



DEPARTMENT OF PHYSICS

UNIVERSITY OF CAPE TOWN

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*(Inverse) Magnetic Catalysis from the
Properties of the QCD Coupling in a Magnetic
Field.*

Luis A. Hernandez

International workshop
QCD challenges at the LHC: from pp to AA.
Taxco, January 21 2016.

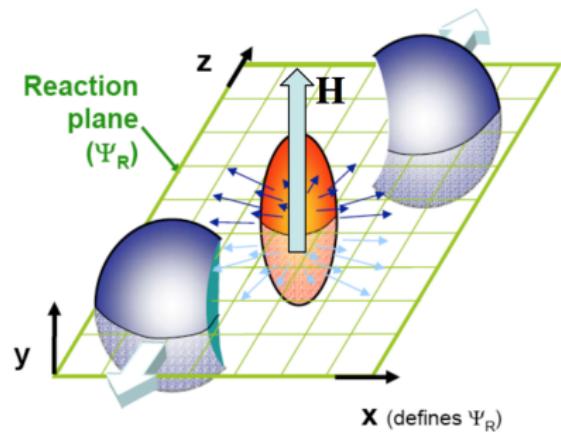
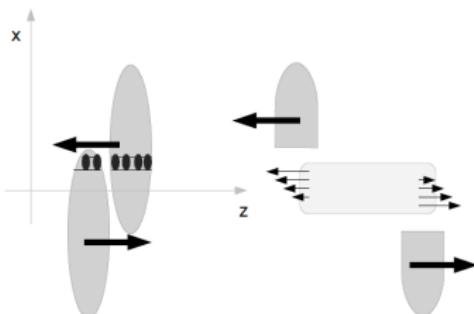
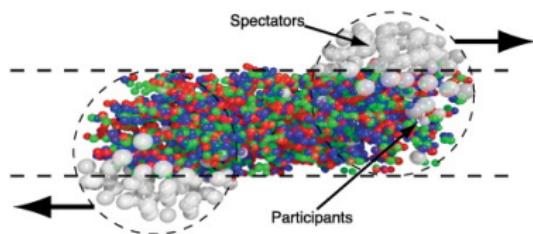
Outline

- Motivation.
- Magnetic fields \Leftrightarrow QCD phase transition.
- Magnetic fields \Leftrightarrow Vacuum and finite temperature systems.
- Magnetic fields \Leftrightarrow Coupling constants.
- Conclusions.

Nature in the presence of Magnetic fields

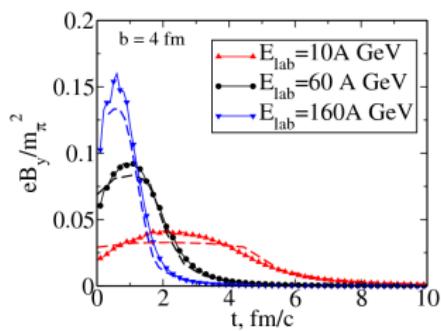
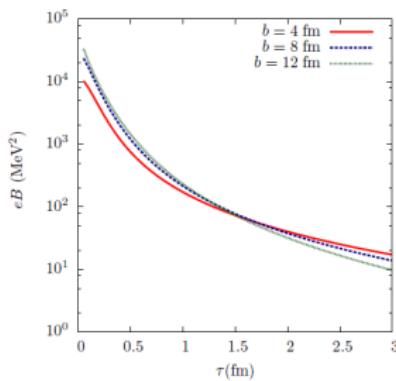
- The earth's magnetic field → 0.6 Gauss.
- A common hand-held magnet → 100 Gauss.
- The strongest steady magnetic fields achieved so far in the laboratory → 4.5×10^5 Gauss.
- Surface field of magnetars → 10^{15} Gauss.
- Heavy Ion Collisions, the strongest magnetic field ever achieved in the laboratory → 10^{18} Gauss $\approx m_\pi^2$
($m_\pi = 135$ MeV).

Magnetic fields generated in HIC.



D. E. Kharzeev, Prog.Part.Nucl.Phys. 75 (2014).

Estimation of magnetic fields strength

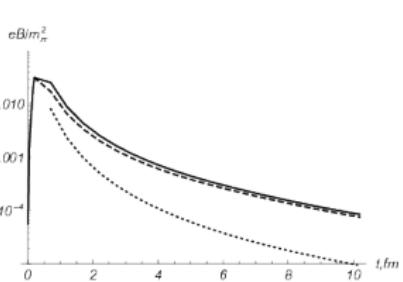


V. Skokov *et al.*, Int.J.Mod.Phys.

A24 (2009) 5925-5932.

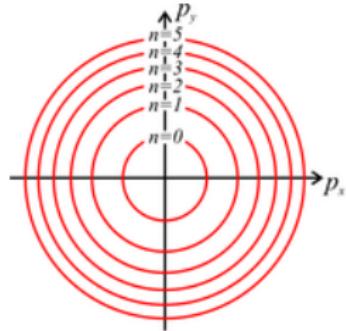
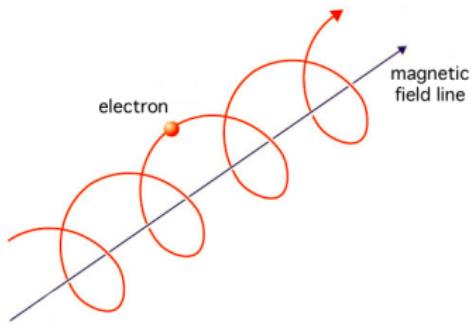
D. Kharzeev *et al.*, Nucl.Phys. A803

(2008) 227-253.



K. Tuchin, arXiv:1508.06925

Charged particles in presence of magnetic fields.



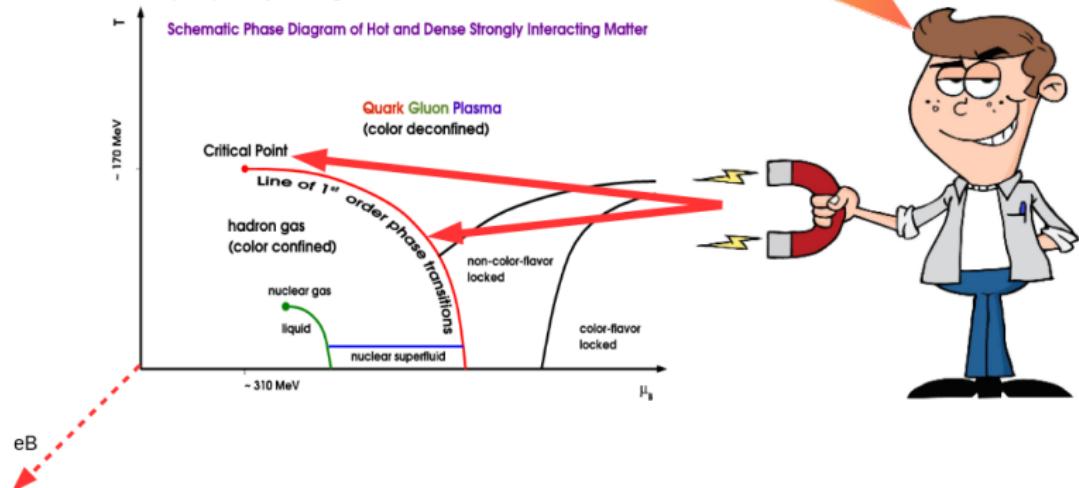
- The physical system loses isotropy.
- Reduction from 3D to 2D.

QCD phase transition & $eB \neq 0$.

Is the T_c modified by strong magnetic fields?

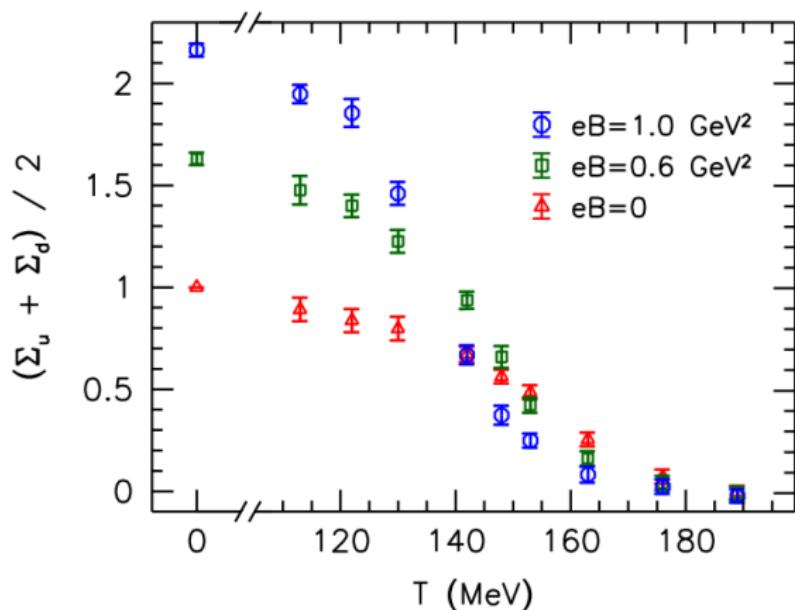
Does the CEP move due to the strength of the magnetic field?

PoS HIGH-PTLHC08 (2008) 027 by T. Csorgo.

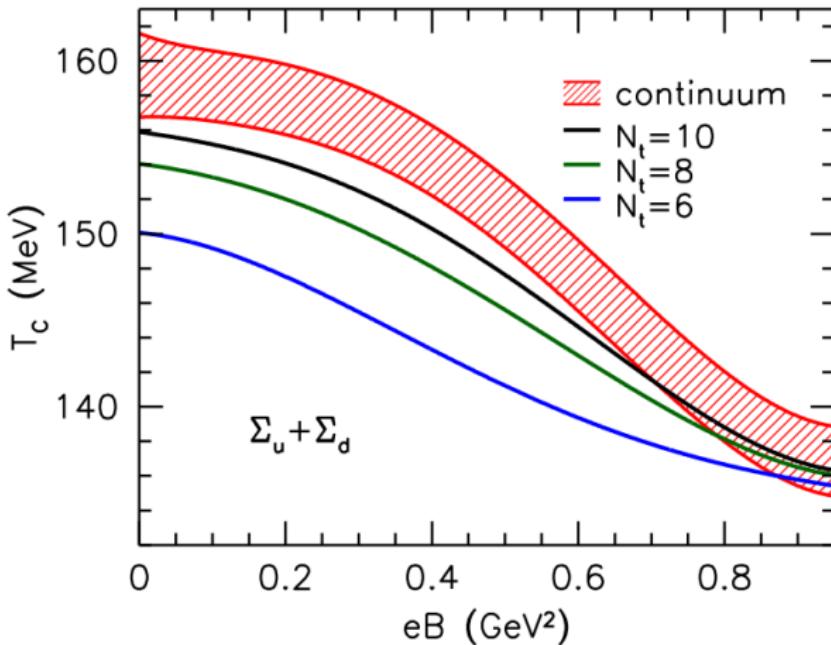


The Magnetized QCD Phase Diagram

Magnetic catalysis & Inverse Magnetic catalysis

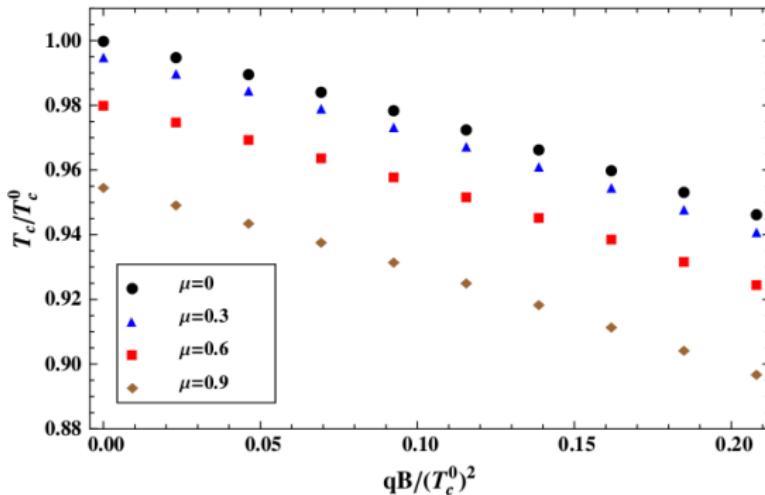


Magnetic catalysis & Inverse Magnetic catalysis



Some theoretical models.

Linear Sigma Model Coupled to quarks.



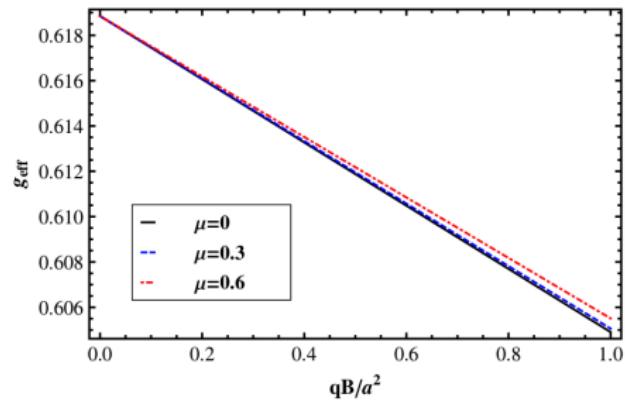
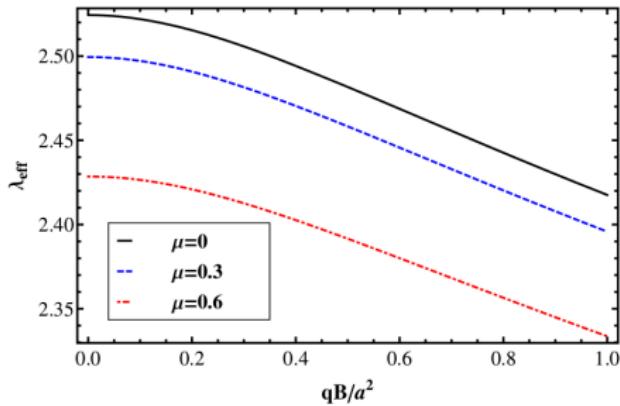
A. Ayala *et al.*, Phys.Rev. D92 (2015) 9, 096011

Beyond mean field approximation.

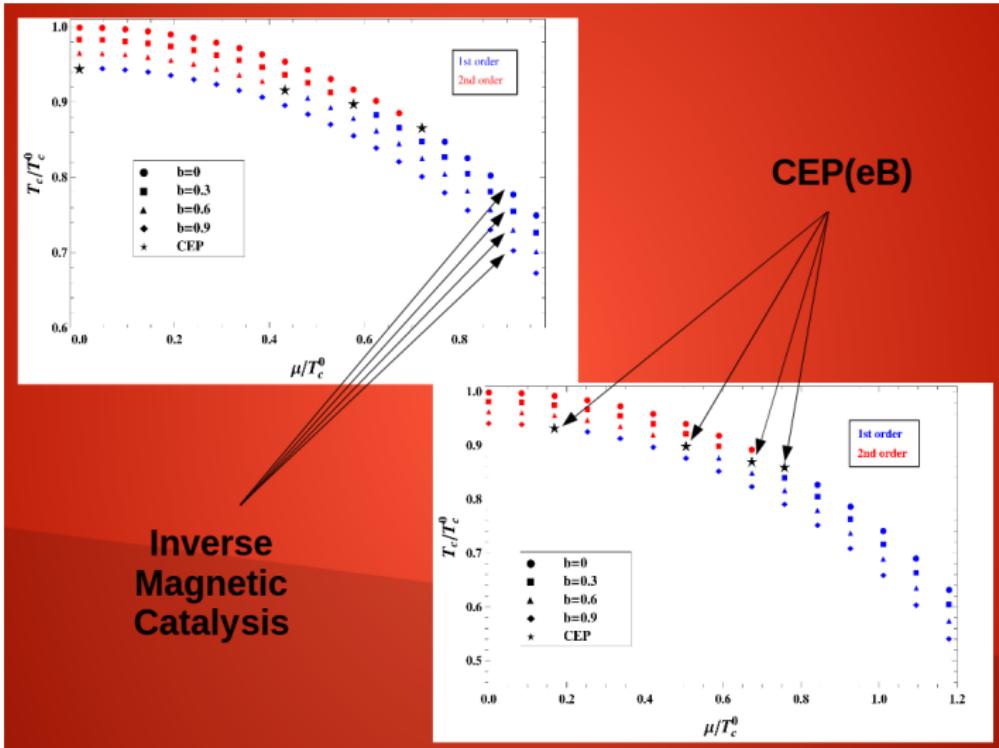
- More than one-loop correction.
- First correction to coupling constants.

Some theoretical models.

First order correction to coupling constants.



Magnetized Effective QCD Phase Diagram.



Physics behind (inverse) magnetic catalysis.

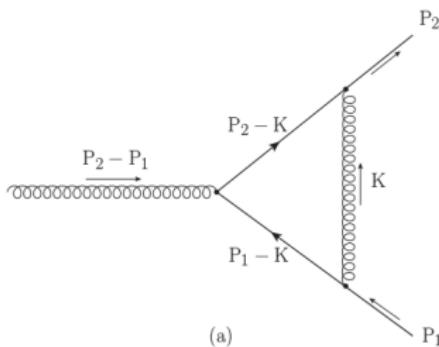
Is the coupling constant the key to understand the (inverse) magnetic catalysis?

- Working within QCD.
- In the framework of perturbation theory.
 - High temperature ($T^2 > eB$).
 - High virtuality ($q^2 > eB$).

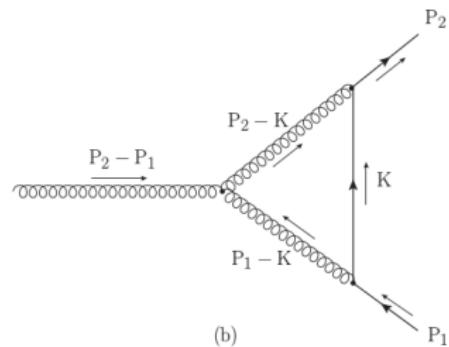
Physics behind (inverse) magnetic catalysis.

The thermo-magnetic correction to the quark-gluon vertex in the presence of a weak magnetic field was computed.

The vacuum one-loop quark-gluon vertex correction at zero temperature in the presence of magnetic field was computed.



QED-like contribution.

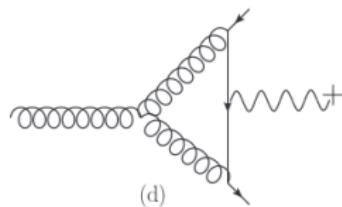
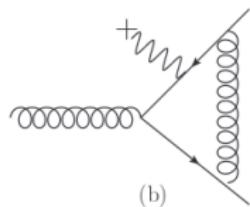
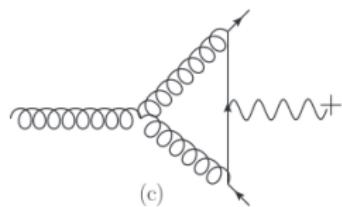
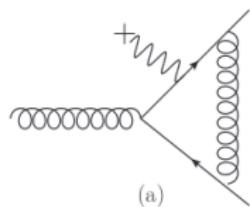


Pure QCD contribution.

Weak field approximation.

Magnetic fields modify the behaviour of Quarks.

$$S_B(K, m_f) = \frac{(m_f - K)}{K^2 + m_f^2} - i \frac{\gamma_1 \gamma_2 (qB)(m_f - K_{\parallel})}{(K^2 + m_f^2)^2}$$
$$+ \frac{2(qB)^2 K_{\perp}^2}{(K^2 + m_f^2)^4} \left[(m_f - K_{\parallel}) + \frac{K_{\perp}(m_f^2 + K_{\parallel}^2)}{K_{\perp}^2} \right].$$



HTL.

thermo-magnetic correction of g .

$$\delta\Gamma_{\mu}^{(a)} = ig^2(C_F - \frac{C_A}{2})(qB)T \sum_n \int \frac{d^3k}{(2\pi)^3} \gamma_{\nu} \left[\gamma_1 \gamma_2 K_{\parallel} \gamma_{\mu} K \tilde{\Delta}(P_2 - K) + K \gamma_{\mu} \gamma_1 \gamma_2 K_{\parallel} \tilde{\Delta}(P_1 - K) \right] \gamma_{\nu} \Delta(K) \tilde{\Delta}(P_2 - K) \tilde{\Delta}(P_1 - K)$$

$$\delta\Gamma_{\mu}^{(b)} = 2ig^2 \frac{C_A}{2}(qB)T \sum_n \int \frac{d^3k}{(2\pi)^3} \left[K \gamma_1 \gamma_2 K_{\parallel} \gamma_{\mu} - 2\gamma_{\nu} \gamma_1 \gamma_2 K_{\parallel} \gamma_{\nu} K_{\mu} + \gamma_{\mu} \gamma_1 \gamma_2 K_{\parallel} K \right] \tilde{\Delta}(K)^2 \Delta(P_1 - K) \Delta(P_2 - K) \tilde{\Delta}(P_1 - K)$$

with

$$\Delta(K) = \frac{1}{\omega_n^2 + k^2}$$

$$\tilde{\Delta}(K) = \frac{1}{\tilde{\omega}_n^2 + k^2 + m^2}$$

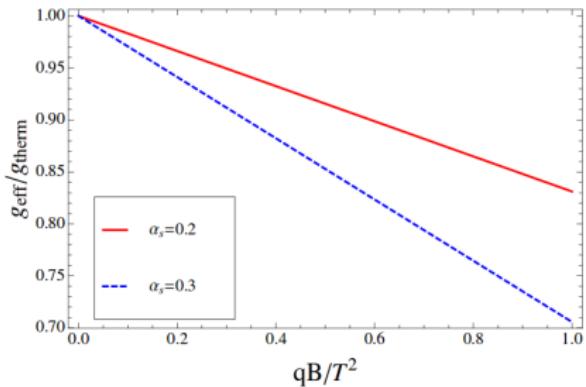
HTL .

Result

$$\begin{aligned}\delta\Gamma_\mu &= \delta\Gamma_\mu^{(a)} + \delta\Gamma_\mu^{(b)} \\ &= g^2 C_F \left(\frac{qB}{6\pi^2 T^2} \right) \left[\log(2) - \frac{\pi T}{2m} \right] \vec{\gamma}_\parallel \Sigma_3\end{aligned}$$

Then

$$g_{eff} = g \left[1 - \frac{m^2}{T^2} + \delta\Gamma \right]$$



Vacuum

Magnetic correction of g .

$$\begin{aligned} \delta\Gamma_{(a)}^\mu &= g^2(qB) \left(C_F - \frac{C_A}{2} \right) \int \frac{d^4k}{(2\pi)^4} \frac{1}{k^2} \\ &\times \left[\gamma^\nu \frac{\not{p}_2 - \not{k}}{(p_2 - k)^2} \gamma^\mu \frac{\gamma_1 \gamma_2 (\not{p}_1 - \not{k})_\parallel}{(p_1 - k)^4} \gamma_\nu + \gamma^\nu \frac{\gamma_1 \gamma_2 (\not{p}_2 - \not{k})_\parallel}{(p_2 - k)^4} \gamma^\mu \frac{\not{p}_1 - \not{k}}{(p_1 - k)^2} \gamma_\nu \right] \end{aligned}$$

$$\begin{aligned} \delta\Gamma_{(b)}^\mu &= -2g^2(qB) \frac{C_A}{2} \int \frac{d^4k}{(2\pi)^4} \frac{1}{k^4} \left[g^{\mu\nu} (2p_2 - p_1 - k)^\rho \right. \\ &+ g^{\nu\rho} (2k - p_2 - p_1)^\mu + g^{\rho\mu} (2p_1 - k - p_2)^\nu \left. \right] \gamma_\rho \frac{\gamma_1 \gamma_2 \not{k}_\parallel}{(p_2 - k)^2 (p_1 - k)^2} \gamma_\nu \end{aligned}$$

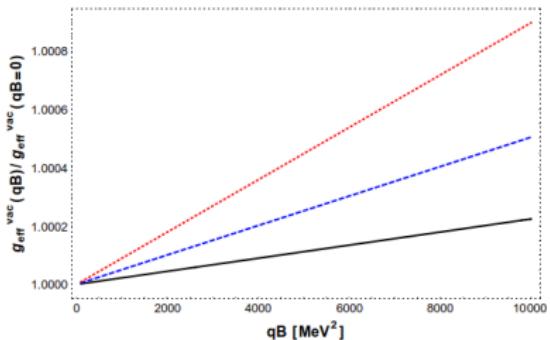
Vacuum

Result

$$\begin{aligned}\delta\Gamma^\mu &= \delta\Gamma_{(a)}^\mu + \delta\Gamma_{(b)}^\mu \\ &= -g^2 \frac{1}{3\pi^2} \left(\frac{q \vec{\Sigma} \cdot \vec{B}}{Q^2} \right) \left[[1 - \log(4)]C_F - \frac{[7 + 3\log(4)]}{10}C_A \right]\end{aligned}$$

Then

$$g_{eff} = g [1 + \delta\Gamma]$$



Conclusions.

- Our results show that the geometrical effect produced by the magnetic field at high temperature, whereby quarks and anti-quarks get closer on average, is accompanied by the decrease of their effective interaction due to the asymptotic freedom.
- In contrast, at $T = 0$ such geometrical effect does not take place. This because the color charge associated to gluons produces a kind of screening of the color charge associated to quarks.

Many Thanks!!!