow \gamma\gamma^* \rightarrow 4\ell$$

$$h^0 \rightarrow WW^* \rightarrow \nu\ell\nu\ell$$

$$\ell = e, \mu$$

$$h^0 \rightarrow \tau\tau, \begin{cases} \tau \rightarrow \ell\nu\bar{\nu} \\ \tau \rightarrow had. + \bar{\nu} \end{cases}$$

$$h^0 \rightarrow b\bar{b}$$

Upper bound

$$h^0 \rightarrow \mu\mu$$

$$h^0 \rightarrow \mu\mu(ee)$$

$$h^0 \rightarrow Z\gamma \quad Z \rightarrow \ell\ell$$

“Τησ δατα αρε φουνδ το βε χομπατιβλε ωιτη τησ Στανδαρδ Μοδελ εΞεπχατιονσ φορ α Ηιγγσ βοσον ατ α μασσ οφ 125.36 Γεσ... Τογεταιηερ τηεσ αχχουντ φορ αππροξιματελψ 88 % οφ αλλ δεχαψσ οφ α ΣΜ Ηιγγσ βοσον. *ATLAS*, Eur. Phys. J. C (2016) 76:6

Exploring electroweak symmetry breaking at the LHC.

$$\mathcal{L}_{ME+h} = \mathcal{L}_{Fermionic} + \mathcal{L}_{Gauge} - V(\Phi)$$

Higgs self-couplings

$$\begin{aligned} V(\Phi) &= -\lambda v^2 \Phi^\dagger \Phi + \lambda (\Phi^\dagger \Phi)^2 \\ &= \lambda v^2 h^2 + \lambda v h^3 + \frac{1}{4} \lambda h^4 \end{aligned}$$

Higgs Decay modes at LHC

Direct measurement

$$h^0 \rightarrow ZZ^* \rightarrow 4\ell$$

$$h^0 \rightarrow \gamma\gamma^* \rightarrow 4\ell$$

$$h^0 \rightarrow WW^* \rightarrow \nu\ell\nu\ell$$

$$\ell = e, \mu$$

$$h^0 \rightarrow \tau\tau, \begin{cases} \tau \rightarrow \ell\nu\bar{\nu} \\ \tau \rightarrow had. + \bar{\nu} \end{cases}$$

$$h^0 \rightarrow b\bar{b}$$

Upper bound

$$h^0 \rightarrow \mu\mu$$

$$h^0 \rightarrow \mu\mu(ee)$$

$$h^0 \rightarrow Z\gamma \quad Z \rightarrow \ell\ell$$

“Τησ δατα αρε φουνδ το βε χομπατιβλε ωιτη τησ Στανδαρδ Μοδελ εξεπεχτατιονσ φορ α Ηιγγσ βοσον ατ α μασσ οφ 125.36 Γεσ... Τογεταιηερ τηεσ αχχουντ φορ αππροξιματαελψ 88 % οφ αλλ δεχαψσ οφ α ΣΜ Ηιγγσ βοσον.”
ATLAS, Eur. Phys. J. C (2016) 76:6

Exploring electroweak symmetry breaking at the LHC.

“Σεαρχηεσ φορ βοτη ρεσοναντ ανδ νονρεσοναντ Ηιγγσ βοσον παιρ προδυχτιον... Νο επιδενχε οφ τηειρ προδυχτιον ισ οβσερπειδ”
ATLAS, Phys. Rev. D 92, 092004 (2015)

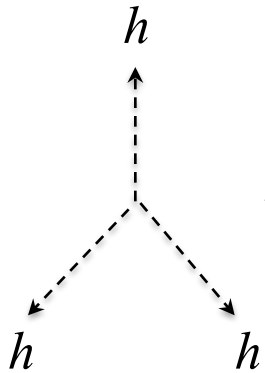
Higgs self-couplings

$$\begin{aligned} V(\Phi) &= -\lambda v^2 \Phi^\dagger \Phi + \lambda (\Phi^\dagger \Phi)^2 \\ &= \lambda v^2 h^2 + \lambda v h^3 + \frac{1}{4} \lambda h^4 \end{aligned}$$

HHH AT ONE LOOP LEVEL IN SM

Self-coupling hhh at tree and one loop level in SM

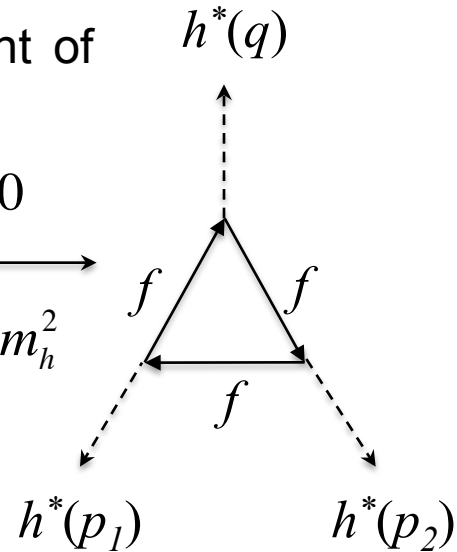
At tree level the self-coupling hhh is independent of the higgs bosons 4-momentum.



$$\sim \frac{3m_h^2}{v}$$

$$q + p_1 + p_2 = 0$$

$$q^2 \neq p_1^2 \neq p_2^2 \neq m_h^2$$



At one loop level and with off-shell Higgs, we need more diagrams to satisfy the Bose symmetry.

$$+ \begin{pmatrix} q \\ p_1 \leftrightarrow p_2 \end{pmatrix} + \begin{pmatrix} q \leftrightarrow p_1 \\ p_2 \end{pmatrix} + \begin{pmatrix} q \leftrightarrow p_2 \\ p_1 \end{pmatrix} + \begin{pmatrix} p_1 \leftrightarrow p_2 \\ q \leftrightarrow p_1 \end{pmatrix} + \begin{pmatrix} p_1 \leftrightarrow p_2 \\ q \leftrightarrow p_2 \end{pmatrix}$$

$$= \Gamma_{h^*h^*h^*}(q^2, p_1^2, p_2^2) = \frac{3m_h^2}{v} \lambda_{hhh}$$

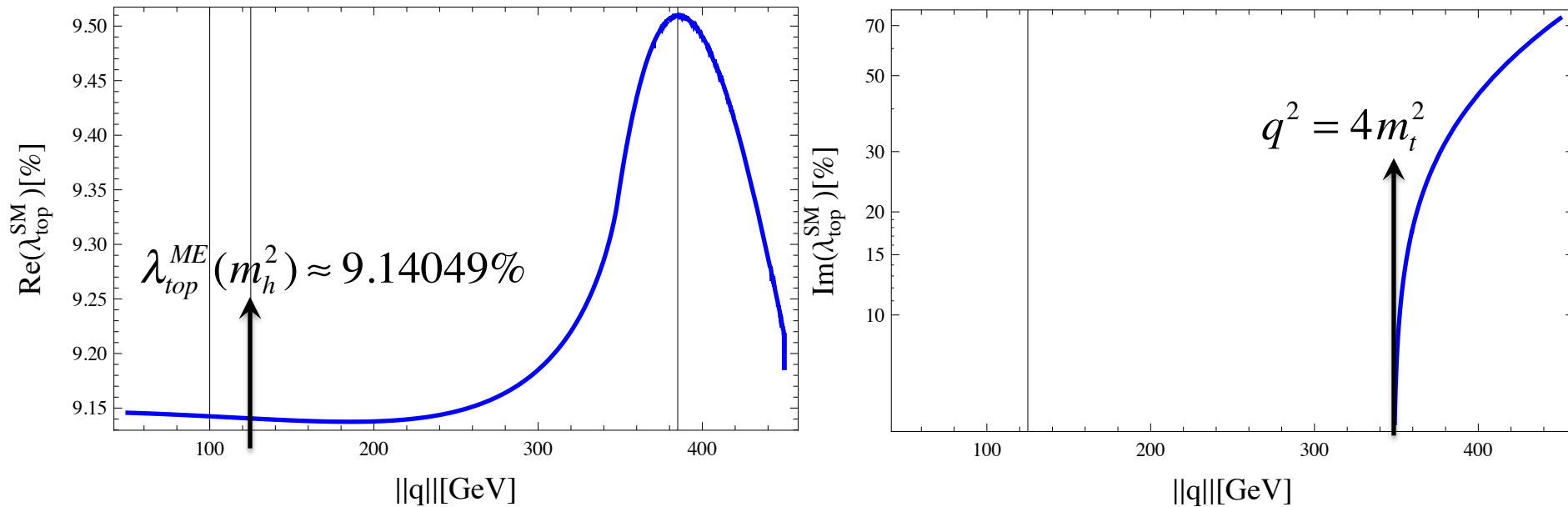
Here $\lambda_{\eta\eta\eta}$ is the correction to self-coupling hhh .

HHH AT ONE LOOP LEVEL IN SM

Self-coupling hhh at tree and one loop level in SM

The leading one loop contribution of the top quark to self-coupling is derived as:

$$\lambda_t^{ME}(q^2) = \frac{gm_t^2 s_t^2 N_c}{3!16\pi^2 m_W^2} \left[18 - \int_{x=0}^1 \int_{y=0}^{1-x} \Xi_t^{ME}(x, y, s_t, s_q) dx dy \right]$$



Previous results:

$$\lambda_{\text{top}}^{ME}(m_h^2) \approx 9.14693\%$$

Phys. Lett. B558, 157 (2003)

$$\lambda_{\text{top}}^{ME}(m_h^2) \approx 9.8221\%$$

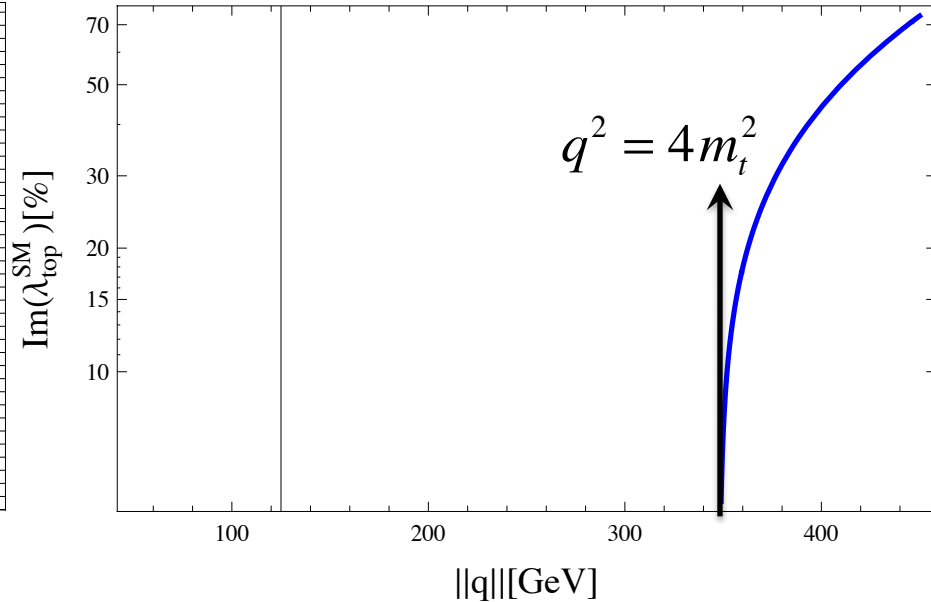
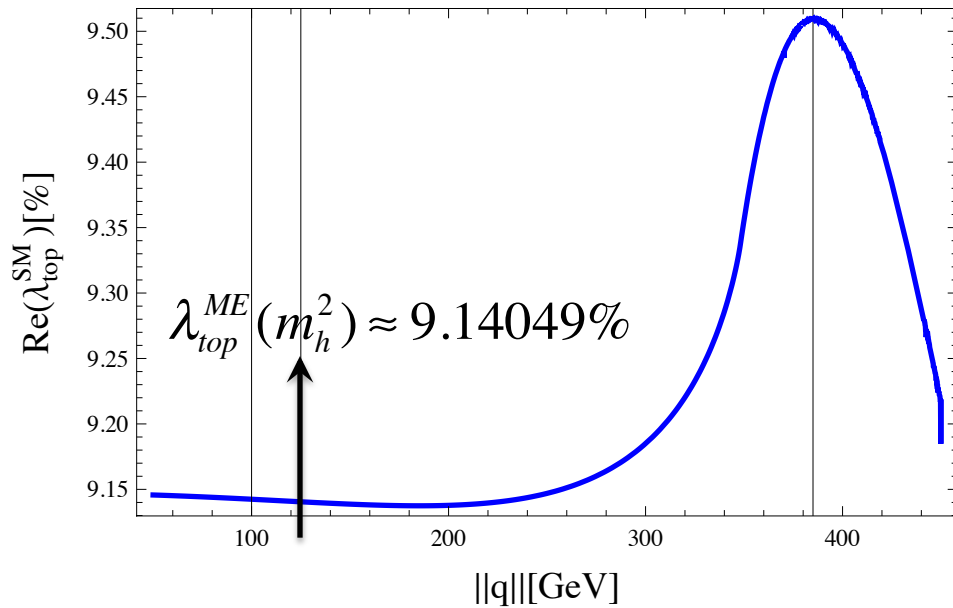
Eur. Phys. J. C23, 163 (2002)

HHH AT ONE LOOP LEVEL IN SM

Self-coupling hhh at tree and one loop level in SM

The leading one loop contribution of the top quark to self-coupling is derived as:

$$\lambda_t^{ME}(q^2) = \frac{gm_t^2 s_t^2 N_c}{3!16\pi^2 m_W^2} \left[18 - \int_{x=0}^1 \int_{y=0}^{1-x} \Xi_t^{ME}(x, y, s_t, s_q) dx dy \right]$$



Another contributions:

$$\lambda_b^{ME}(m_h^2) \approx 3.4 \times 10^{-6}\% + i6.4 \times 10^{-9}\%$$

$$\lambda_\tau^{ME}(m_h^2) \approx 2.5 \times 10^{-8}\% + i7.9 \times 10^{-12}\%$$

$$\lambda_{Z+W}^{ME}(m_h^2) \approx 0.0726\%$$

$$\lambda_h^{ME}(m_h^2) \approx 1.8397\%$$

LITTLEST HIGGS MODEL WITH T-PARITY

Main ideas

1.- Little higgs models are based on a collective symmetry breaking pattern.

$$\text{(Global) } SU(5) \xrightarrow{f \sim O(\text{TeV})} SO(5)$$

$$\text{(Gauged) } [SU(2)_1 \otimes U(1)_1] \otimes [SU(2)_2 \otimes U(1)_2] \longrightarrow SU(2) \otimes U(1) \xrightarrow{v_{SM}} U(1)_{EM}$$

2.- There are 14 Goldstone bosons, and are parametrized by a nonlinear sigma model.

3.- After the first symmetry breaking there are four heavy gauge boson: Z_H, A_H, W_H^\pm .

4.- After the second symmetry breaking we have the SM fields, and seven additional scalars fields.

5.- The LHM had an alternative to hierarchy problem.

6.- The heavy photon is a dark matter candidate. *Phys. Rev. D 88, 075018 (2013)*

7.- A natural way to define the action of T-parity on the gauge fields is:

$$W_{1\mu}^a \leftrightarrow W_{2\mu}^a \quad B_{1\mu}^a \leftrightarrow B_{2\mu}^a$$

8.- The constraints from Higgs couplings results from the 8 TeV run at the LHC ... exclude f up to 694 GeV. *JHEP02 (2014) 053*

LITTLEST HIGGS MODEL WITH T-PARITY

Mirror fermions

1.- Little higgs models are based on a collective symmetry breaking pattern.

$$\text{(Global) } SU(5) \xrightarrow{f \sim O(\text{TeV})} SO(5)$$

$$\text{(Gauged) } [SU(2)_1 \otimes U(1)_1] \otimes [SU(2)_2 \otimes U(1)_2] \longrightarrow SU(2) \otimes U(1) \xrightarrow{v_{SM}} U(1)_{EM}$$

For each SM $SU(2)_L$ fermion doublet, a fermion doublet under $SU(2)_1$ another under $SU(2)_2$ are introduced. The T-parity even linear combination is associated with the SM, while the T-odd combination is given a mass of order the scale f .

The mirror fermion acquire mass through the $SU(5)$ and T invariant Yukawa interaction.

$$\mathcal{L}_{mirror} = -\kappa_{ii} f (\bar{\Psi}_2^i \xi + \bar{\Psi}_1^i \Sigma_0 \Omega \xi^\dagger \Omega) \Psi_R^i$$

Then the masses and the Higgs coupling for the u mirror quark and the mirror neutrino are.

$$m_{\ell_H} = m_{d_H} = \sqrt{2} \kappa_{ii} f$$

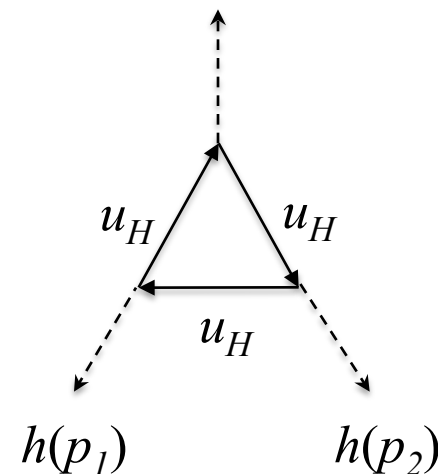
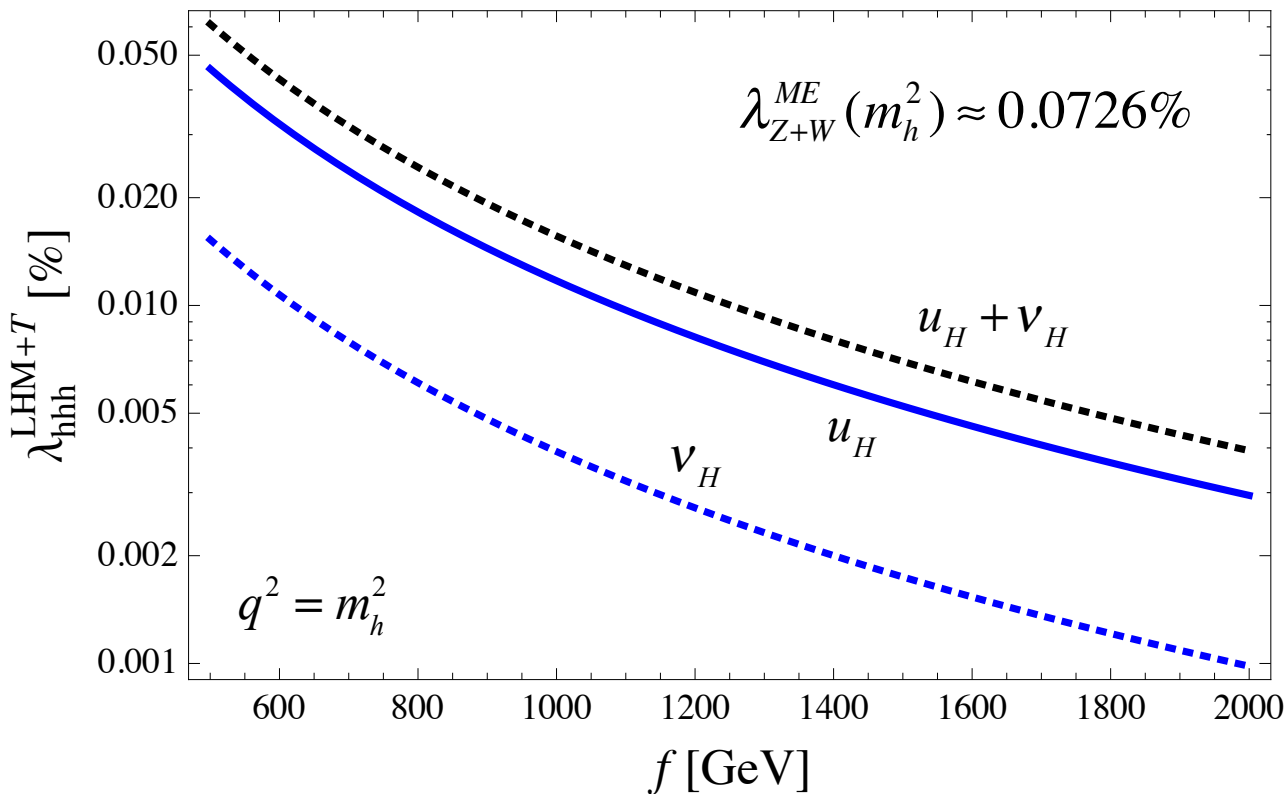
$$m_{\nu_H} = m_{u_H} = \sqrt{2} \kappa_{ii} f \left(1 - \frac{v^2}{8f^2} \right)$$

d_H and l_H do not have direct couplings with the Higgs boson.

$$h \bar{u}_H u_H \sim h \bar{\nu}_H \nu_H \sim \frac{i \kappa_{ii}}{2\sqrt{2}} \left(\frac{v}{f} \right)$$

HHH AT ONE LOOP IN LHM WITH T-PARITY

$$\lambda_{u_H, v_H}^{LHM+T}(q^2) = \frac{N_c \kappa_{ii}^2 m_{u_H, v_H} m_W}{3\pi^2 2^{\frac{11}{2}} 3! g m_h^2} \frac{v^3}{f^3} \left[18 - \int_{x=0}^1 \int_{y=0}^{1-x} \Xi_{u_H}^{ME}(x, y, s_{u_H, v_H}, s_q) dx dy \right] h^*(q)$$



Making the decay $q_H \rightarrow V_H q$ kinematically allowed, this corresponds to $\kappa_{ii} > 0.45$.

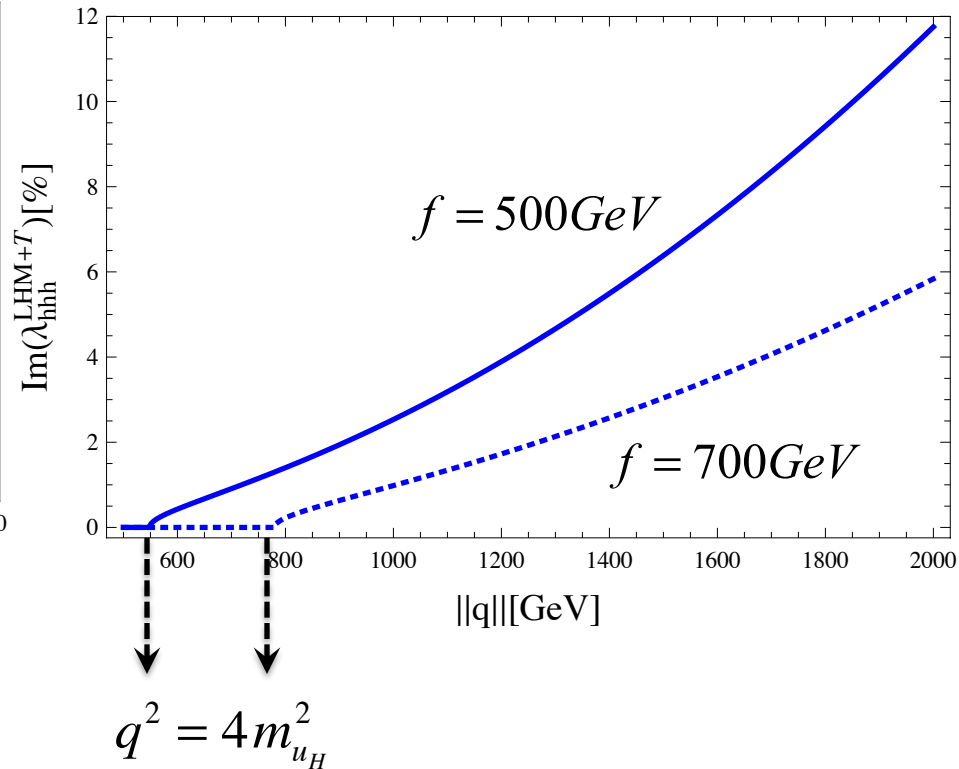
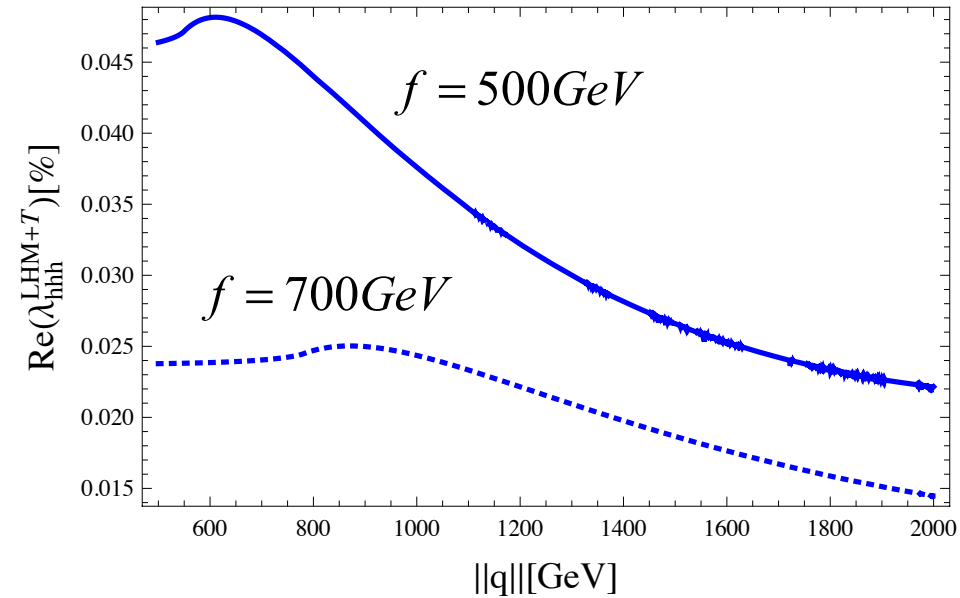
JHEP02 (2014) 053

$$m_{v_H} = m_{u_H} = \sqrt{2} \kappa_{ii} f \left(1 - \frac{v^2}{8f^2} \right)$$

$$h \bar{u}_H u_H \sim h \bar{v}_H v_H \sim \frac{i \kappa_{ii}}{2\sqrt{2}} \left(\frac{v}{f} \right)$$

HHH AT ONE LOOP IN LHM WITH T-PARITY

$$\lambda_{u_H, \nu_H}^{LHM+T}(q^2) = \frac{N_c \kappa_{ii}^2 m_{u_H, \nu_H} m_W}{3\pi^2 2^{\frac{11}{2}} 3! g m_h^2} \frac{v^3}{f^3} \left[18 - \int_{x=0}^1 \int_{y=0}^{1-x} \Xi_{u_H}^{ME}(x, y, s_{u_H, \nu_H}, s_q) dx dy \right]$$



$$m_{\nu_H} = m_{u_H} = \sqrt{2} \kappa_{ii} f \left(1 - \frac{v^2}{8f^2} \right)$$

$$h\bar{u}_H u_H \sim h\bar{\nu}_H \nu_H \sim \frac{i\kappa_{ii}}{2\sqrt{2}} \left(\frac{v}{f} \right)$$

CONCLUSIONS

1.- *To test the electroweak symmetry breaking sector of the SM, not only should the couplings of the Higgs boson to gauge boson and fermions be measured, but also the self-couplings of the Higgs boson.*

2.- *In the SM, the one-loop contribution to self-couplings hhh is $\approx 11\%$, where the top quark contribution is $\approx 9.14\%$*

3.- *In the LHM+T, the one-loop contribution of mirror fermions to self-coupling hhh is $\approx 2.3 \times 10^{-2}$ for $f=700$ GeV.*